

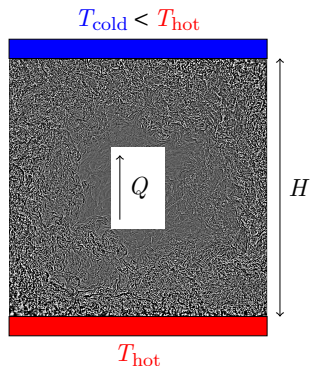


Turbulent transitions in thermal convection

Équipe Convection – Laboratoire de physique ENSL

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Model System : the Rayleigh-Bénard cell



Control parameters

- ▶ Thermal forcing *vs* diffusion :

$$Ra = \frac{g\alpha(T_{\text{hot}} - T_{\text{cold}})H^3}{\nu\kappa}$$

- ▶ Viscous *vs* Thermal diffusion :

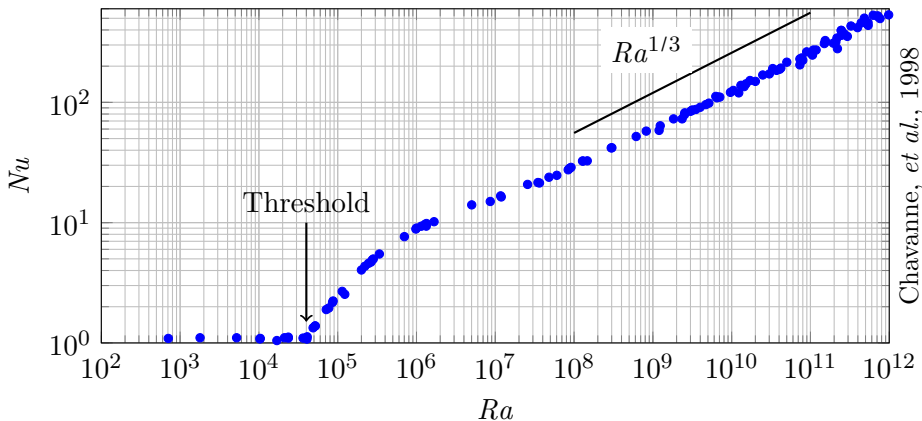
$$Pr = \frac{\nu}{\kappa}$$

System Response

- ▶ Normalized thermal flux :

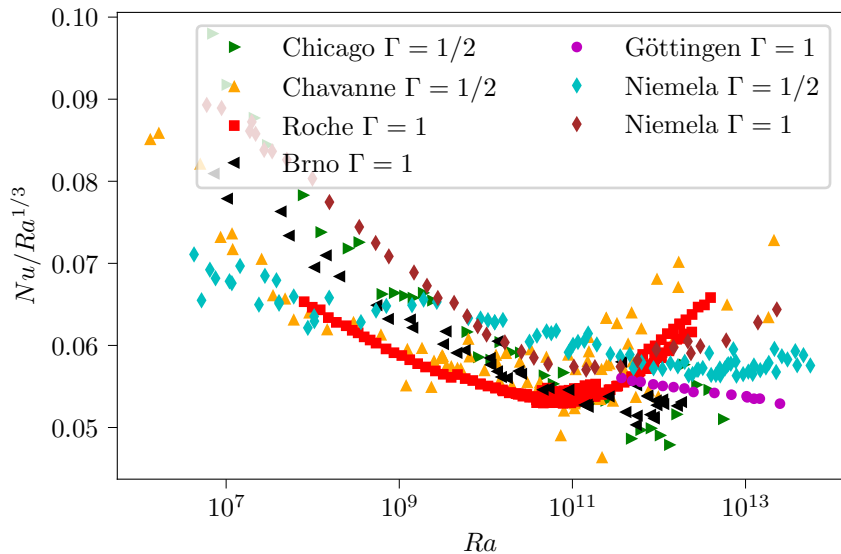
$$Nu = \frac{QH}{\lambda(T_{\text{hot}} - T_{\text{cold}})}$$

Scaling laws



$$\underbrace{\left(\frac{QH}{\lambda(T_{\text{hot}} - T_{\text{cold}})} \right)}_{Nu} = \underbrace{\left(\frac{g\alpha(T_{\text{hot}} - T_{\text{cold}})H^3}{\nu\kappa} \right)^{1/3}}_{Ra^{1/3}}$$

$Ra > 10^{12}$?



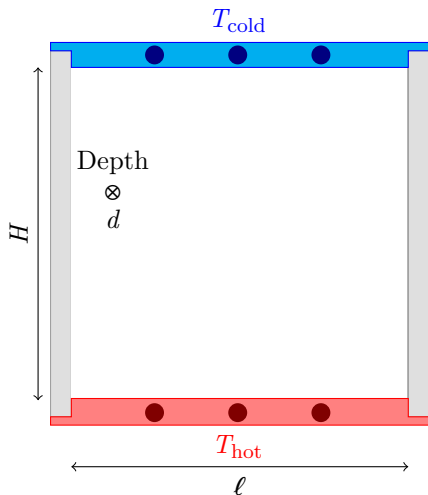
Open questions

- ▶ What is this transition ?
Kraichnan asymptotic regime : $Nu \propto Ra^{1/2}$?
Transition to turbulence in the boundary layer ?
- ▶ What triggers or inhibits this transition ?
- ▶ How to explain to quantitative difference between experiments ?

Experimental apparatus

- ▶ Large infrastructure and/or special working fluid needed to reach $Ra > 10^{12}$
- ▶ Alternative approach : use plate roughness to trigger the transition to turbulence in the boundary layer

Experimental setup : rectangular cell



► **Room temperature**

$$H = \ell = 41.5 \text{ cm}, d = 10 \text{ cm}$$

Water :

$$3 < Pr < 7; 10^9 < Ra < 10^{11}$$

Fluorocarbon :

$$10 < Pr < 15; 10^{10} < Ra < 10^{12}$$

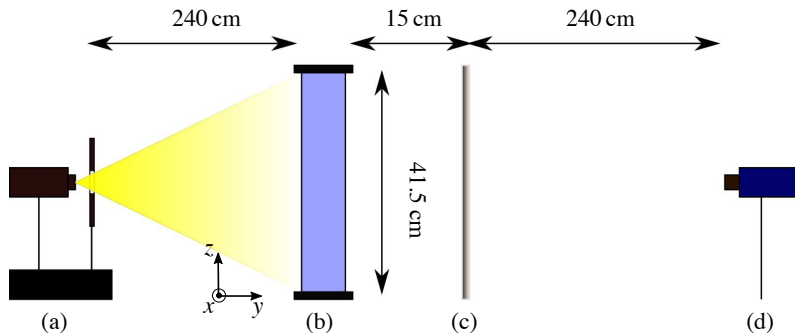
► **Liquid helium**

$$H = 4.8 \text{ cm}, \ell = 5 \text{ cm}, d = 1.5 \text{ cm}$$

$$0.5 < Pr < 0.71;$$

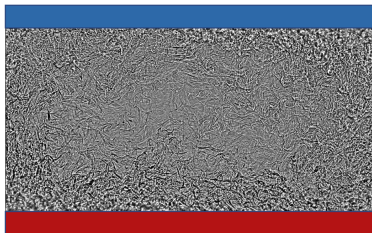
$$3 \times 10^9 < Ra < 3 \times 10^{10}$$

Non-invasive technique : shadowgraphy

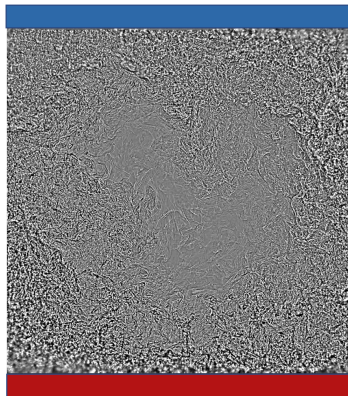


$$\frac{I(x, z)}{I_0} \approx \int_0^d \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) (\ln n) dy$$

Shadowgraph in Fluorocarbon

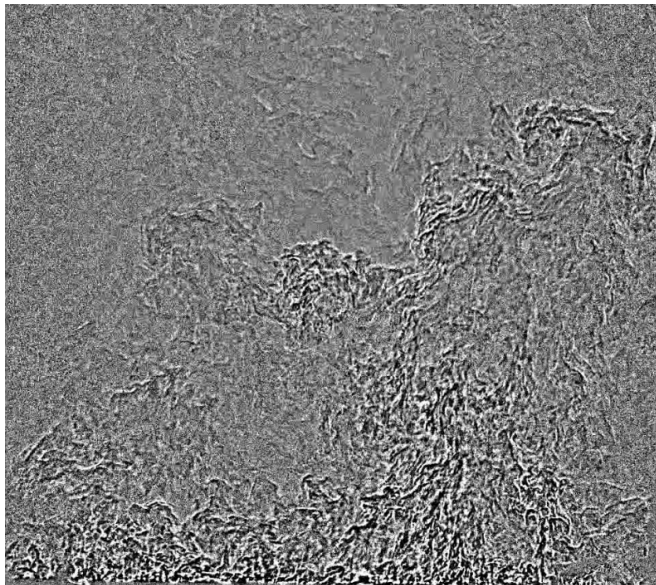


- $\Gamma = 2$
- $Ra = 9 \cdot 10^{10}$

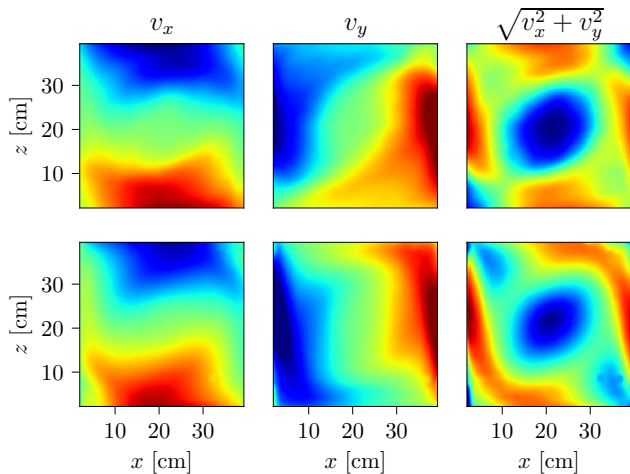


- $\Gamma = 1$
- $Ra = 2 \cdot 10^{12}$

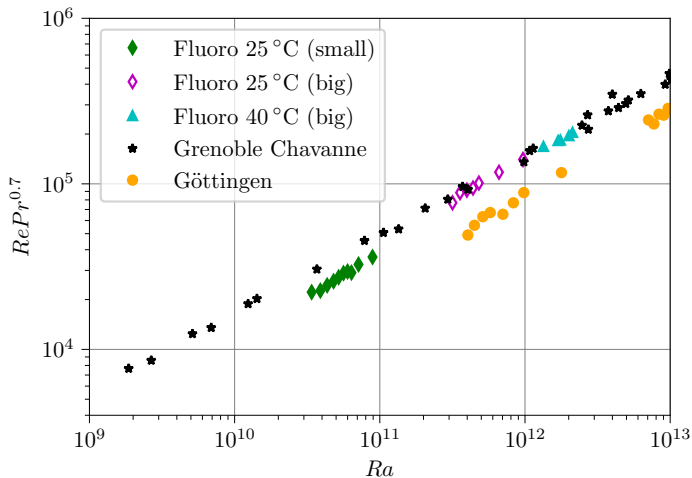
Shadowgraph in liquid helium



CIV on shadowgraph (top) vs PIV (bottom)



Reynolds number scaling



Friction coefficient

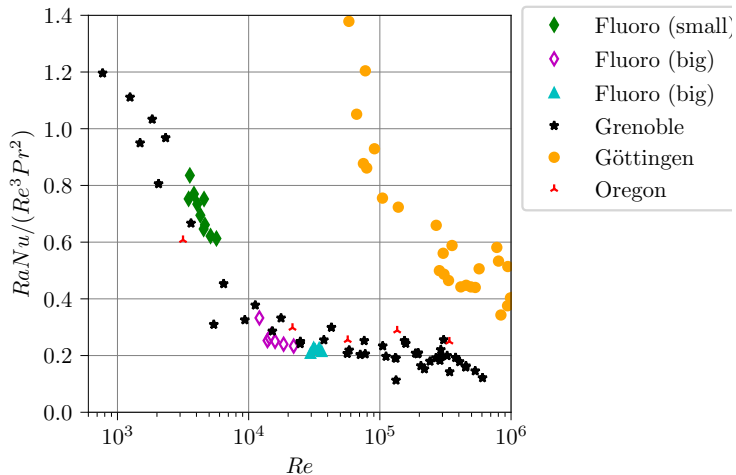
- ▶ Balance of kinetic energy

$$\epsilon = \frac{\nu^3}{h^4} (Nu - 1) Ra Pr^{-2}$$

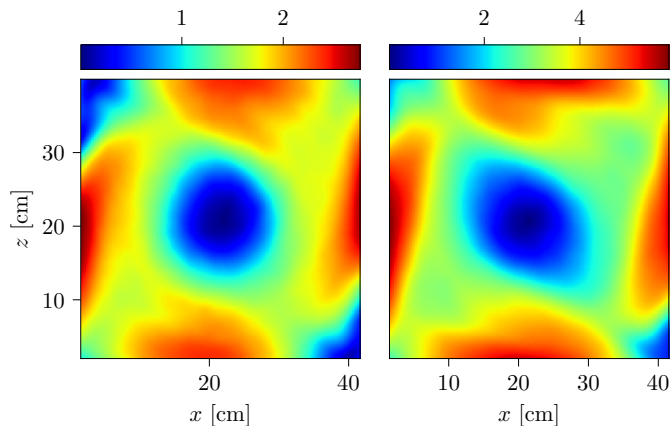
- ▶ Coefficient de friction

$$\frac{(Nu - 1) Ra Pr^{-2}}{Re^3} = \frac{\nu (\mathbf{grad} \mathbf{u})^2}{u^3 / h}$$

Friction coefficient

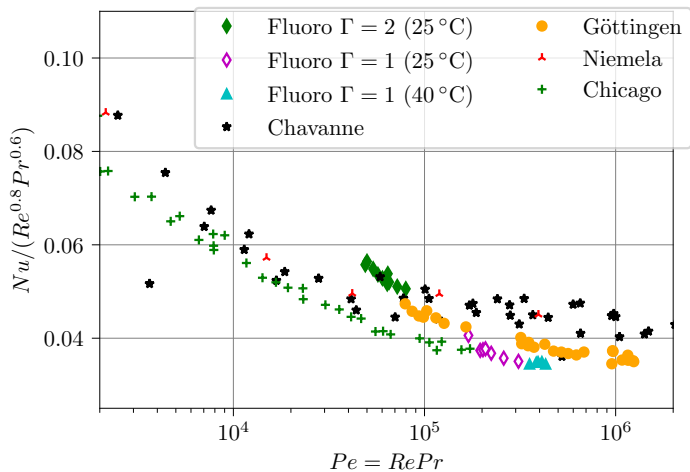


Mean field at the transition



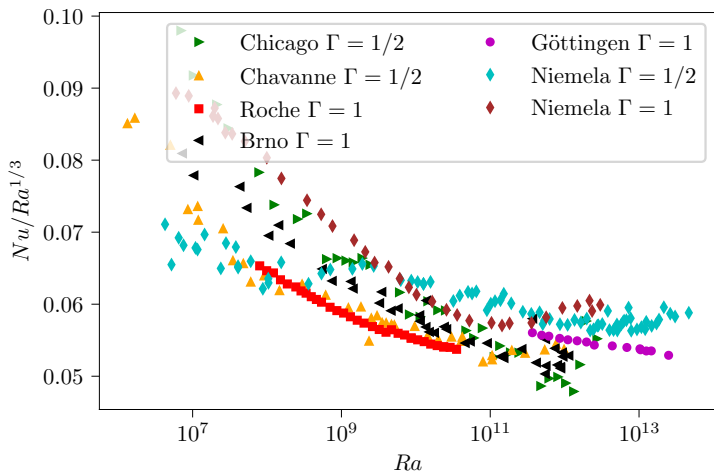
Velocity fields in the $\Gamma = 1$ cell with FC770. Left : $Ra = 3.6 \times 10^{11}$.
Right : $Ra = 2.1 \times 10^{12}$.

Péclet number as the control parameter



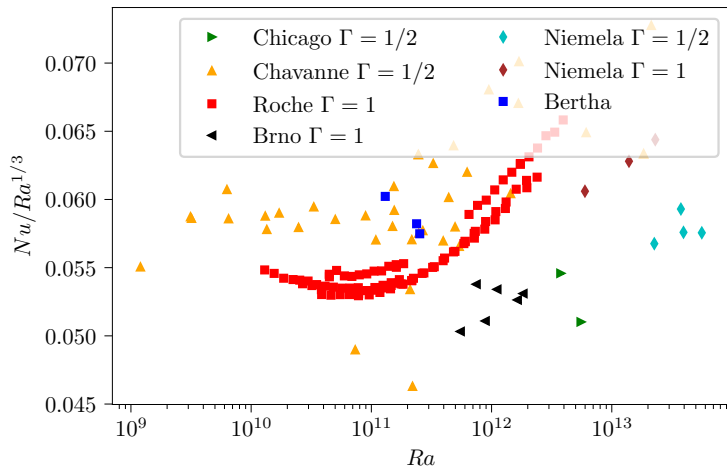
$Pr \leq 1$ in Boussinesq conditions

$$\Delta\lambda/\lambda < 0.1, \Delta\mu/\mu < 0.1, \Delta\alpha/\alpha < 0.1, \Delta c_p/c_p < 0.1, \Delta\kappa/\kappa < 0.1, \\ \Delta\rho/\rho < 0.1, \alpha\Delta T < 0.1$$



$Pr \geq 1$ in Boussinesq conditions

$$\Delta\lambda/\lambda < 0.1, \Delta\mu/\mu < 0.1, \Delta\alpha/\alpha < 0.1, \Delta c_p/c_p < 0.1, \Delta\kappa/\kappa < 0.1, \\ \Delta\rho/\rho < 0.1, \alpha\Delta T < 0.1$$



Conclusions

- ▶ Threshold in terms of Reynolds or Péclet number : role of geometry and Prandtl number ;
- ▶ Transition in the $Nu - Ra$ plane not very sharp ;
- ▶ Transition of friction from $1/\sqrt{Re}$ to plateau ;
- ▶ No dramatic change of mean velocity field.