

Credit: NASA/Swift Dana Berry



UNIVERSITÉ
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Unveiling the structure of ultra-dense matter

*Synergy between nuclear physics and
gravitational waves*

*Francesca Gulminelli - Université de
Caen-Normandie*

Congrès Général des 150 ans de la SFP, 3-7 Juillet, 2023

plan

1. Neutron stars and gravitational waves
2. Gravitational waves and the equation of state of dense matter
3. The contribution of nuclear physics
4. Future perspectives

plan

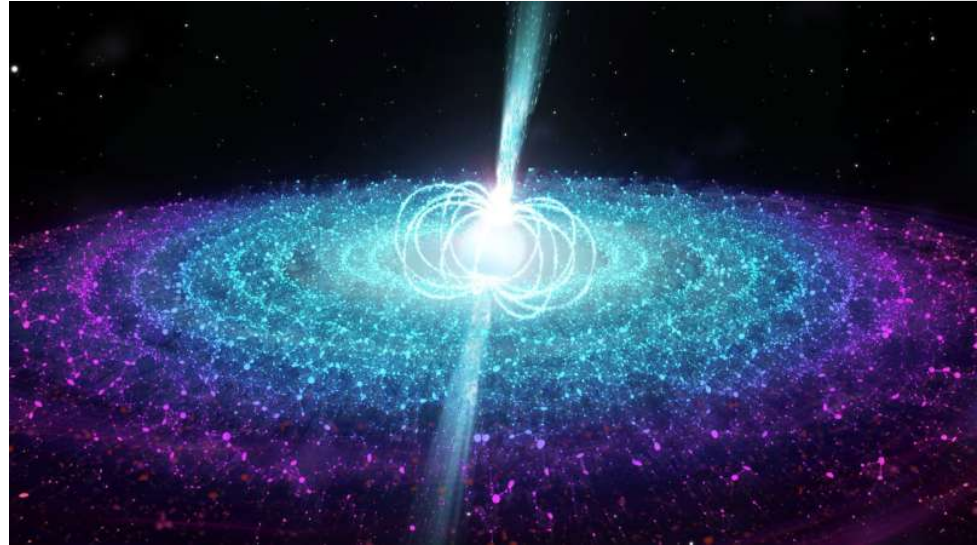
1. **Neutron stars and gravitational waves**
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Neutron stars

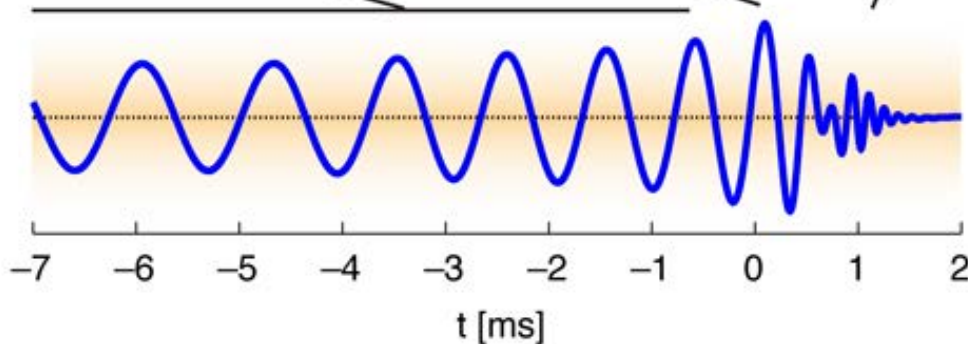
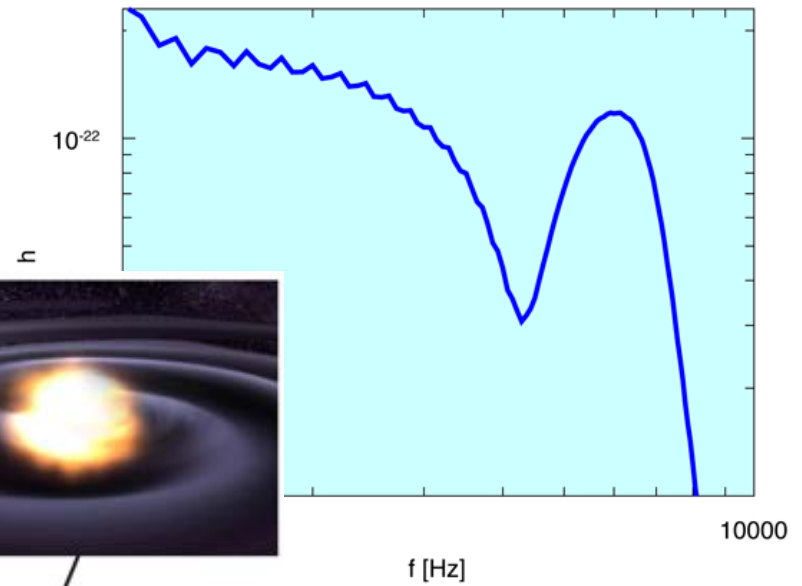
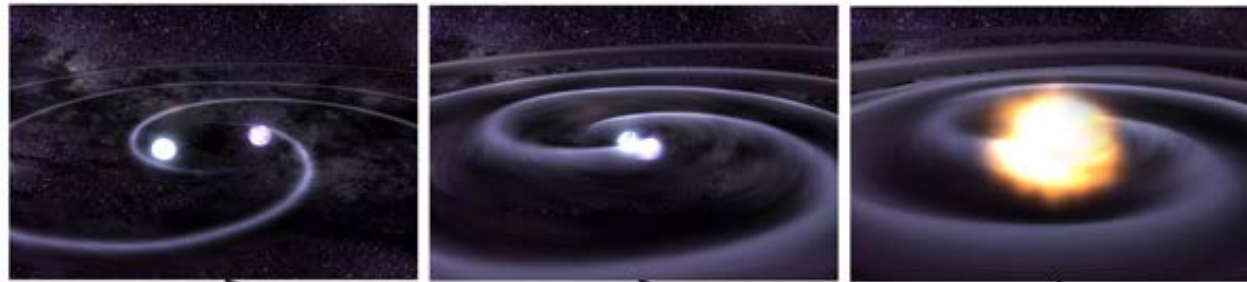
record list

Credit: ICRAR/Univ. Amsterdam.

- The objects which spin the fastest $\nu=716$ Hz
=> $V_{\text{equator}} \sim c/4$
- The highest speed of the galaxy $v=1083$ km/s
- The most intense magnetic fields
 $H > 10^{14}$ Gauss
- The densest matter of the universe
 $\rho \sim 10^{14}$ g/cm³
- The only place where:
 - neutrinos can be trapped
 - quarks might be deconfined



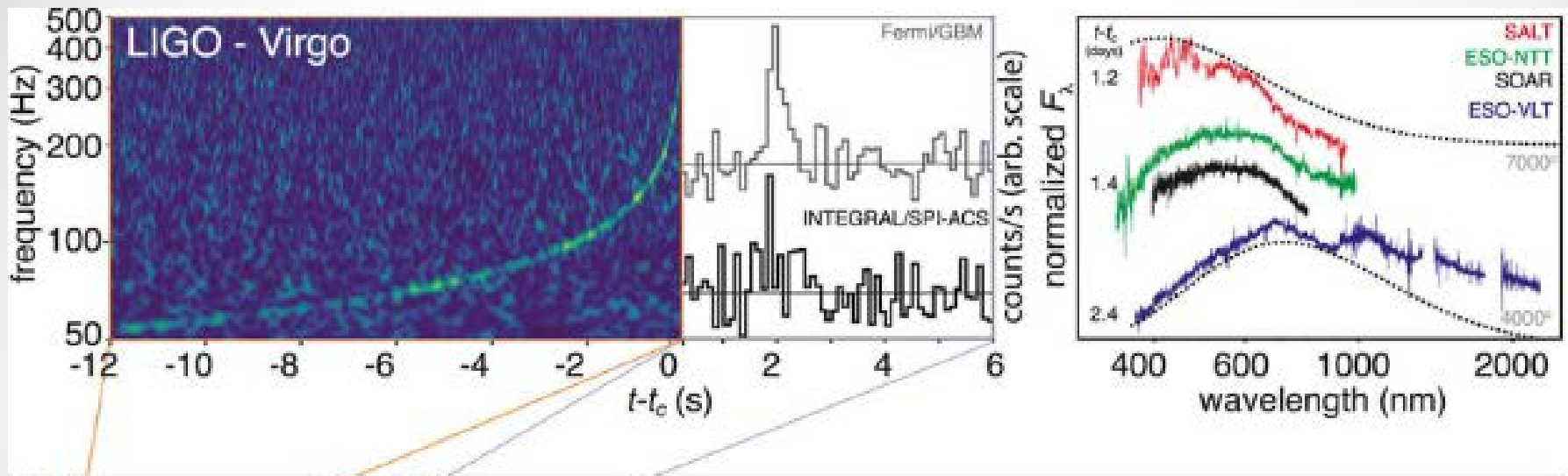
Neutron star binaries: GW sources



$$h_{ij}^{TT} = \frac{2G}{c^4 r} P_{ij}{}^{kl} \ddot{Q}_{kl}$$

Source: NASA & A. Stonebraker, adapted from PRL 104, 141101

GW170817 by LVC



<https://www.ligo.org/detections.php>



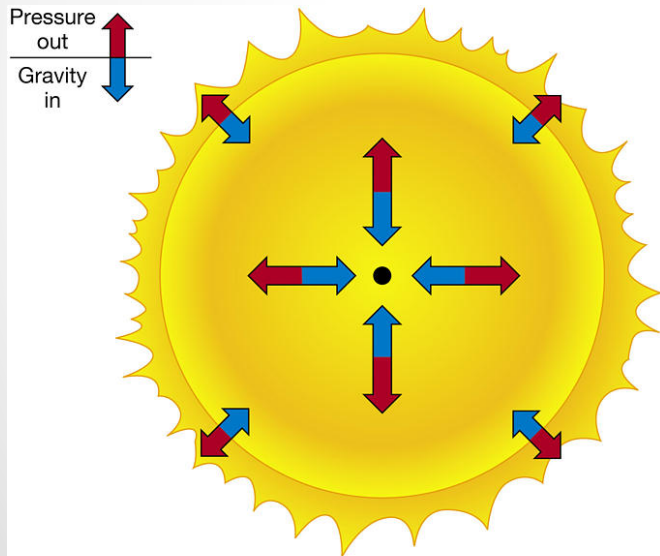
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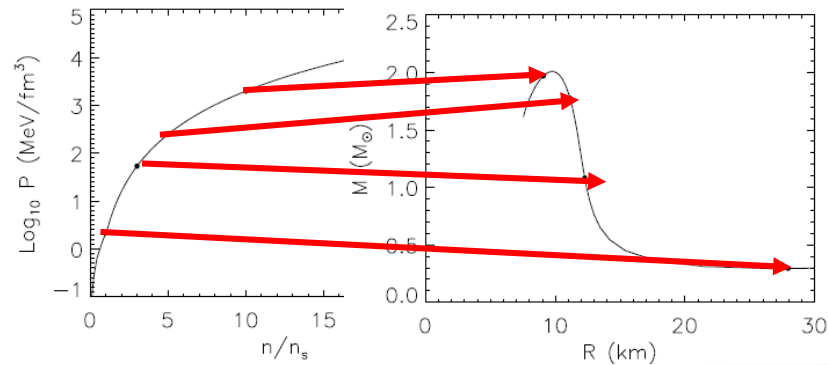
Modeling (neutron) stars: hydrostatics

- Equilibrium of a self-gravitating object
=> Tolman Oppenheimer Volkoff (1939):

$$\frac{dP(\rho)}{dr} = -\frac{G}{r^2} \left[\rho(r) + \frac{P(\rho)}{c^2} \right] \left[m(r) + 4\pi r^3 \frac{P(\rho)}{c^2} \right] \left[1 - \frac{2Gm(r)}{rc^2} \right]^{-1}$$



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Equilibrium radius for each mass

Only depends on the EoS $P(\rho)$

Modeling (neutron) stars: hydrostatics

- Influence of a second body => Thorne and Campolattaro (1967):

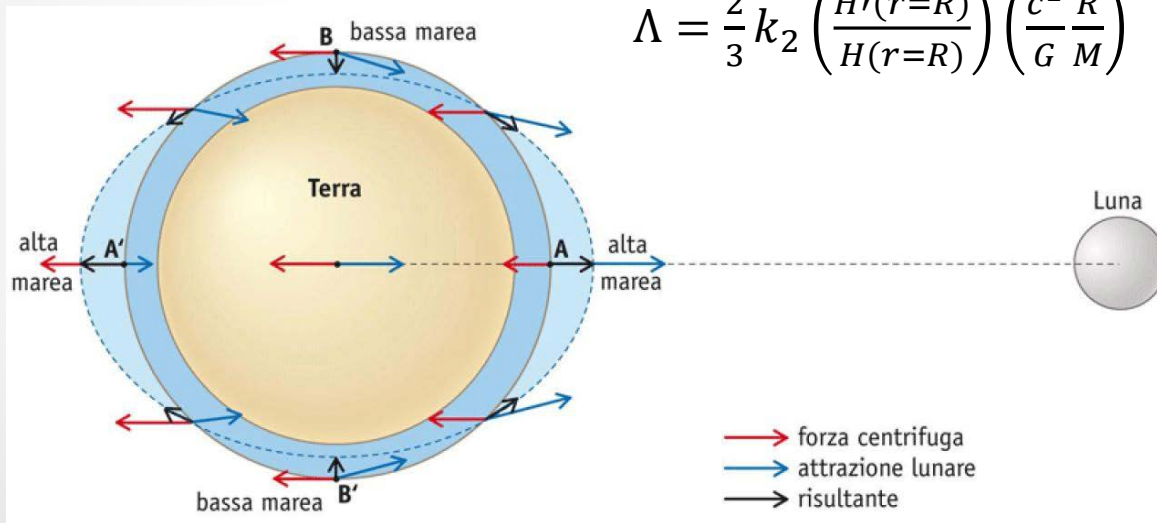
$$\frac{d^2 H(r)}{dr^2} + \frac{dH(r)}{dr} \left[\frac{2}{r} + e^{\lambda(\mathbf{P}(\rho))} \left(\frac{2m(r)}{r^2} + 4\pi r (\mathbf{P}(\rho) - \rho(r)) \right) \right] + H(r) Q(\mathbf{P}(\rho)) = 0$$

$$\forall \rho_c$$

$$\Lambda = \frac{2}{3} k_2 \left(\frac{H'(r=R)}{H(r=R)} \right) \left(\frac{c^2 R}{G M} \right)^5$$

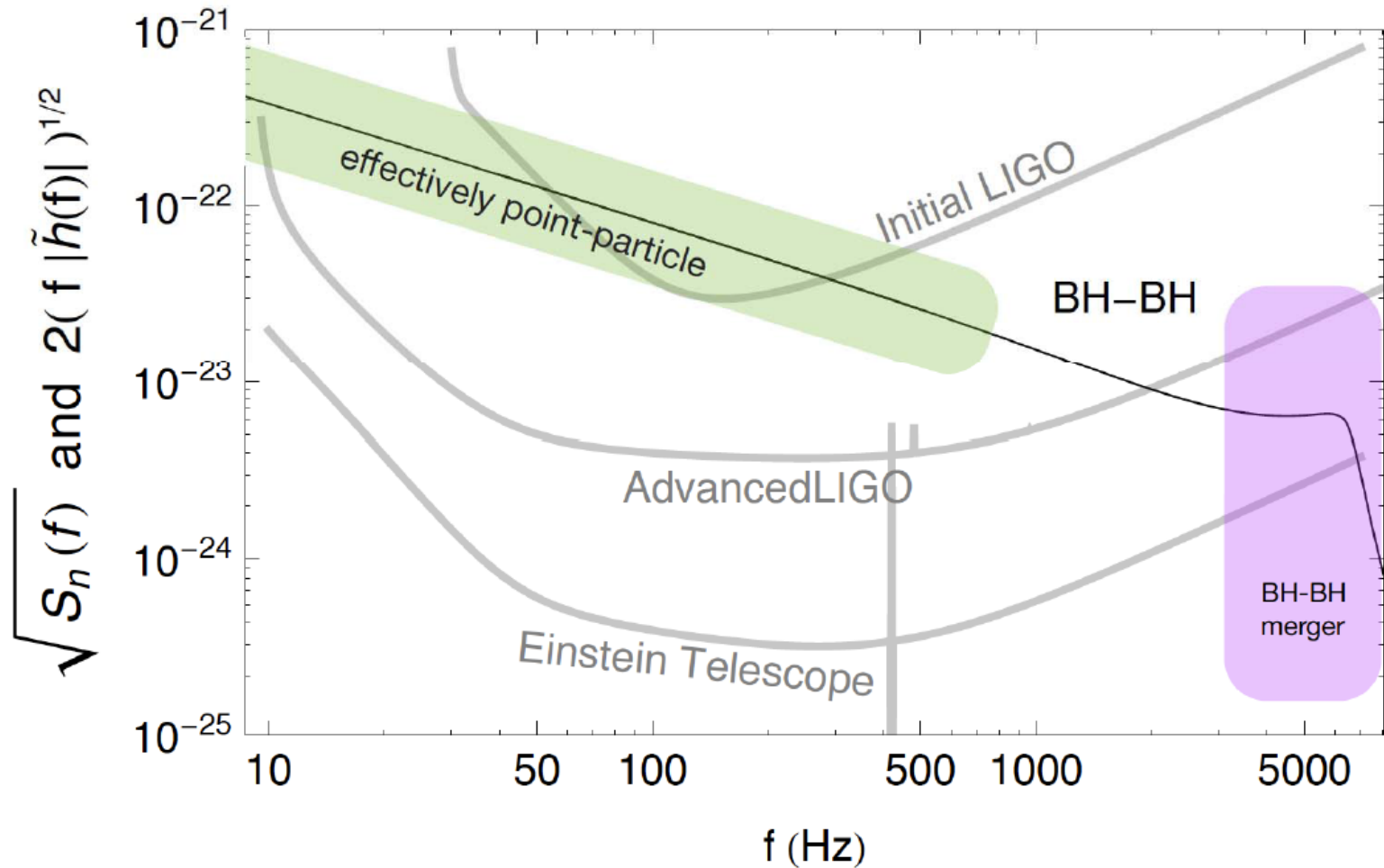
Tidal polarisability

$$\bar{\bar{Q}}^{(i)} = -\Lambda_i \bar{\bar{\mathcal{E}}}^{(j)}$$

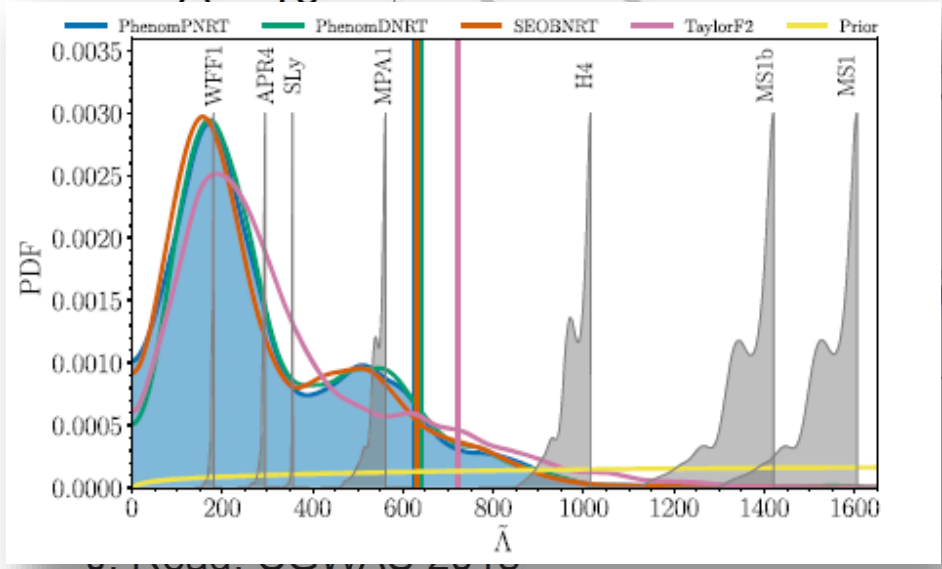
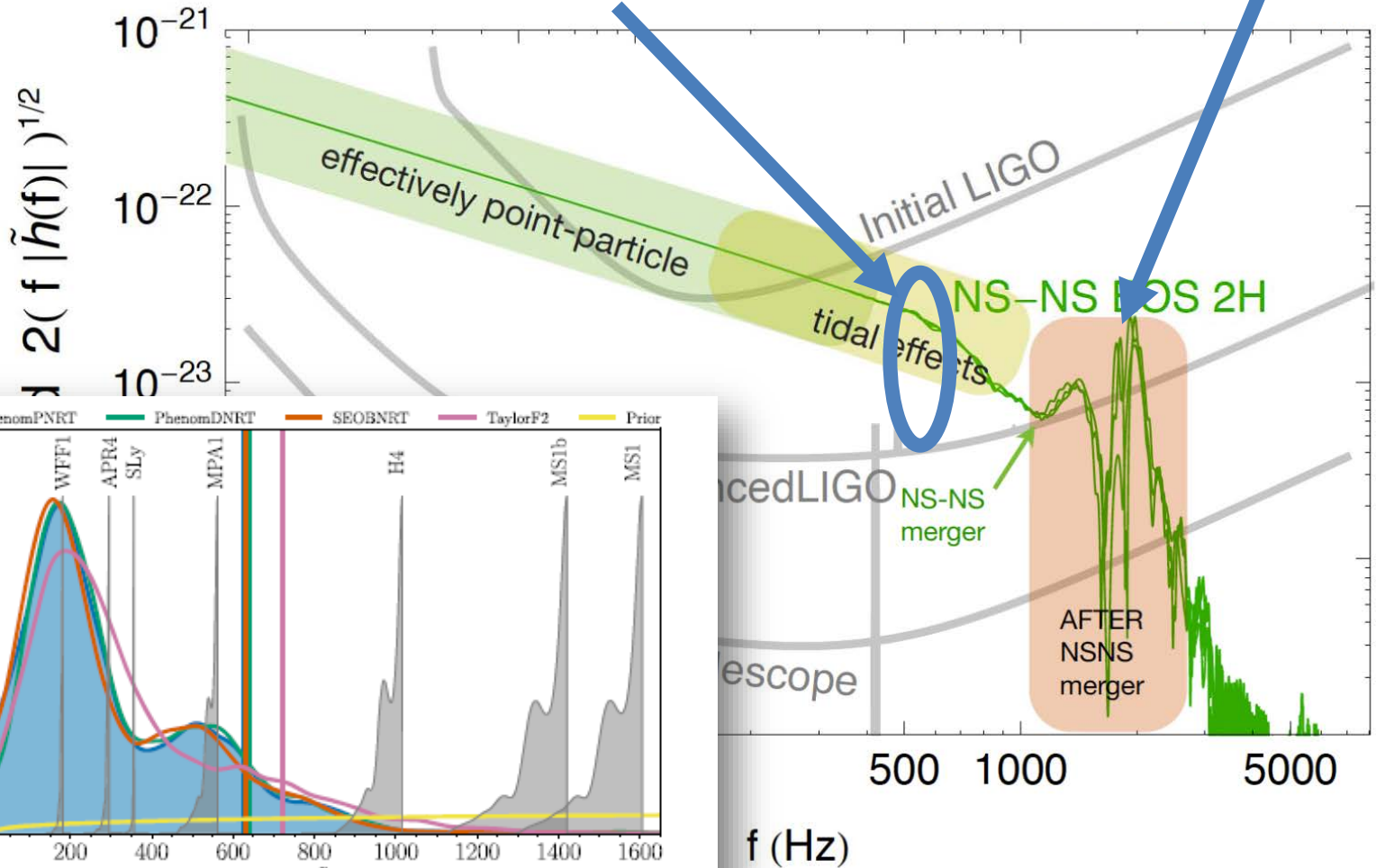
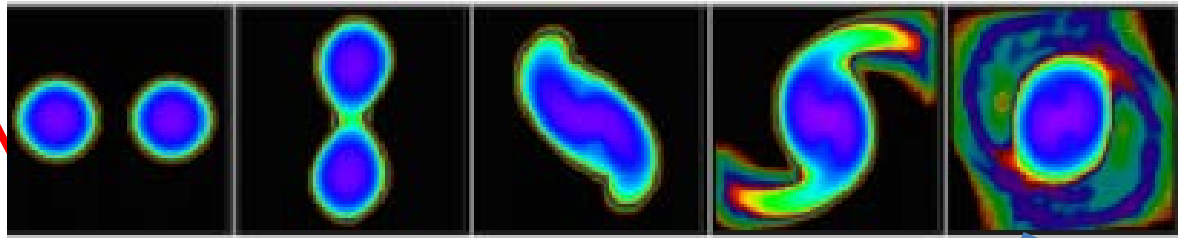


Only depends on the EoS $P(\rho)$

Spectrum of **BBH** inspiral, scale to 1.35-1.35, 45 Mpc



Spectrum of $\tilde{h}(f)$

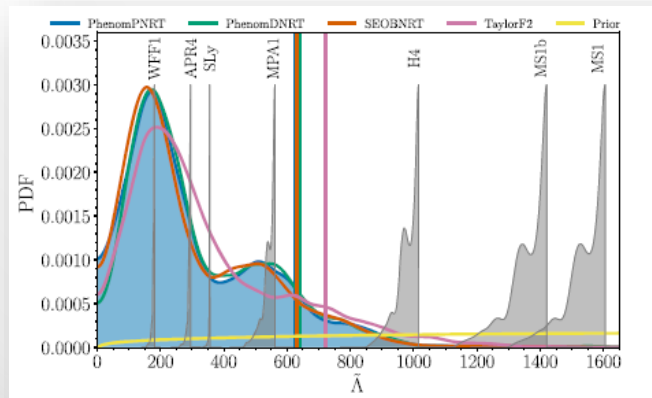
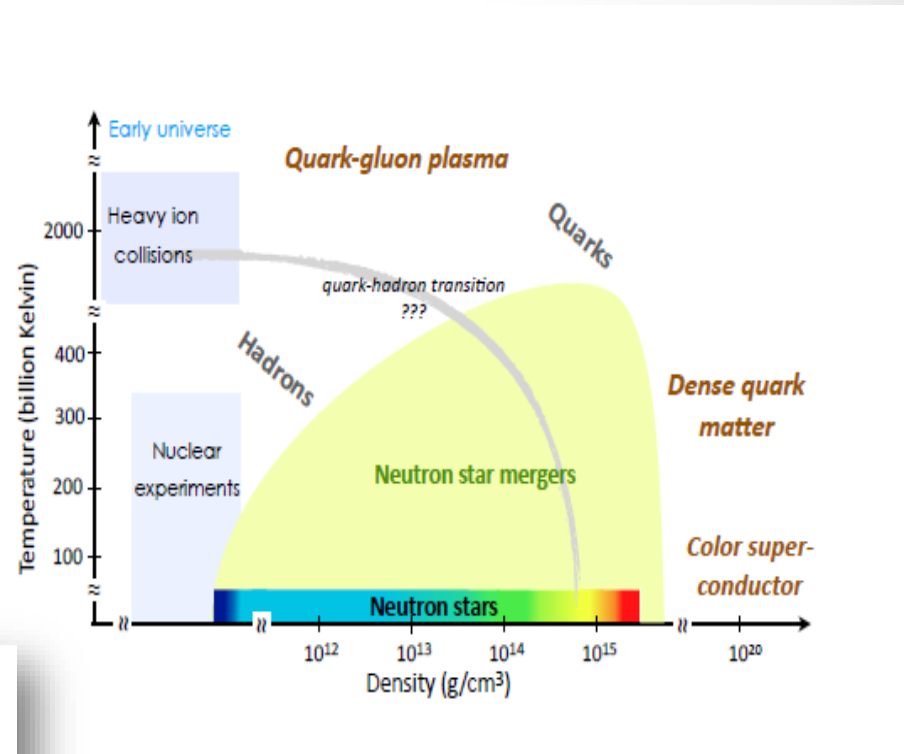


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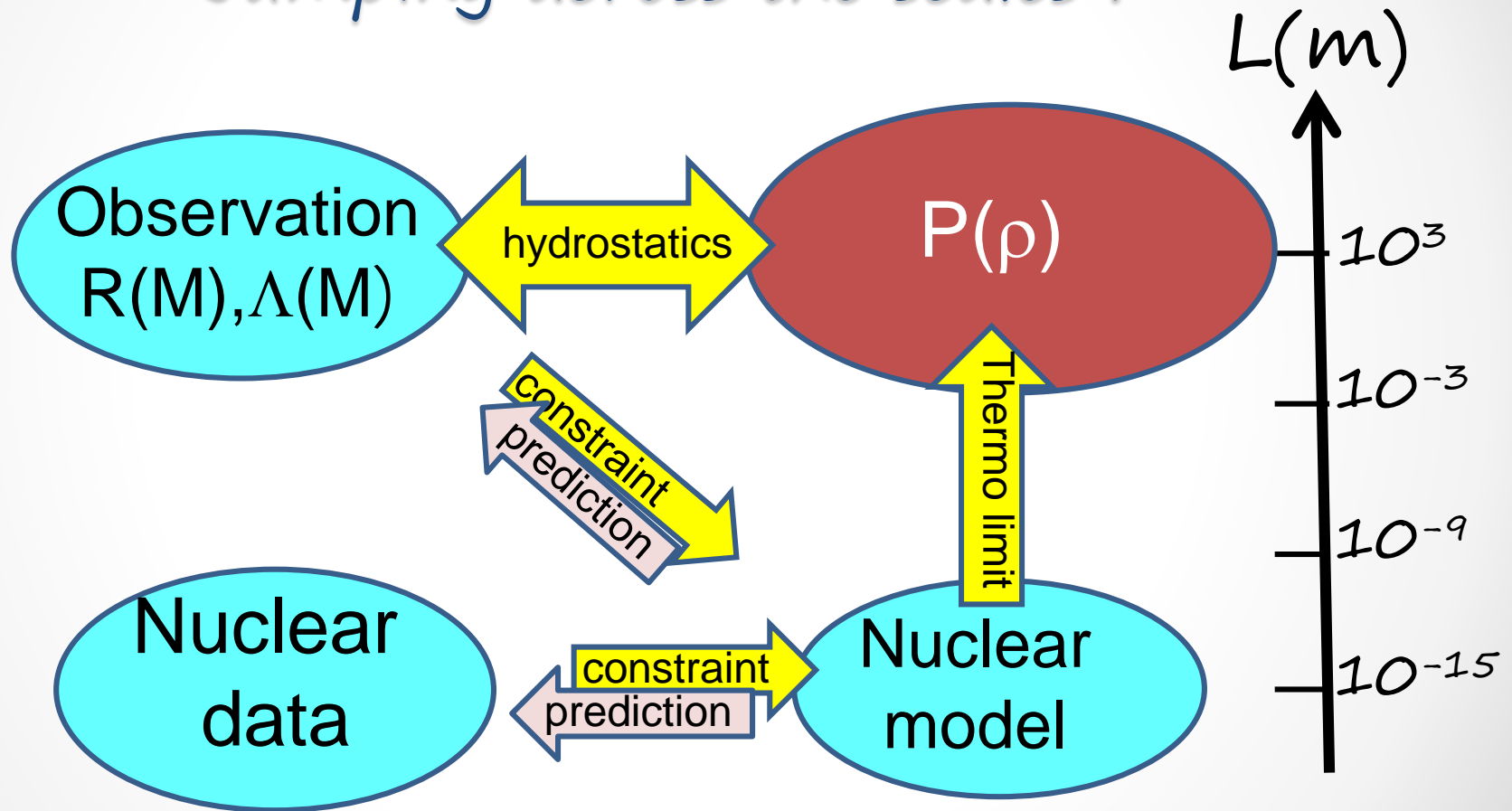
Neutron star EoS and nuclear physics

- $M \sim 1.4 M_{\odot}$, $R \sim 12$ km \Rightarrow density as the one inside the atomic nucleus \Rightarrow nuclear degrees of freedom (hadrons or quarks) in strong interaction
- The strong interaction is described by the standard model of particle physics, but no solution exists for dense matter
- Effective models of neutrons and proton



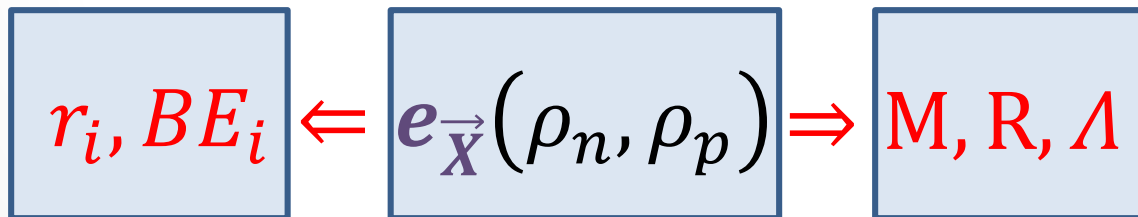
JCAPO3(2020)050
 Science case for the Einstein Telescope

Jumping across the scales ?



Jumping across the scales: the syllabus

- An effective model of the relevant densities (neutrons and protons) gives a prediction for the EoS $P(\rho) = -\rho^2 \left. \frac{\partial e(\rho_n, \rho_p)}{\partial \rho} \right|_{\mu_L=0}$
- A flexible analytic representation $e_{\vec{X}}(\rho_n, \rho_p)$: the variation of the parameter set \vec{X} allows reproducing the effective models and interpolating among them ~ 15 parameters
- The X_i variation explores the equation of state space compatible with the hypothesis of a matter of neutrons and protons



Laboratory
observables

EoS analytic
representation

Astronomical
observables

Bayesian Inference

$$P(\vec{X}|\vec{f}) = \frac{P(\vec{X}) \prod_i P(f_i|\vec{X})}{P(\vec{f})}$$

- f_1 . nuclear data
- f_2 . ab-initio theory

} **Nuclear physics**

- f_3 . max.mass (radio)
- f_4 . tidal polarisability (GW)
- f_5 . radius (X-ray)

} **Astrophysics**

(1) Huang et al, 2016 AME mass table, Angeli&Marinova, ADNDT 2013

(2) χ EFT Drischler et al PRC 2016

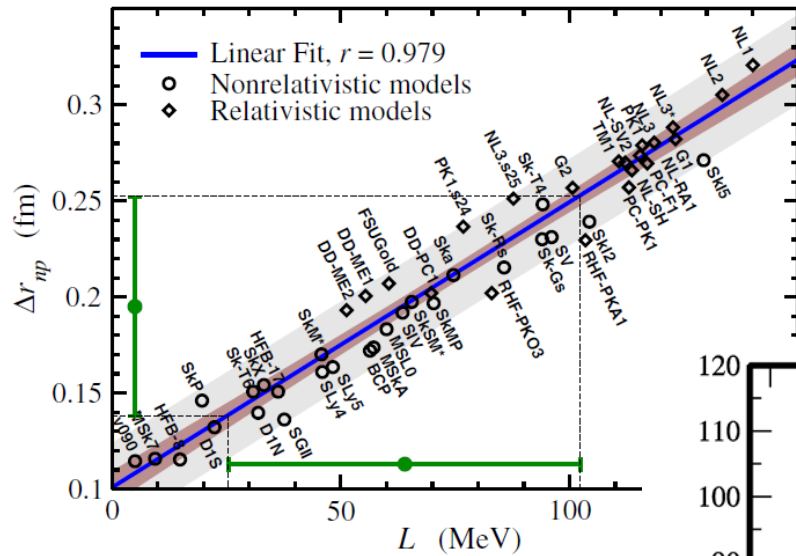
(3) PSR J0348+0432 $M=2.01\pm 0.04 M_{\odot}$

(4) GW170817 $\tilde{\Lambda}(M)$ LVK

(5) PSR J0030+0451, PSR J0740+6620 NICER

Laboratory experiments

example: neutron skin of ^{208}Pb

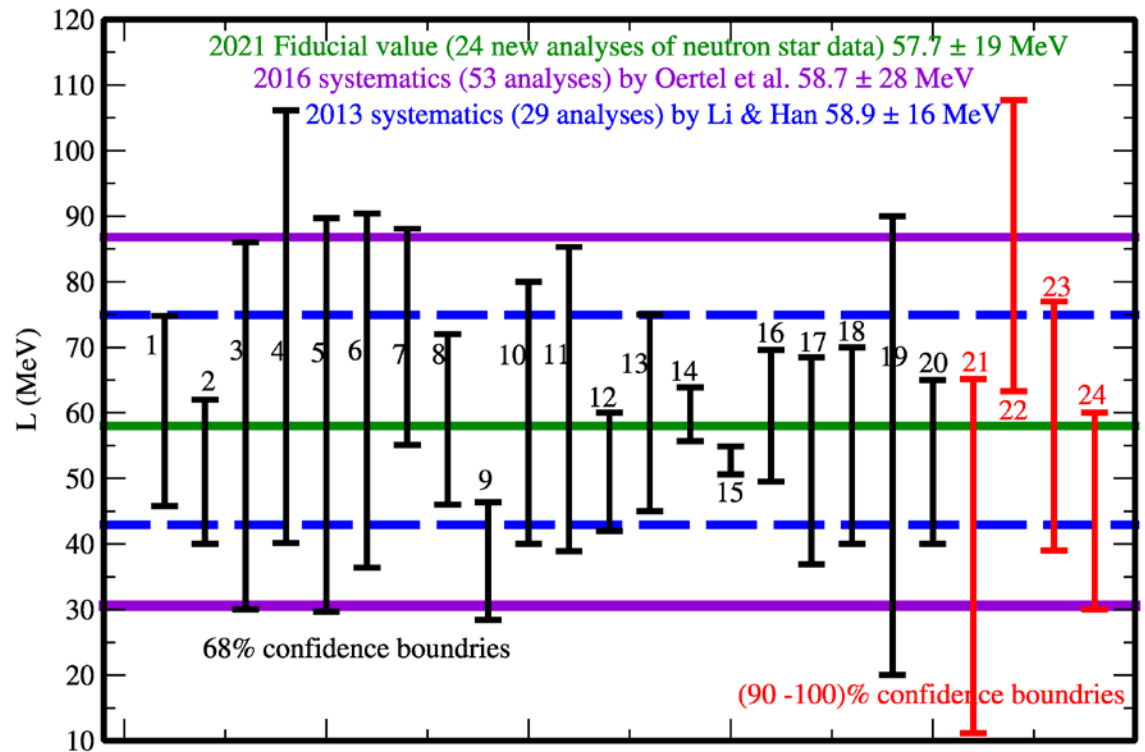


X.Roca-Maza et al PRL 2011

$$\vec{X} = \{E_0, K, E_{sym}, L, K_{sym}, \dots\}$$

$$P(f_i = \Delta r | \vec{X}) \propto e^{-\frac{(\Delta r - \Delta r(\vec{X}))^2}{\sigma_{\Delta r}^2}}$$

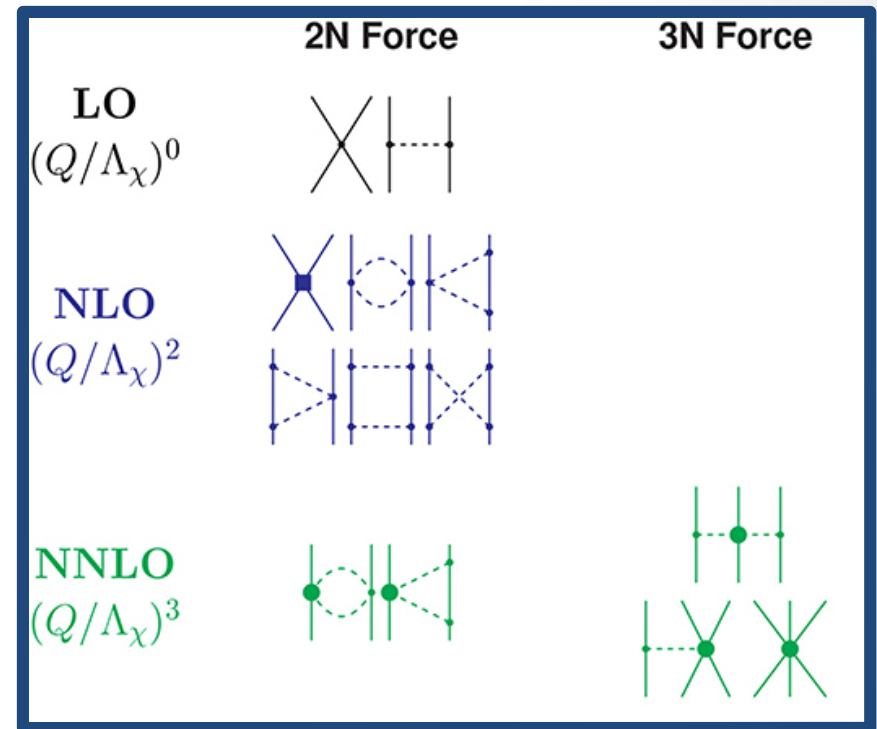
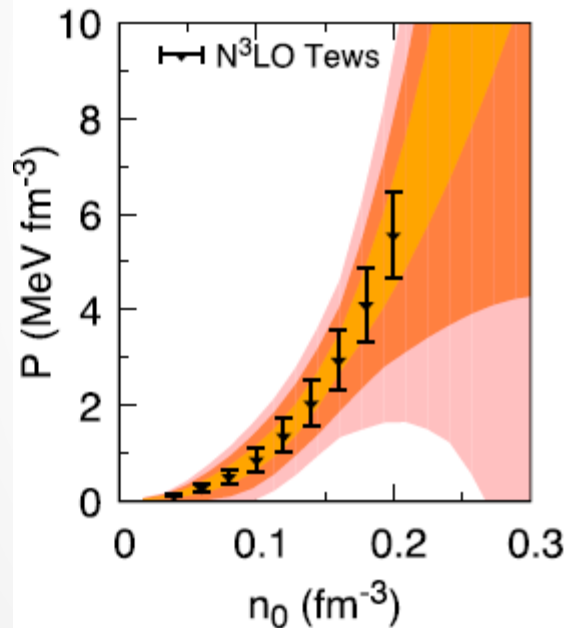
B.An Li et al, Universe 2021, 7(6), 182



Chiral perturbation theory

- The « ab-initio » nuclear theory
- Perturbative expansion : controlled truncation errors
- Moment expansion! Only valid at low density

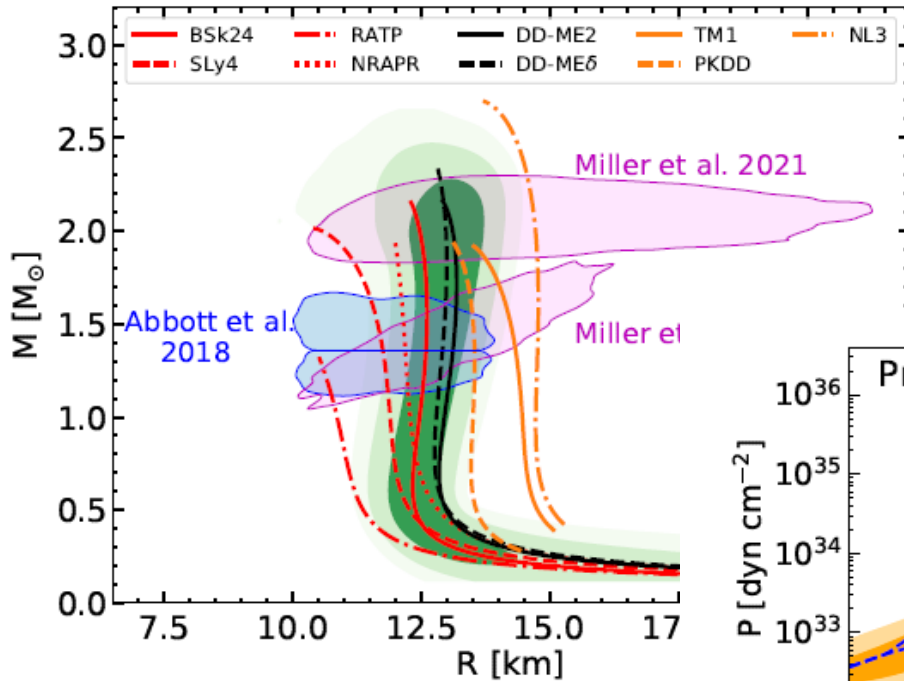
Machleidt R., Int J Mod Phys. (2017)



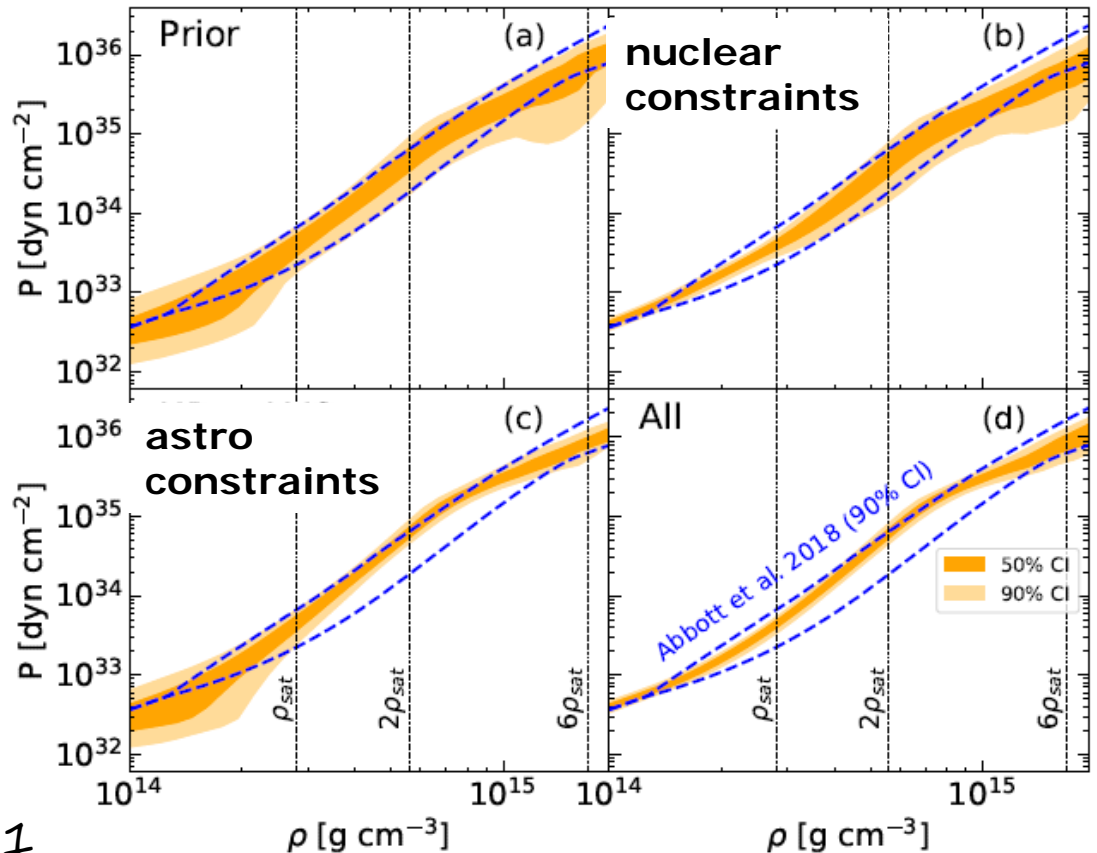
I. Tews, T. Krüger, K. Hebeler, and A. Schwenk, [Phys. Rev. Lett. 110, 032504 \(2013\)](#).

C. Drischler, K. Hebeler, and A. Schwenk, [Phys. Rev. C 93, 054314 \(2016\)](#).

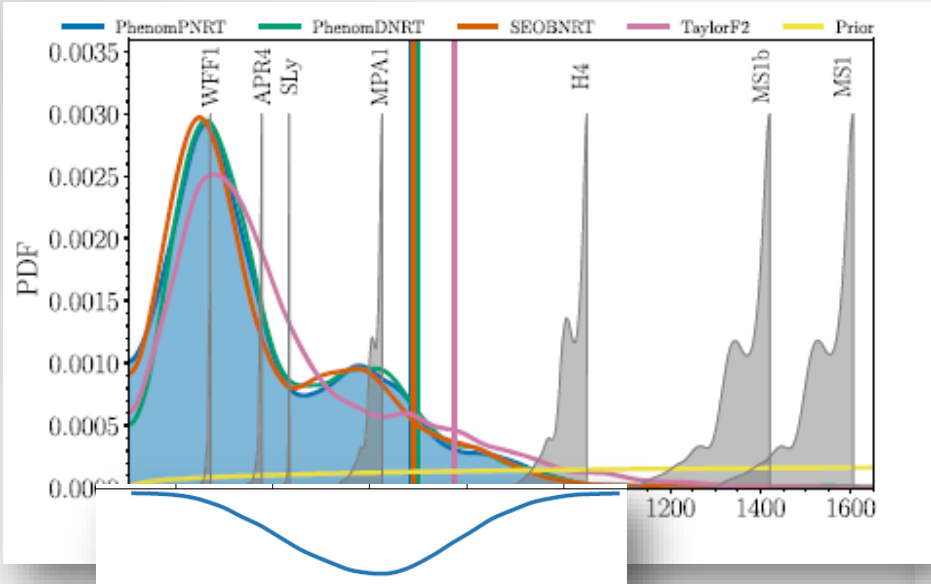
The nuclear physics predictions



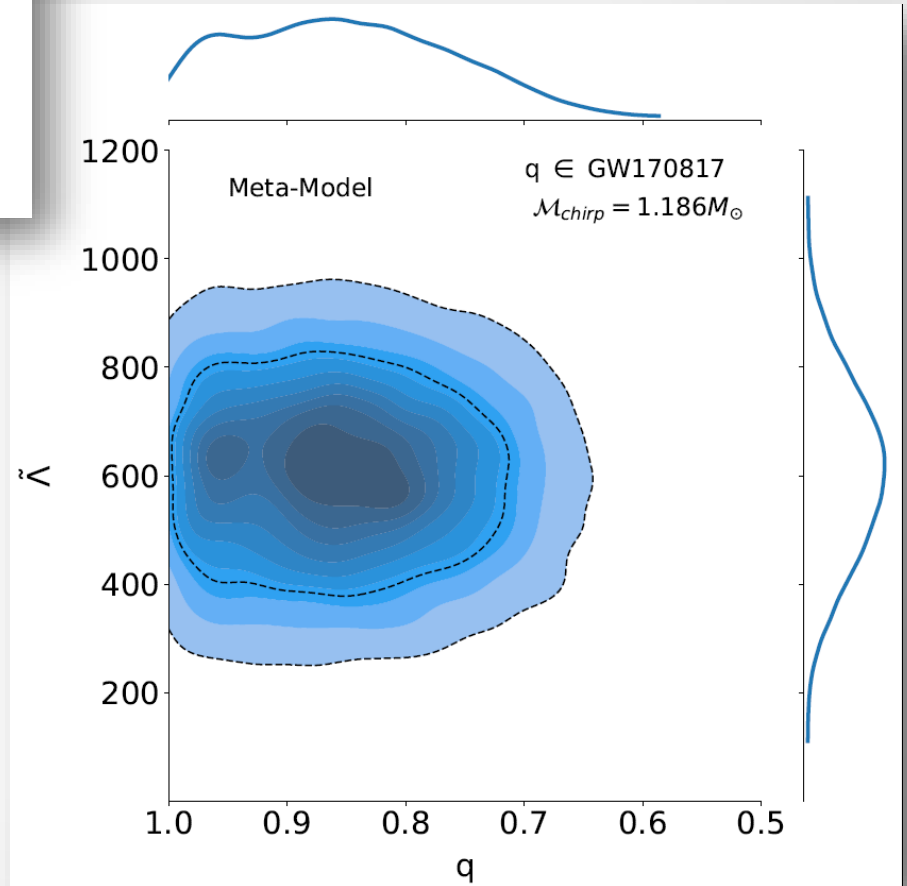
- A neutrons and protons composition is compatible with the observations
- Many models can be excluded



The nuclear physics predictions



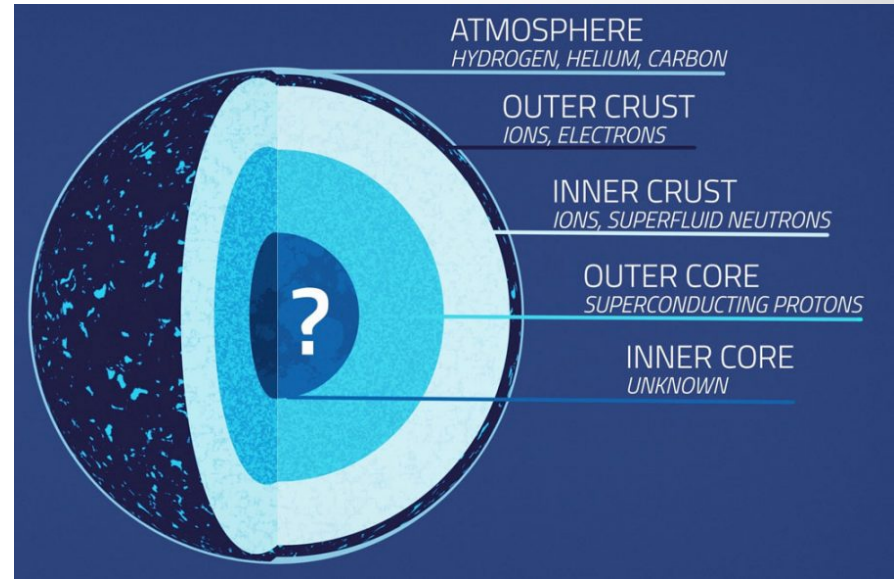
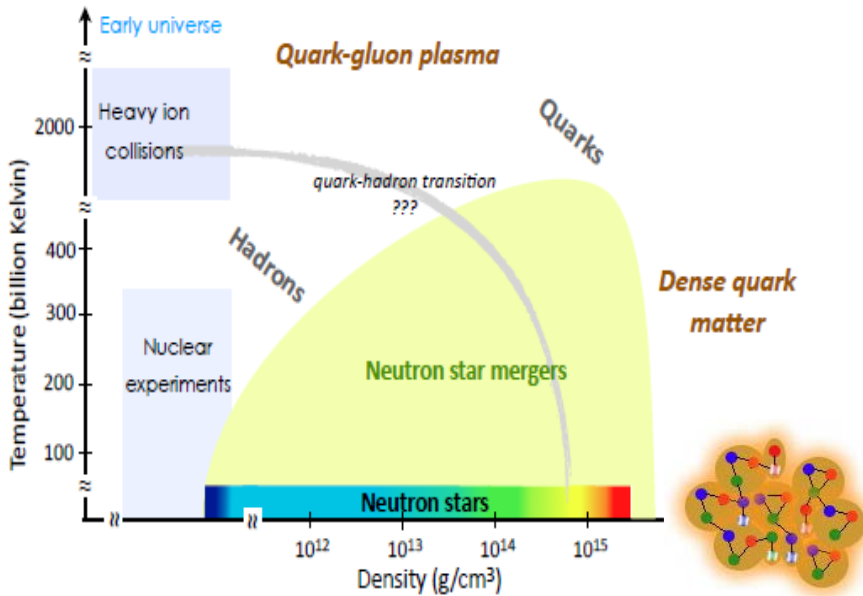
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Quarks in the core of neutron stars?

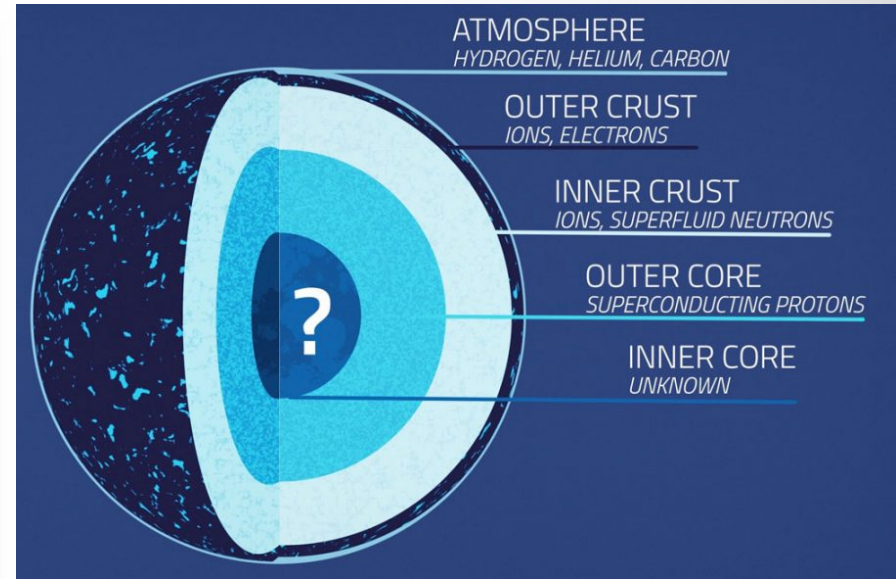
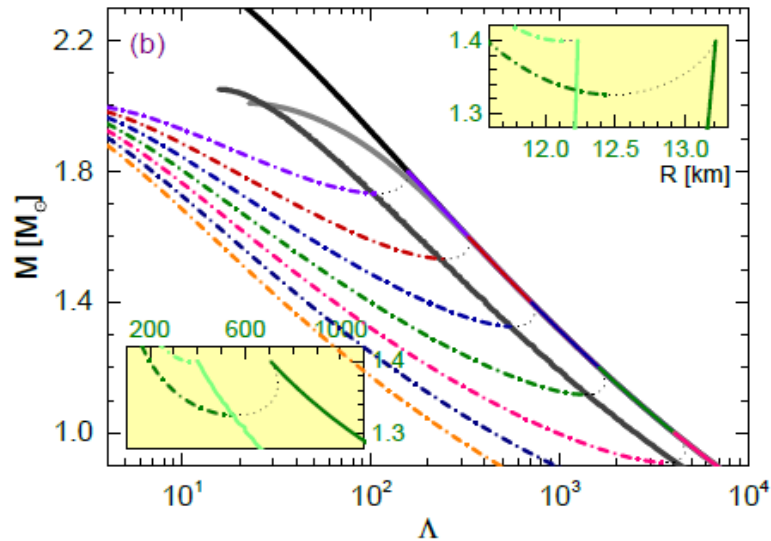
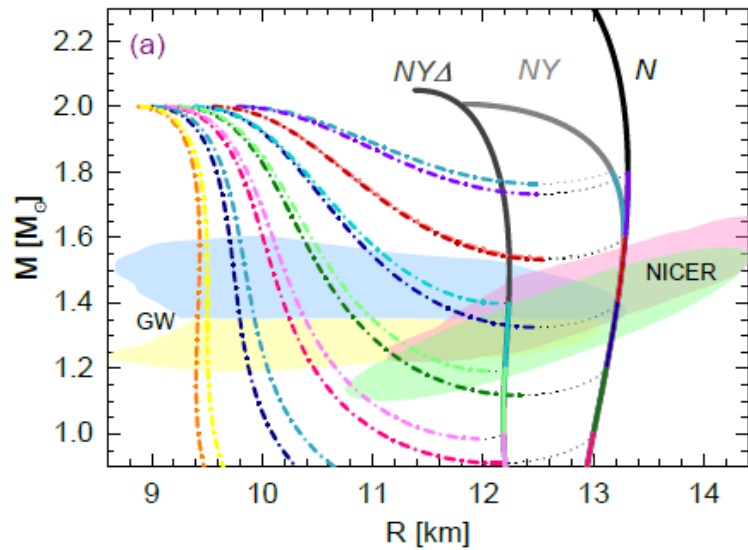


- At high density, matter could be deconfined

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Science case for the Einstein Telescope

...include dans
...nulle: peut-on
...une composition
...nucleonique?

Quarks in the core of neutron stars ?



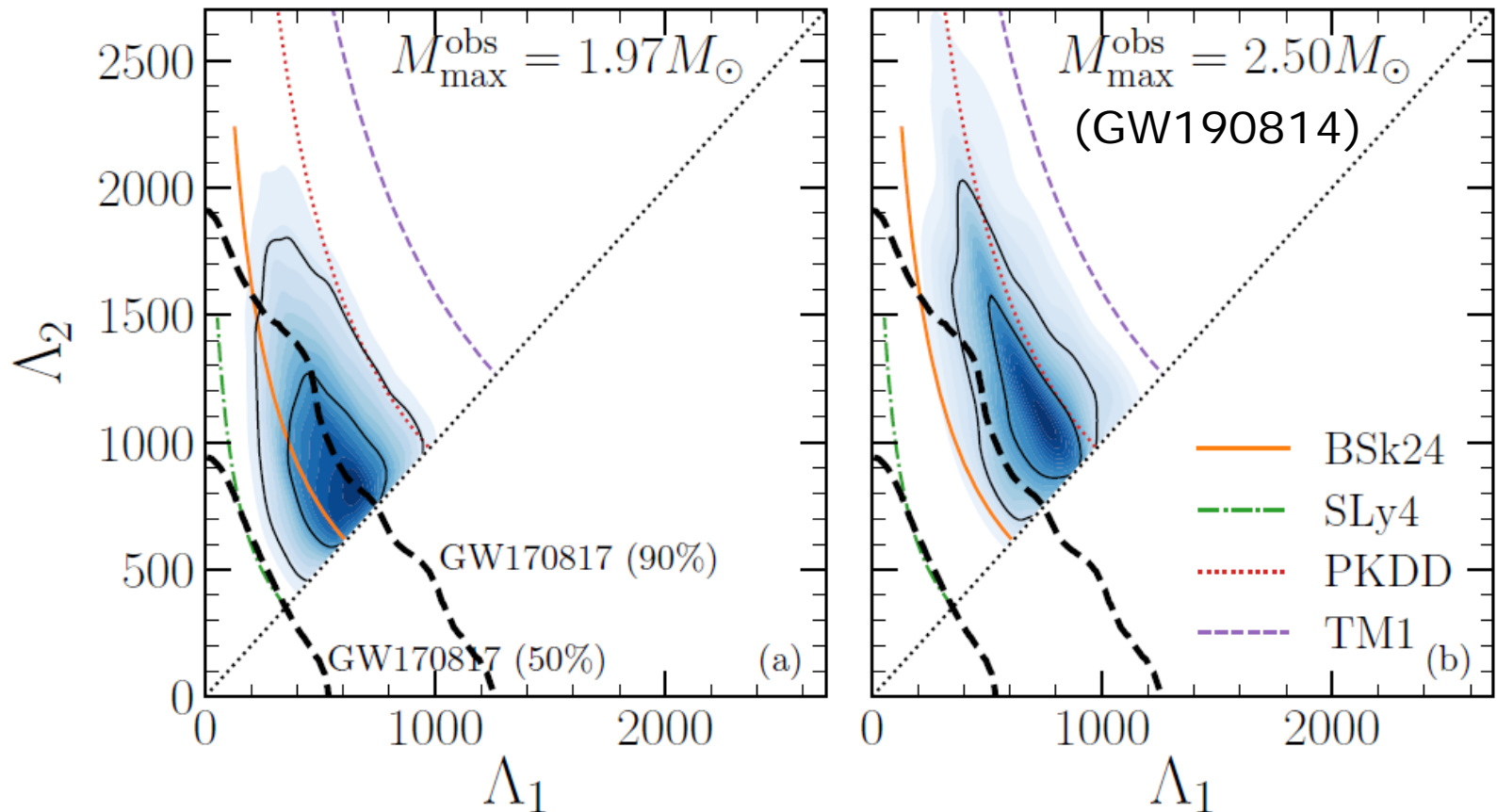
- At high density, matter could be deconfined
- Big uncertainty in the models
- Null hypothesis: can we exclude a nucleonic composition?

J.J.Li, A.Sedrakian, M.Alford,
PRD101 (2020) 063022

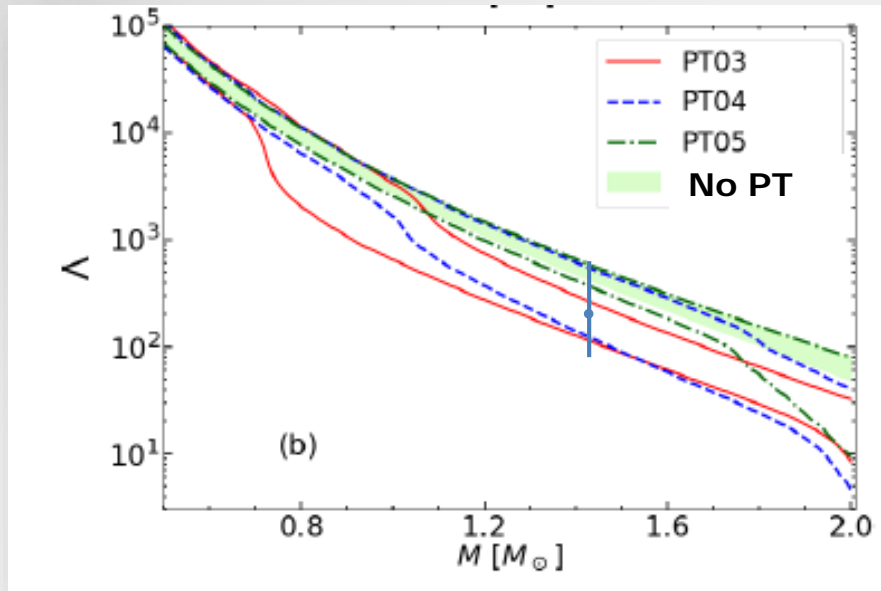
Quarks in the core of neutron stars ?

- A neutrons and protons composition is compatible with the observations
- But a great discovery potential

F.G., A.F.Fantina, NPN 2021

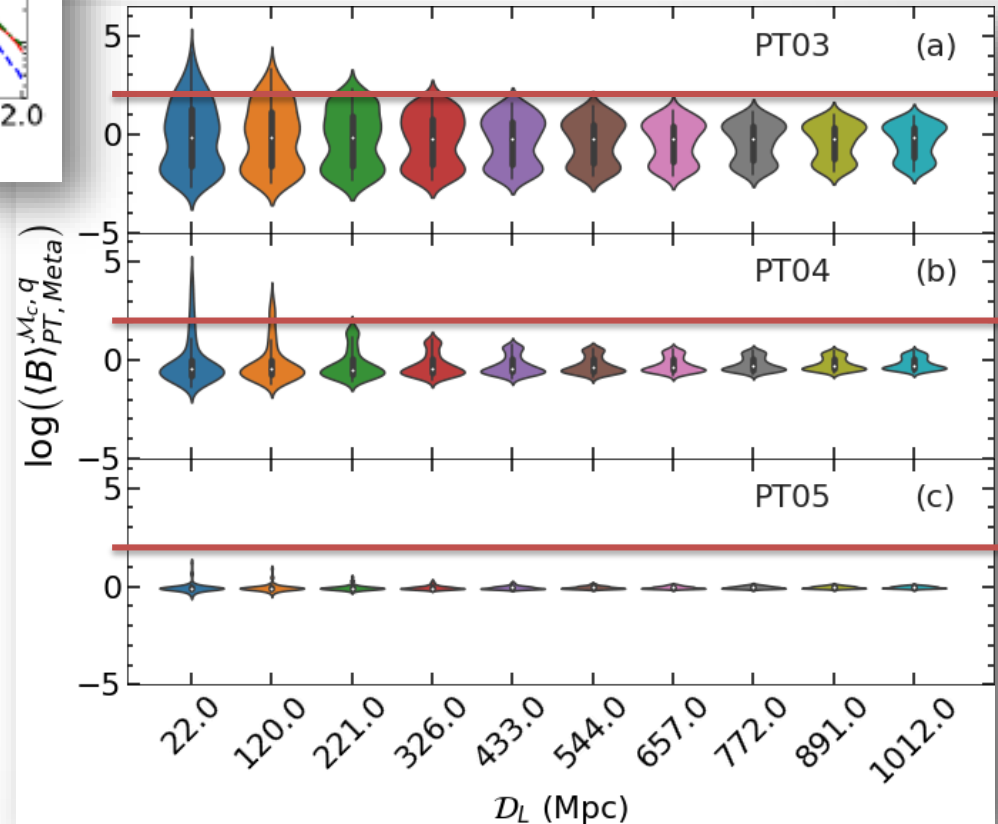


Quarks in the core of neutron stars ?

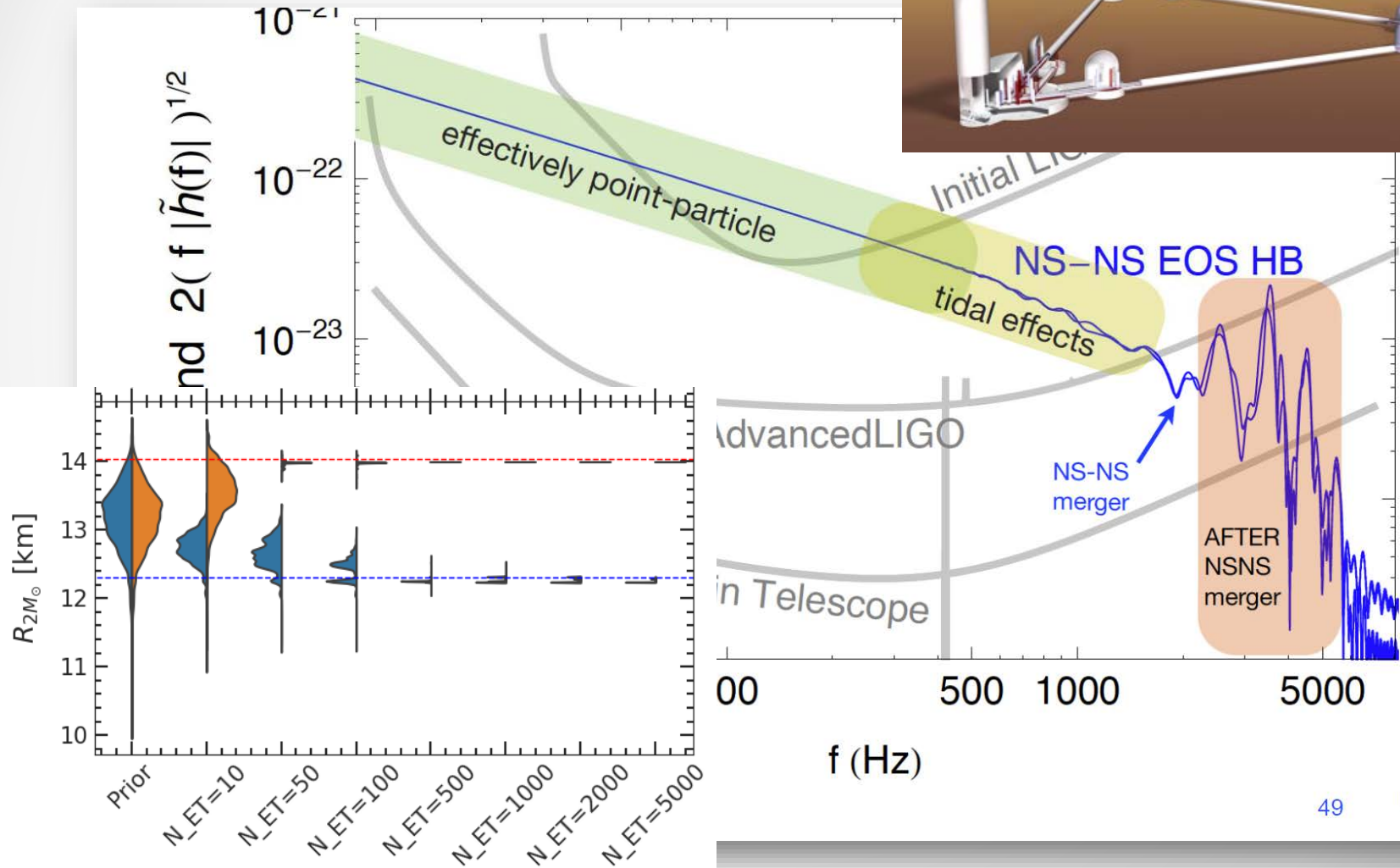
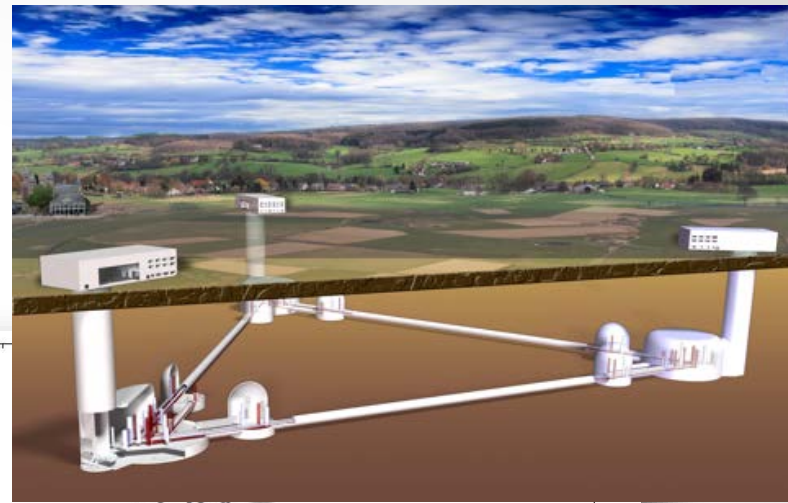


- Only a very close detection would allow identifying deconfined matter, and only if the quark core is large

- Need to reduce the uncertainties!
 - New nuclear observables \Rightarrow GANIL/Spiral2, FRIB
 - Progress in theory
 - New detections with better SNR \Rightarrow Einstein Telescope, Cosmic Explorer

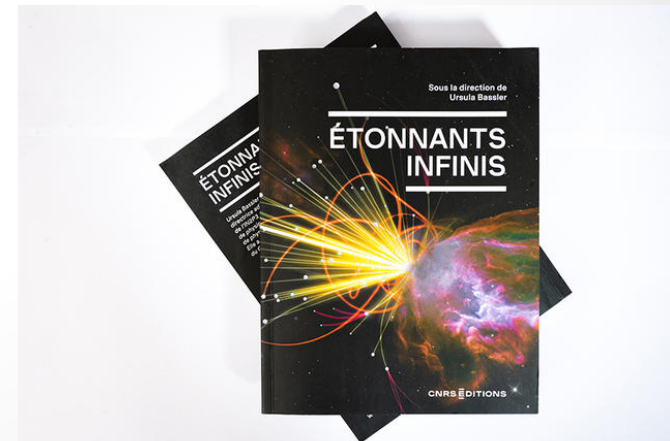


The Einstein Telescope project



Conclusions

- Gravitational waves: a new probe to explore ultra-dense matter
- Important synergies between nuclear physics and astrophysics since the first observation of a neutron star merger GW170817
- No evidence of deconfined matter (quark-gluon plasma) in the neutron star core, but exciting discoveries ahead





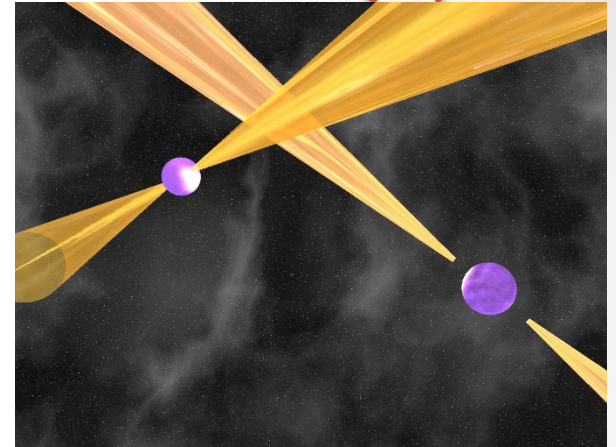
L'histoire de(s) découverte(s)

- **1934:** Prediction Baade&Zwicky
- **1967:** Découverte des pulsars
Bell&Hewish (prix Nobel 1974)



L'histoire de(s) découverte(s)

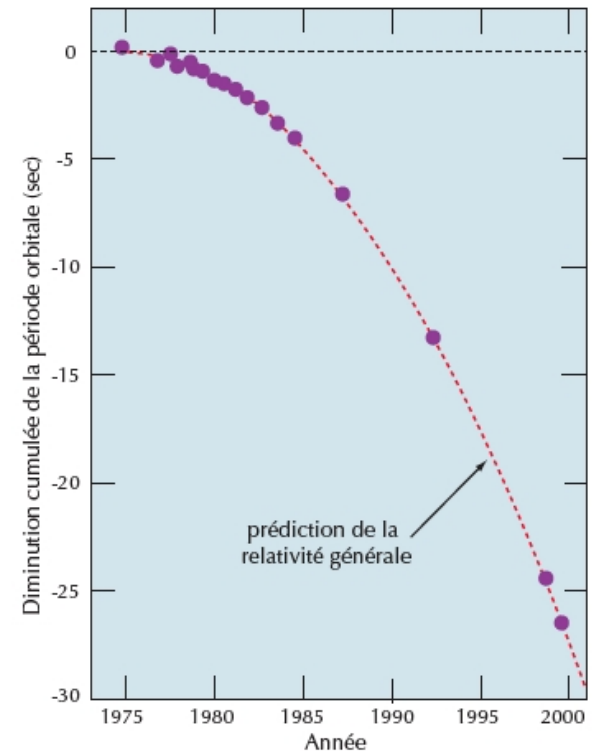
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- **1974:** Première détection d'une binaire
de pulsar (Hulse-Taylor)



Credit: M. Kramer, Univ. Manchester

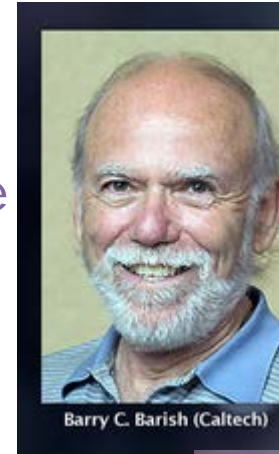
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- **1993:** d'ondes gravitationnelles (OG)
Hulse&Taylor (prix Nobel 1994)



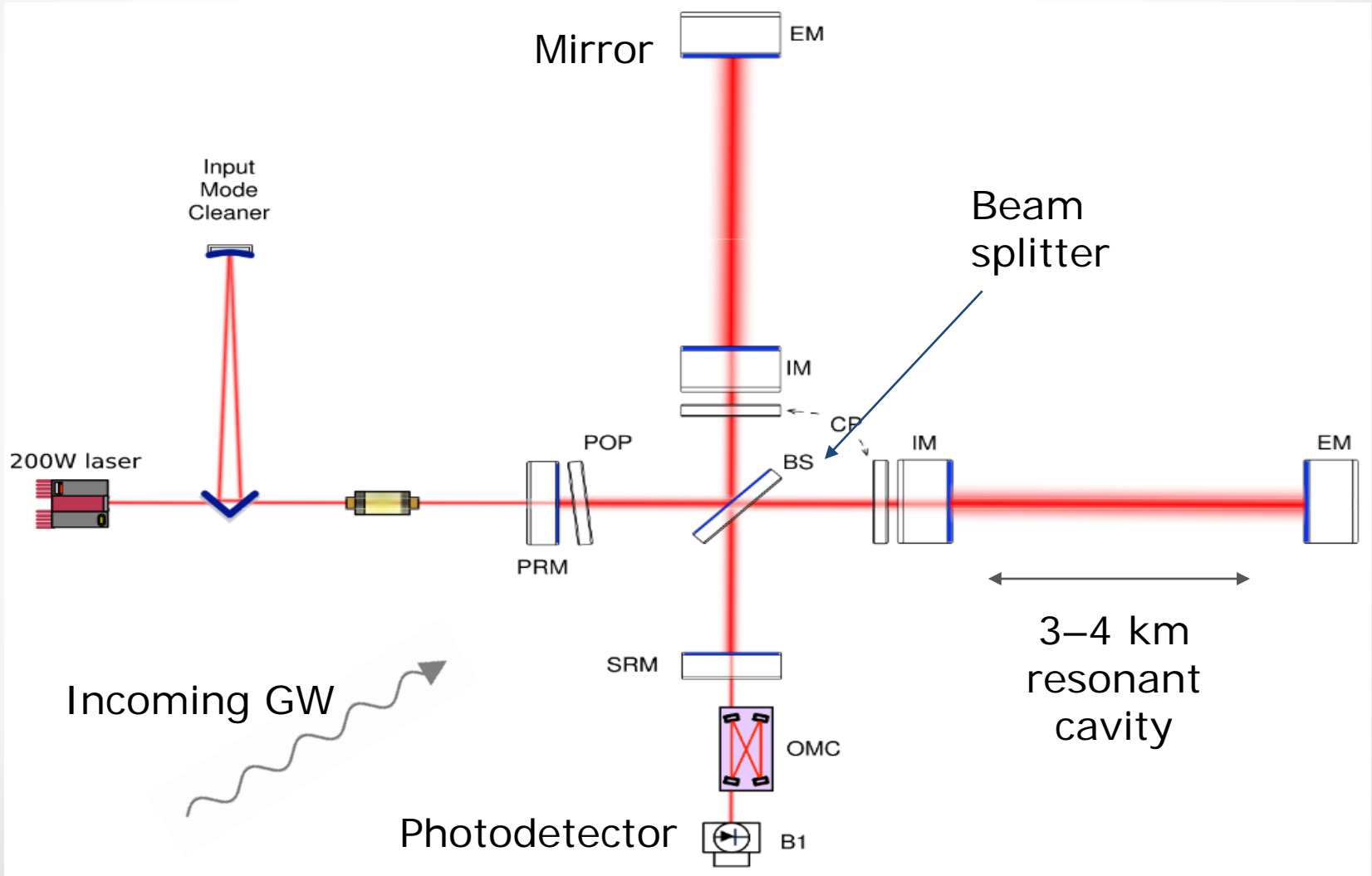
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- **1975-** Mise en évidence indirecte
1993: d'ondes gravitationnelles (OG)
Hulse&Taylor (prix Nobel 1994)
- **2015:** Première détection d'OG d'une binaire de trous noirs (LIGO -Virgo, prix Nobel 2017)
- **2017:** Première détection d'OG d'une binaire d'étoiles à neutrons
- **2023:** Début du quatrième run d'observation LVK

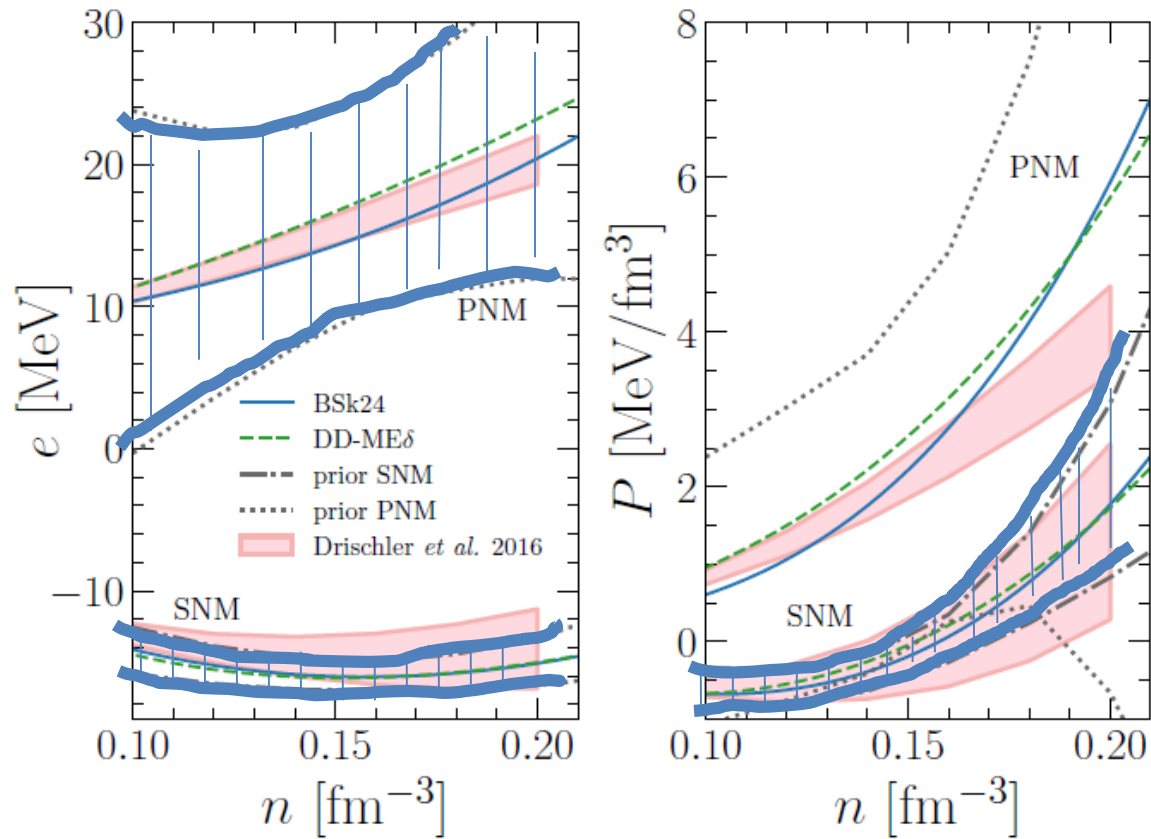


La détection d'OG

interféromètre Michelson-Fabry-Perot



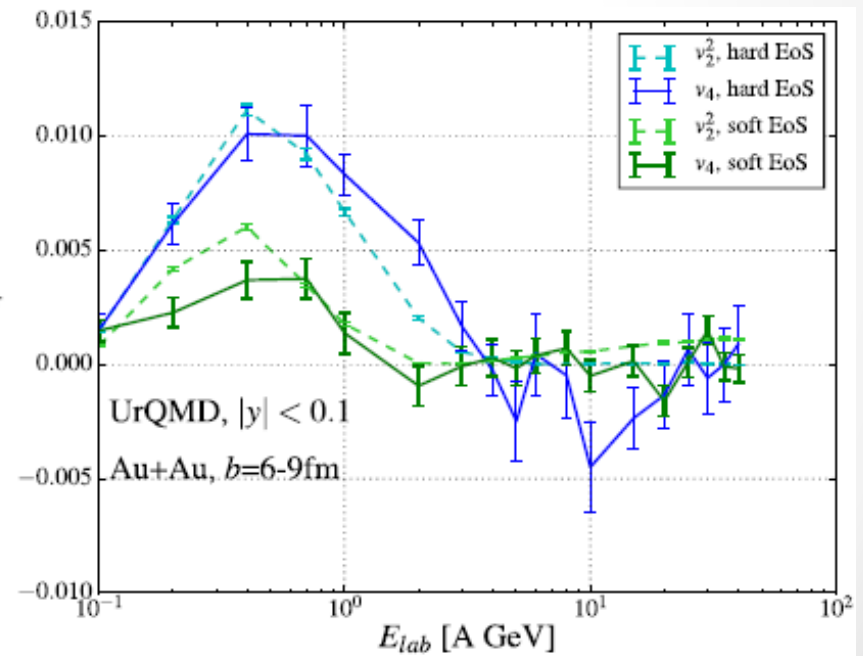
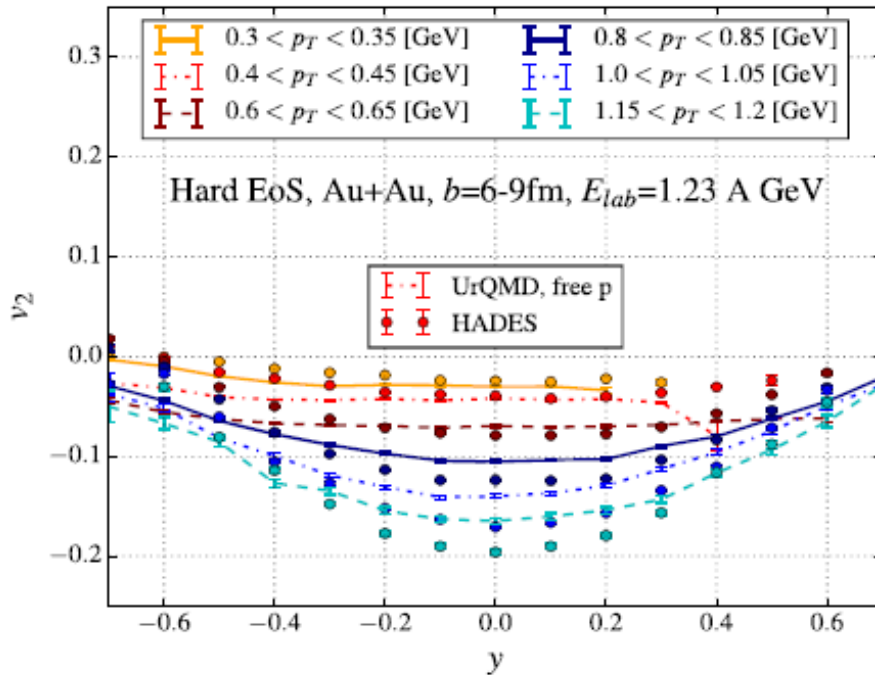
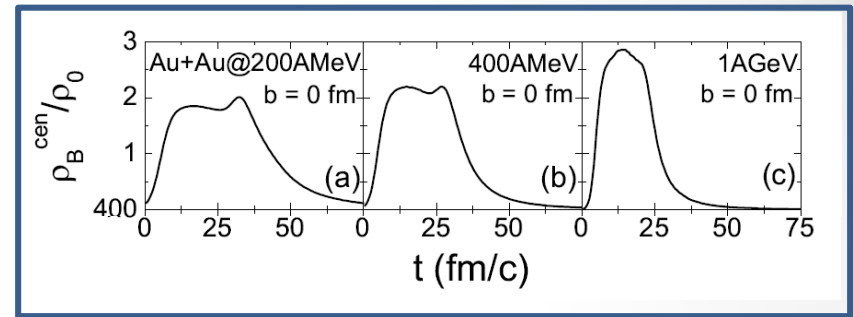
Contraintes expérimentales et théoriques



Tighter constraints from high energy experiments ?

J.Xu, PRC 2013

Elliptic flow @ HADES:
Transport model versus data

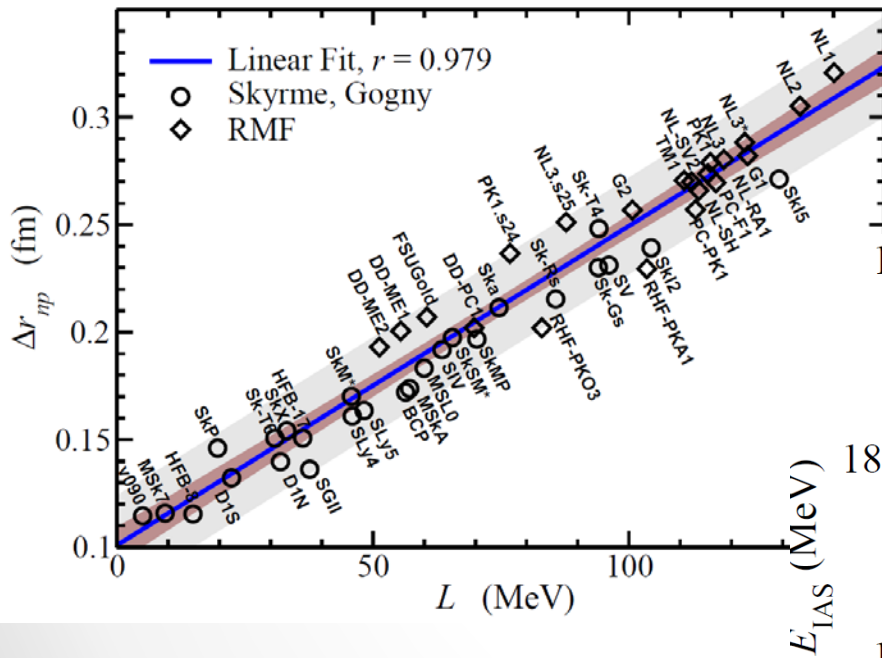


Strategy II: high precision

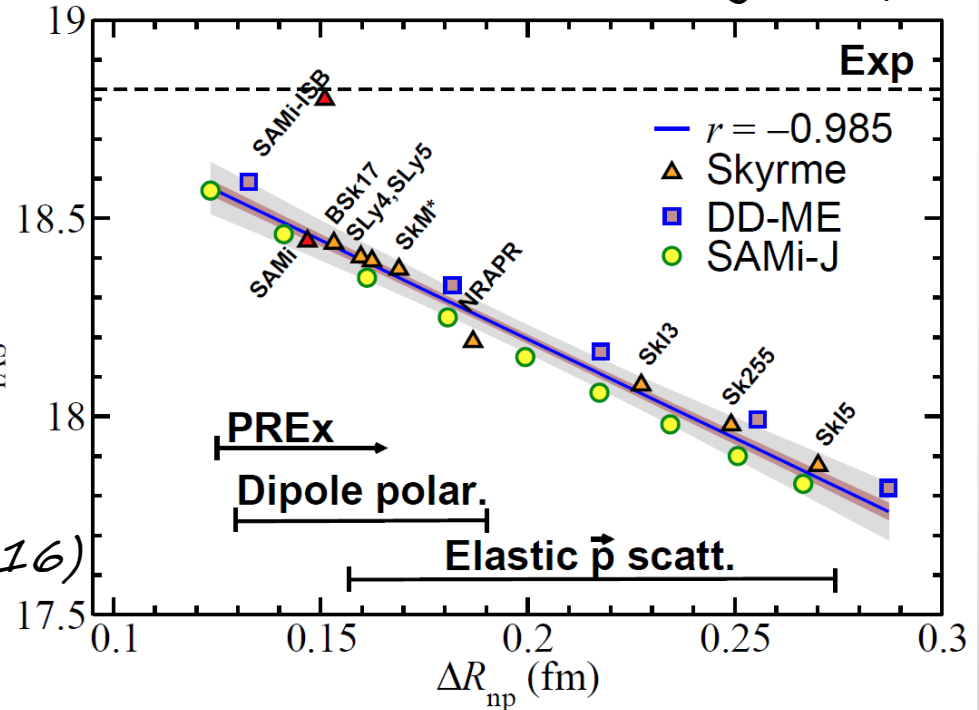
Isovector probes with stable nuclei:

- Skin
- IAS
- GDR response
- Isospin transport in HIC (INDRA/FAZIA)

Exp: $\Delta r_{np} = 0,1318-0,3072$



X.Roca-Maza, G.Colo, H.Sagawa (2018)



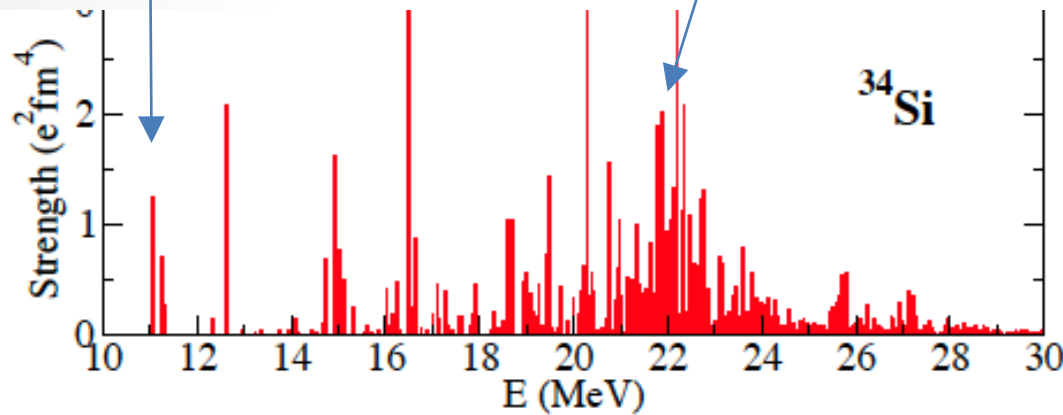
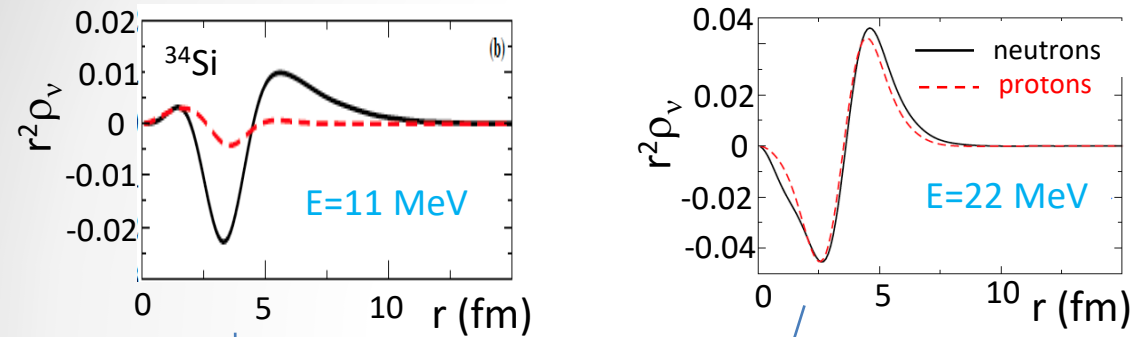
X.Vinas et al (2014)

P.G.Reinhard, W.Nazarewicz (2016)

J.Yang, J.Piekarewicz (2017)

Strategy III: new probes

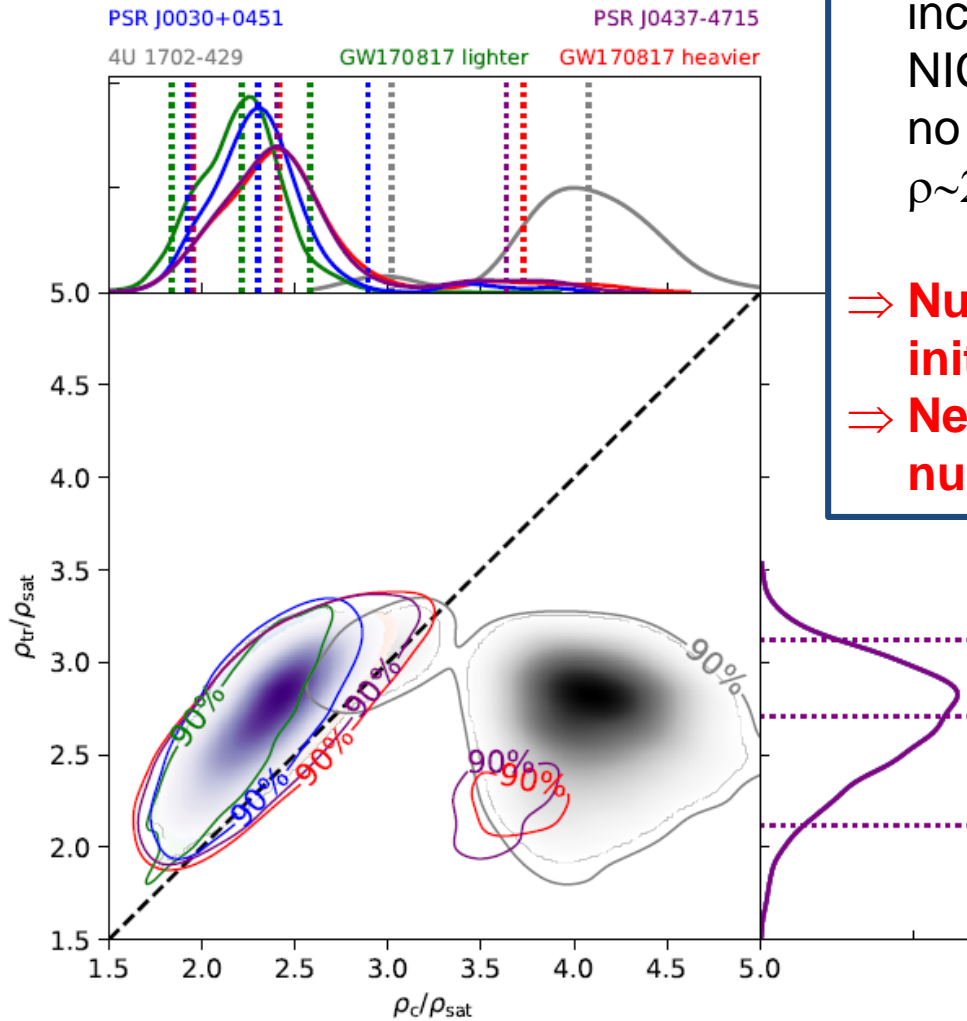
SSRPA calculations



Isoscalar probes in exotic nuclei:

- Soft monopole

D. Gambacurta, Phys. Rev. C 100, 014317 (2019)



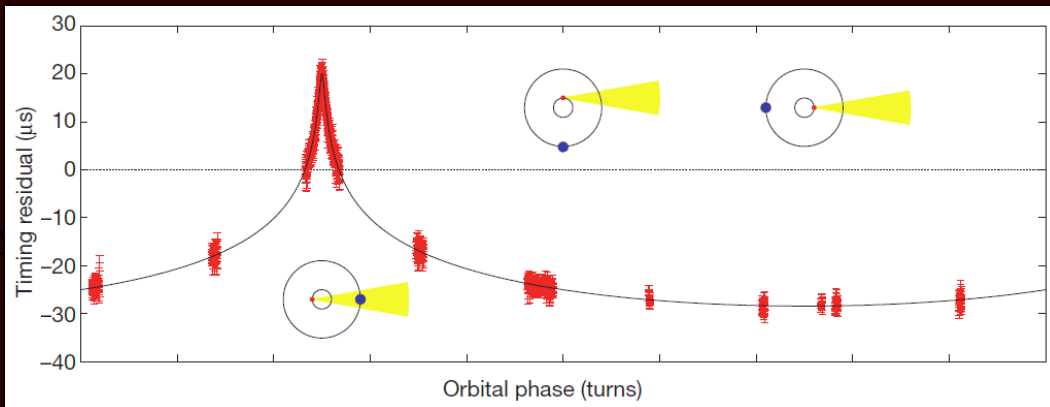
- « Reasonable » agnostic modeling of a 1st order phase transition including both LIGO/VIRGO and NICER data as constraints, predicts no transitions below densities $\rho \sim 2,5\rho_0$

⇒ **Nuclear physics is valid, but ab-initio modeling is not...**
 ⇒ **Need of more constraining nuclear data for $\rho_0 < \rho < 2\rho_0$**

Today: about 2000 Neutron Stars known in the Milky Way and Large Magellanic Cloud

Neutron Stars: today

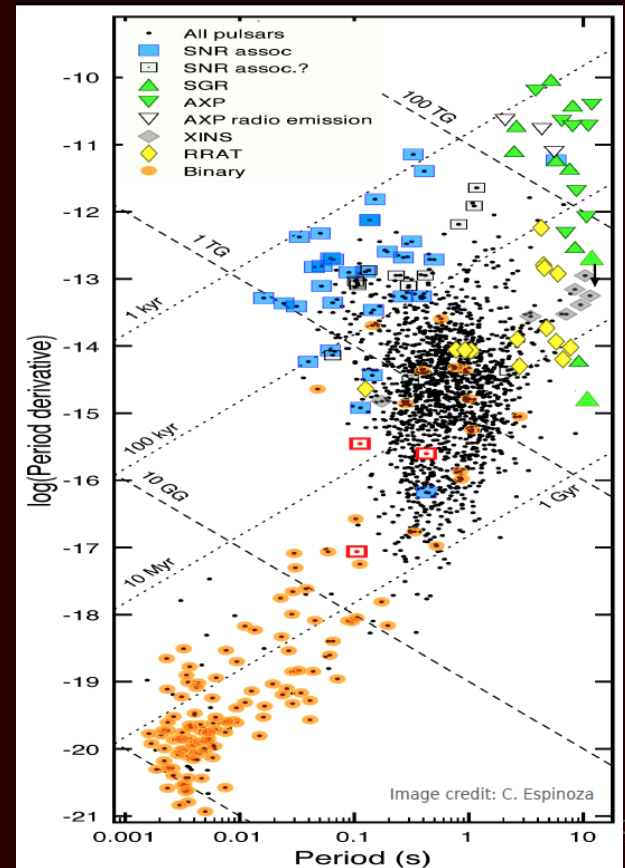
P. Demorest et al., Nature (2010)
 $M(\text{PSR J1614}) = 1.97 \pm 0.04$
J. Antoniadis et al., Science (2013).
 $M(\text{PSR J0348}) = 2.01 \pm 0.04$
H. Cromartie et al. Nature As. (2019)
 $M(\text{PSR J0740}) = 2.14 \pm 0.1$



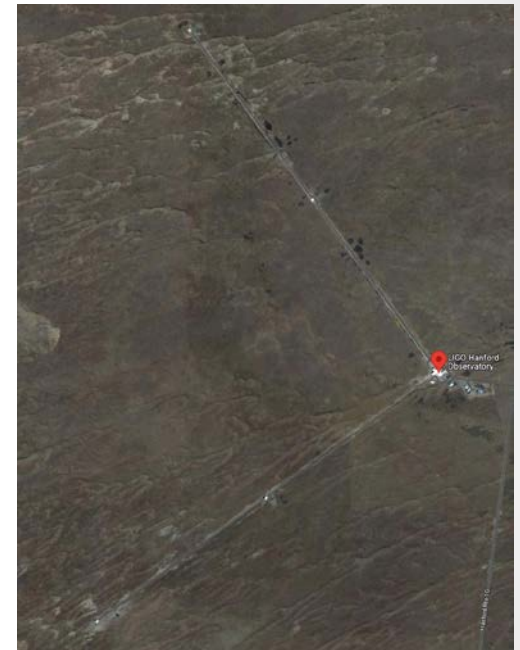
...supermassive objects: challenge for the strong interaction

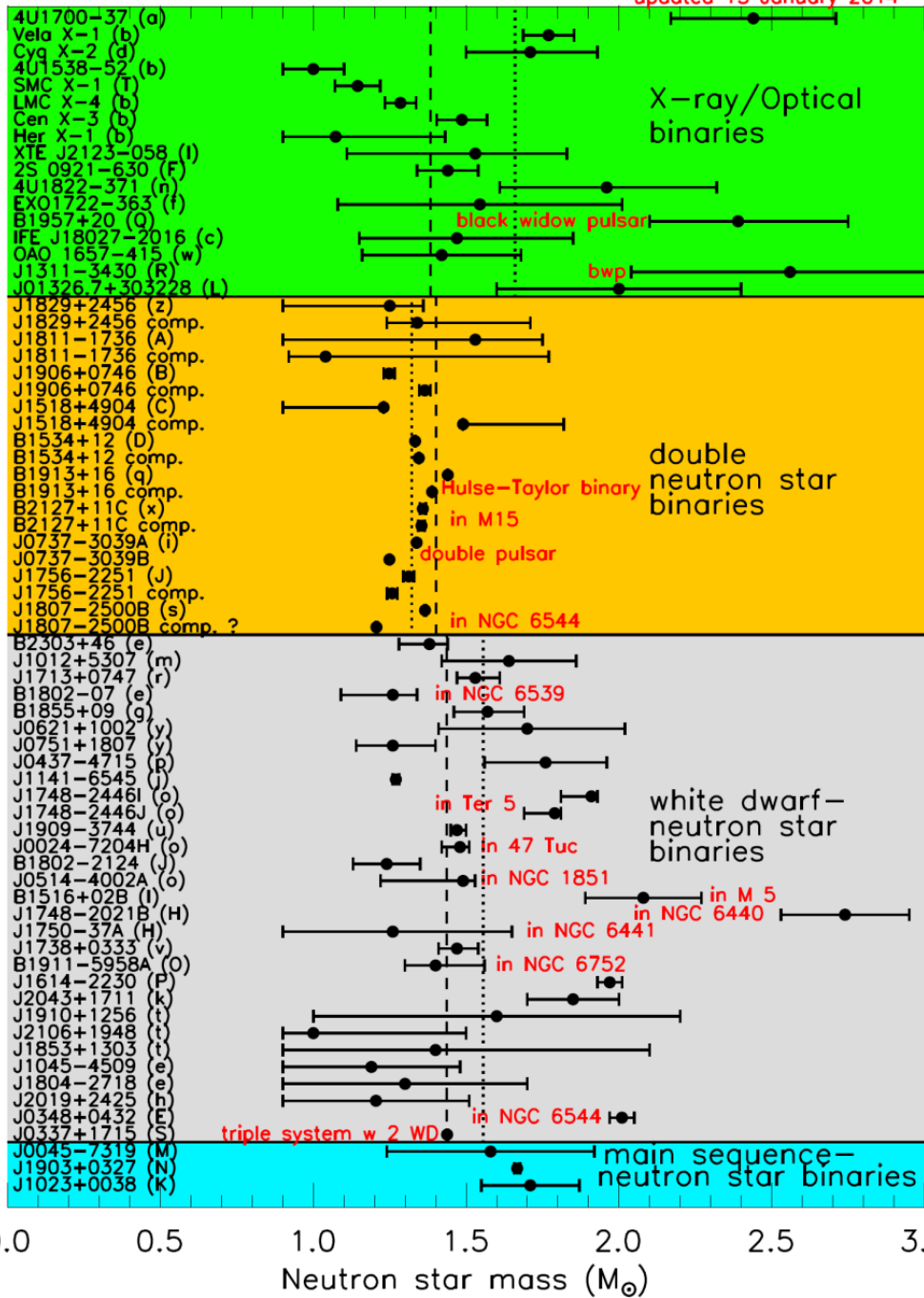
N. Rea et al., ApJ (2013).
 $B(\text{SGR 0418}) = 6 \times 10^{12} \text{ G}$

...SGR, pulsar, magnetars: unified picture

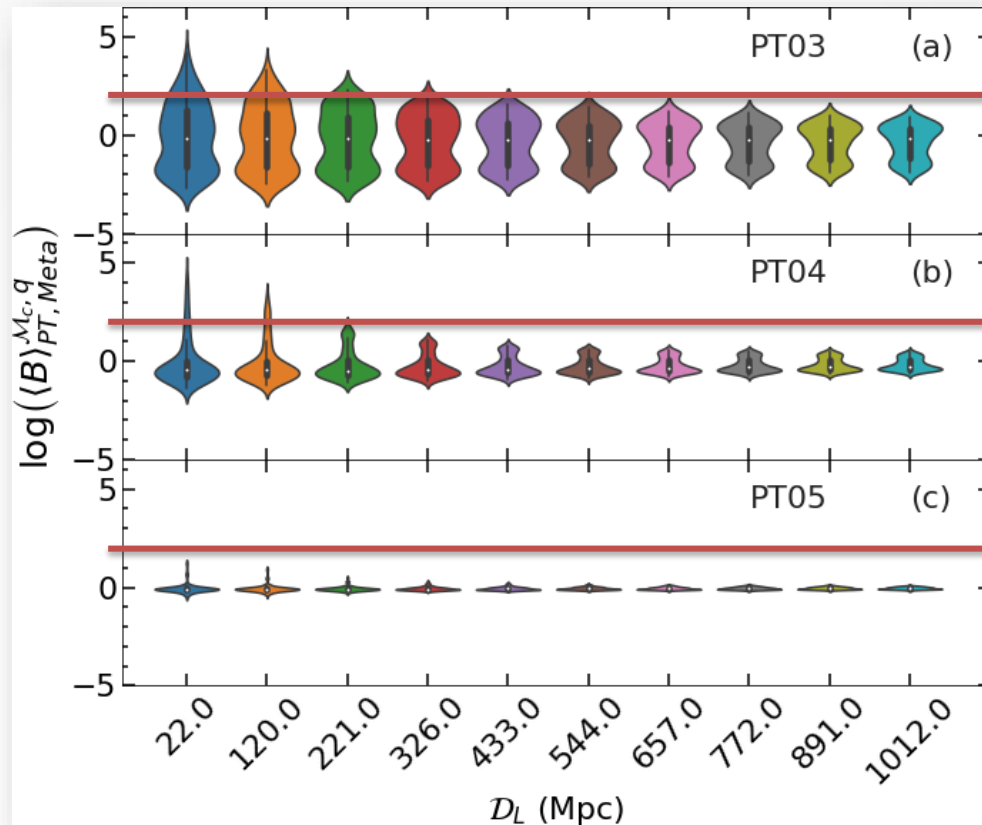


GW detectors: the present network





0.0 0.5 1.0 1.5 2.0 2.5 3.0
Neutron star mass (M_{\odot})

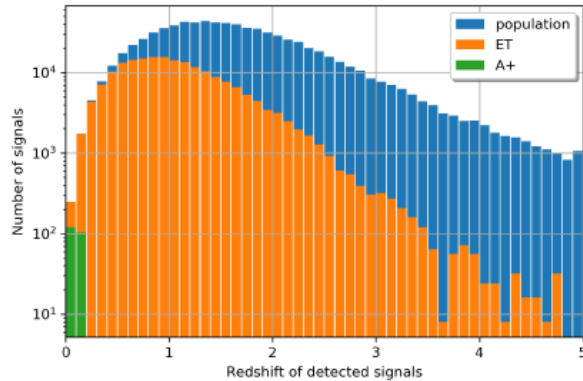


$$\log \left(\langle B \rangle_{PT, nucl}^{M_0, q_0} \right) = \int d\tilde{\Lambda} P_0^{GW}(\tilde{\Lambda}) \log \left[\frac{P_0^{nt}(\tilde{\Lambda})}{P_0^{nucl}(\tilde{\Lambda})} \right]$$

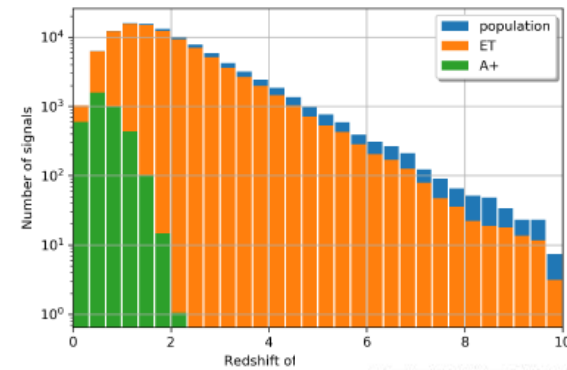
=> High detectability potential of a density jump with G3 detectors for an EARLY phase transition

ET sensitivity

BINARY NEUTRON-STAR MERGERS

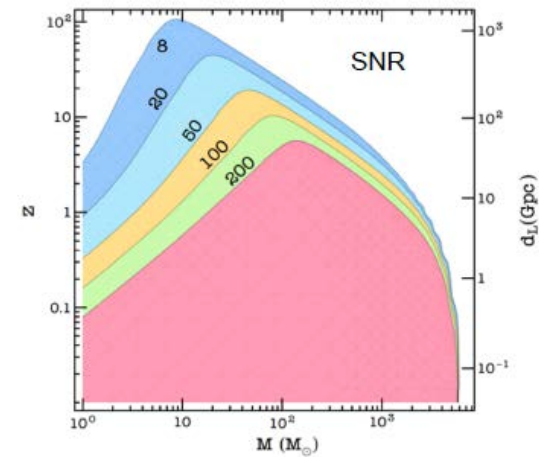
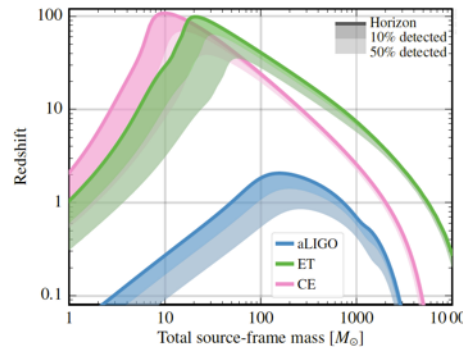


BINARY BLACK-HOLE MERGERS



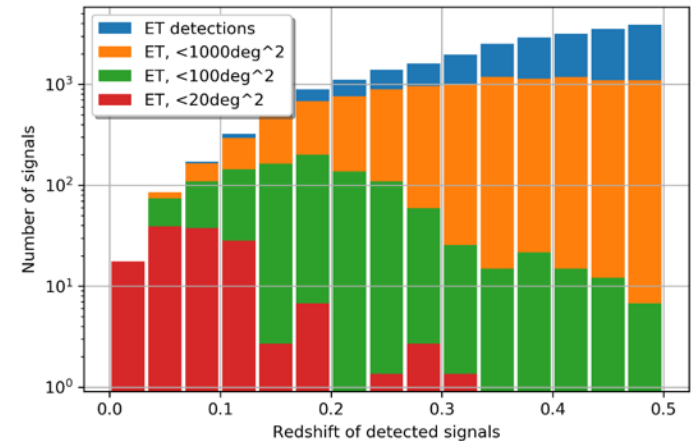
- 10^5 - 10^6 BBH detections per year
- 10^4 - 10^5 BNS detections per year among which ~ 10 - 100 with EM counterparts
- High SNR events
- Overlapping events

~ 1 detection every 30s

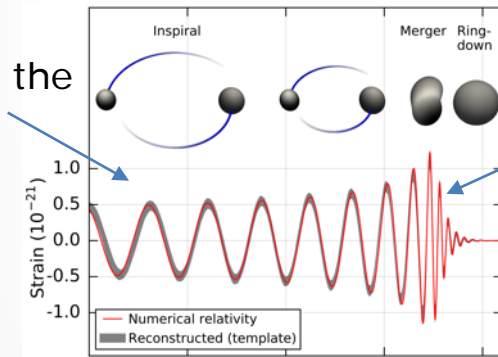


ET sensitivity

- BNS detection with EM counterparts and localization precision $< 20 \text{ deg}^2$: $\sim 10\text{-}100$ per year
- Overlap with many BBH signals
- Potentially, very long signals
- ET will be able to provide alerts few hours before the merger



Identify early the inspiral ...



... and provide alert before the merger phase

- And with ~ 500 BNS-EM detection, we can reach Planck resolution on H_0 measurement

