

Dynamo action sustaining turbulence: a subcritical transition

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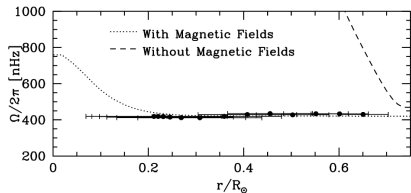
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Astrophysics :



Numerical simulation for the radiative zone of the sun, with and without magnetic field, compared to observations (*Eggenberger et al. 2005*)



Schematic view of the convective (outer) and radiative (inner) zones of the sun (Wikipedia)

Physics : Subcritical transition to turbulence scenario

The Tayler-Spruit dynamo

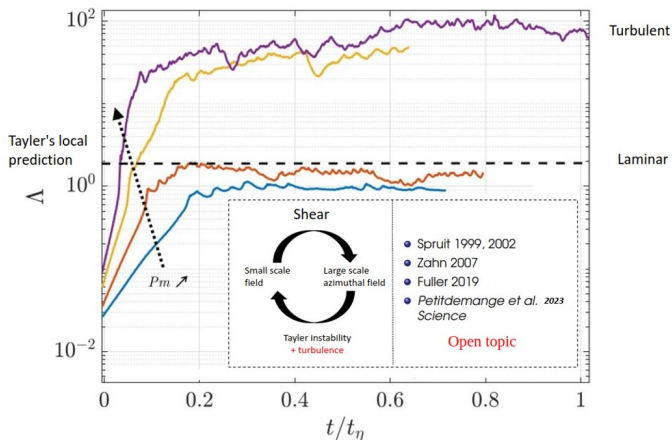
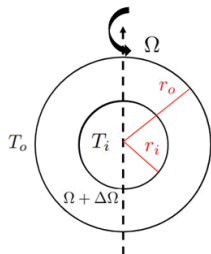


Figure 1 – Time evolution of the dimensionless magnetic energy

$\Delta = \int B^2 / \rho \mu \Omega \eta dV$ for different values of fluid conductivity (Petitdemange et al. 2023)



Solve for $\mathbf{u}, \mathbf{B}, \Theta$

$$\chi = \frac{r_i}{r_o} = 0.35$$

Boussinesq

$$\Delta T = T_o - T_i$$

$$\nu, \kappa, \eta$$

Figure 2 – Geometry

- BC for u : no slip
- BC for B : insulating at $r = r_o$, conducting at $r = r_i$
- Initial B : arbitrary weak
- Imposed (and maintained) inner rotation $\Delta\Omega$

MHD equations for a stably stratified incompressible fluid in rotation
under Boussinesq approximation (PaRoDy code)

$$\begin{aligned}
 \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} &= -\frac{\nabla P}{\rho_0} - 2\boldsymbol{\Omega} \times \mathbf{u} + \nu \Delta \mathbf{u} + \frac{1}{\rho_0 \mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} \\
 &\quad + \alpha \Theta g_0 r \mathbf{e}_r \\
 \nabla \cdot \mathbf{u} &= 0 \\
 \frac{\partial \Theta}{\partial t} + (\mathbf{u} \cdot \nabla) \Theta &= \kappa \Delta \Theta - (\mathbf{u} \cdot \nabla) T_s \\
 \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{u} \times \mathbf{B}) + \eta \Delta \mathbf{B} \\
 \nabla \cdot \mathbf{B} &= 0
 \end{aligned} \tag{1}$$

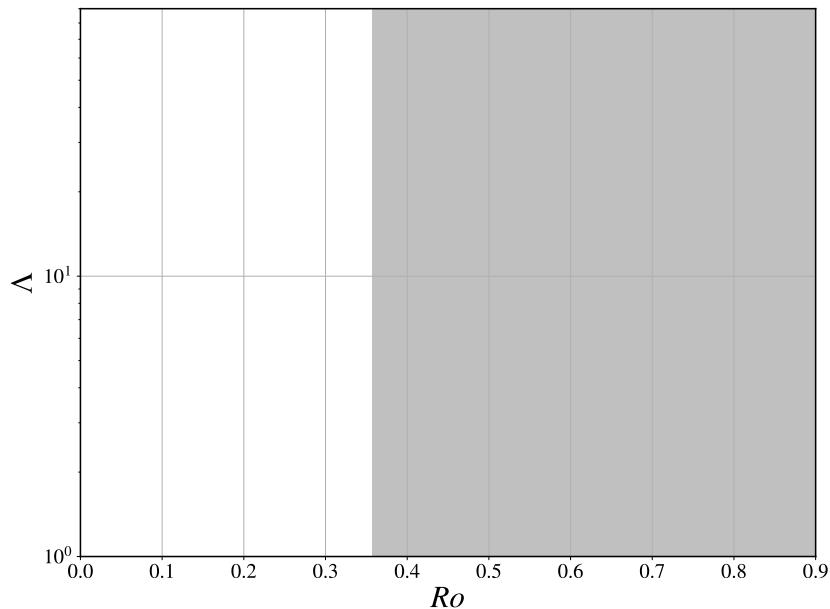
$$Ro = \frac{\Delta \Omega}{\Omega}, \quad Ek = \frac{\nu}{r_o^2 \Omega}, \quad \frac{N}{\Omega} = \sqrt{\frac{\alpha g \Delta T}{(r_o - r_i) \Omega^2}}, \quad Pr = \frac{\nu}{\kappa}, \quad P_m = \frac{\nu}{\eta} \tag{2}$$

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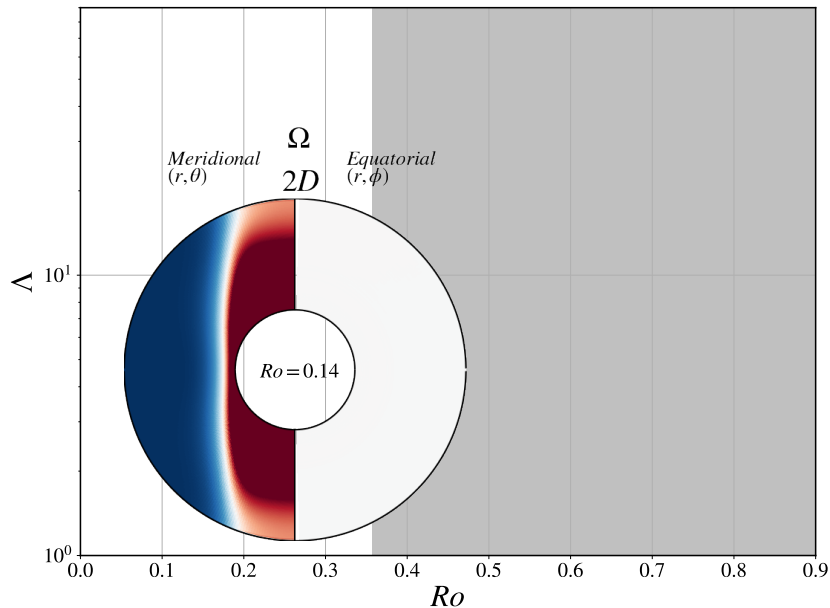
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- Fixed $Ek = 10^{-5}$, $Pr = 0.1$ and $\frac{N}{\Omega} = 1.24$
- Varying Ro and Pm
- Look at $\Lambda = \int \frac{B^2}{\rho\mu\eta\Omega} dV$

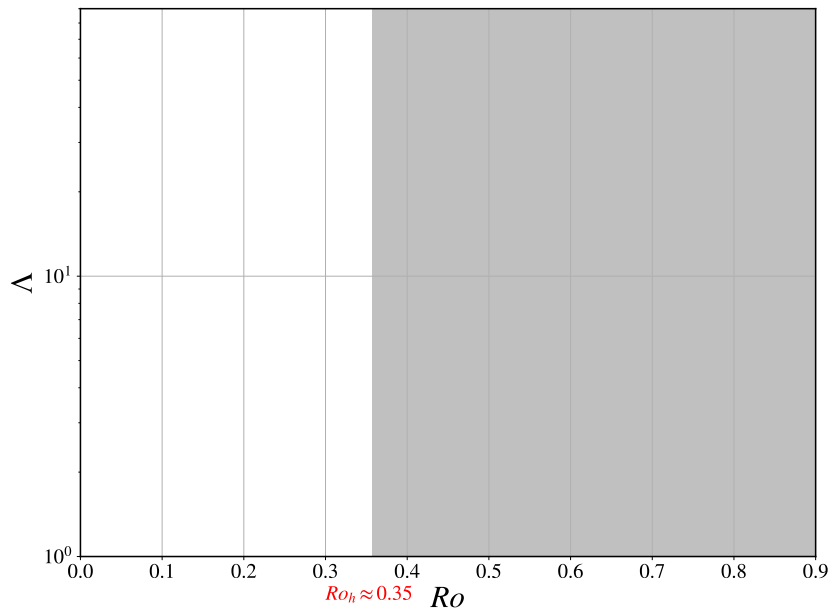
Bifurcations



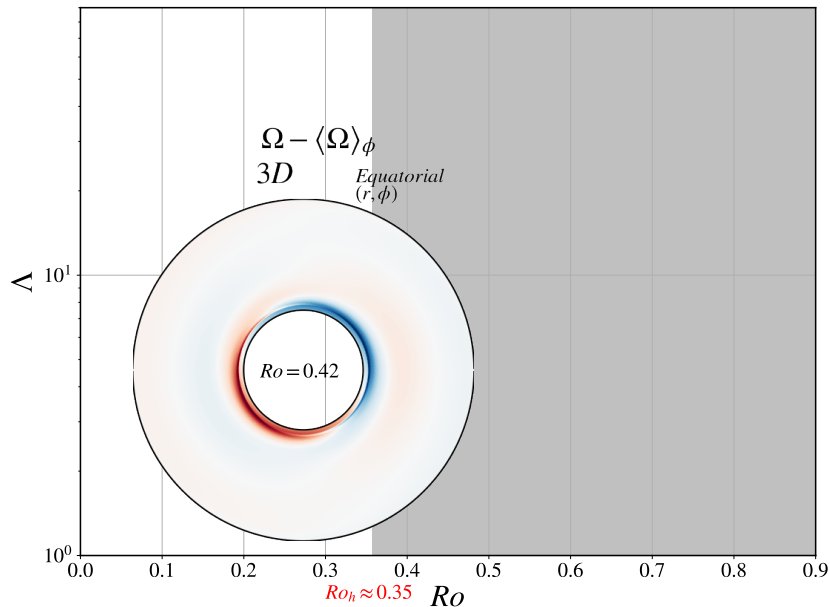
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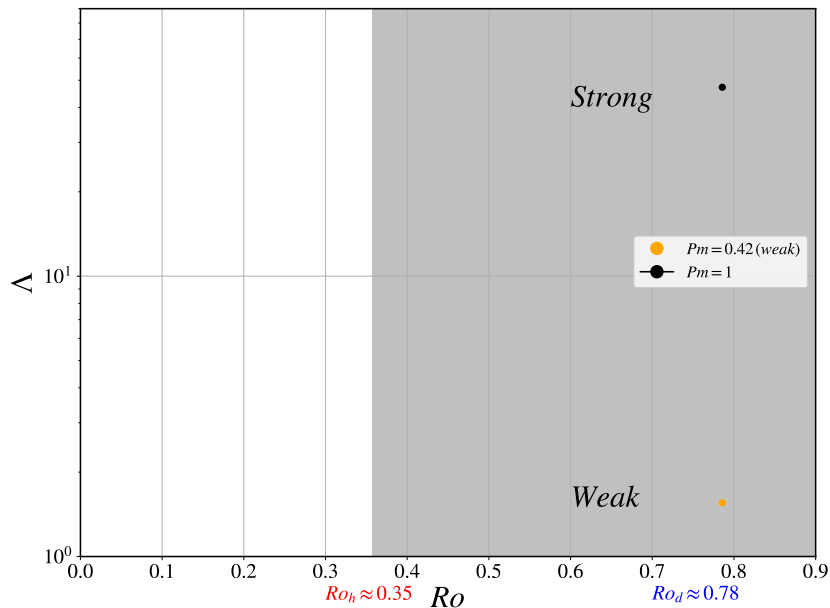
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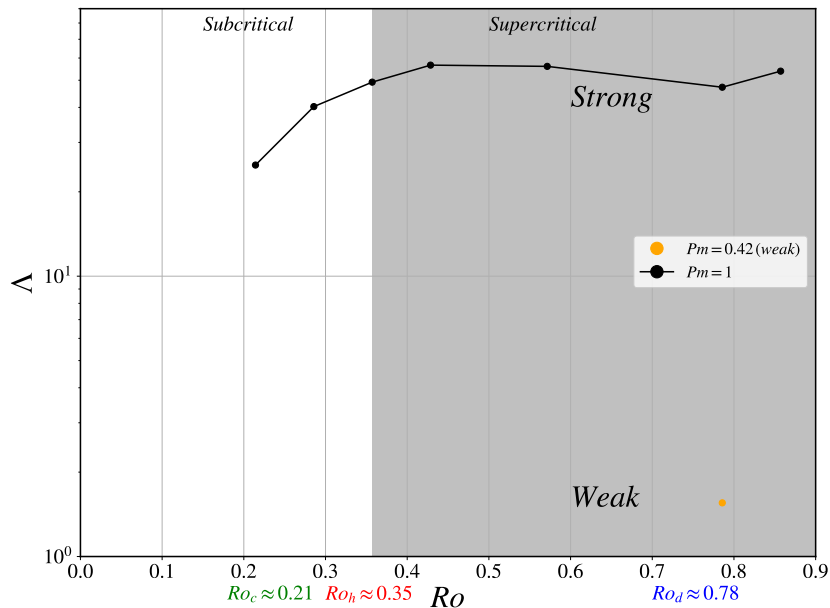
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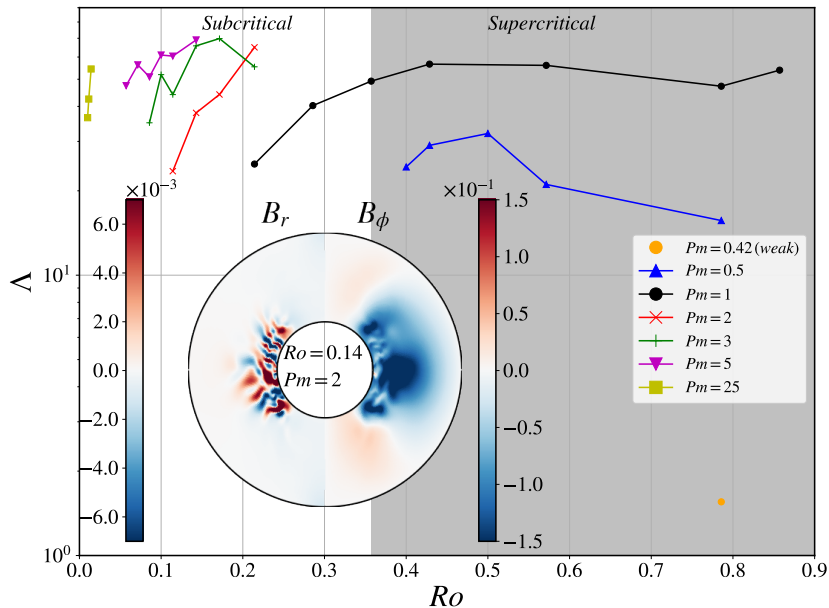
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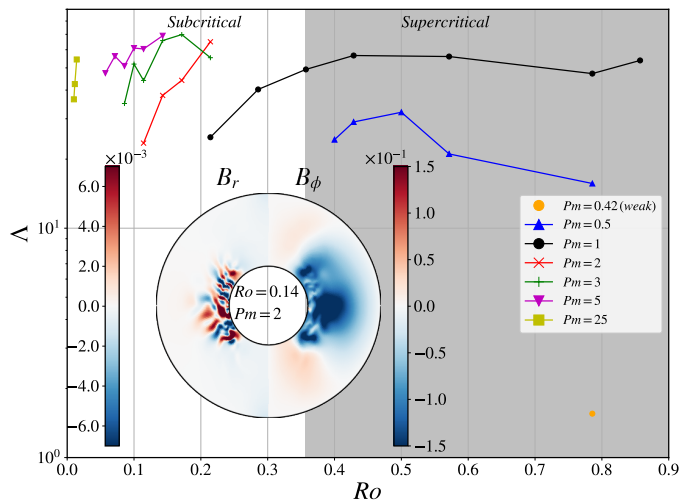


Figure 3 – Bifurcation diagram of the Elsasser number $\Lambda = \int B^2 / \rho \mu \Omega \eta dV$ versus Ro for $Ek = 10^{-5}$, $N/\Omega = 1.24$, $Pr = 0.1$ (Daniel et al. 2023).

Bifurcations

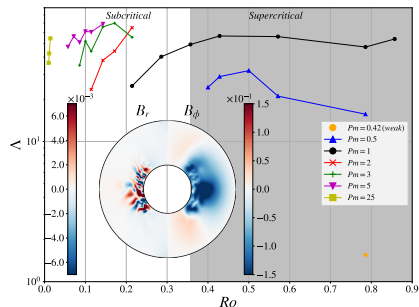


Figure 4 – Bifurcation diagram of the Elsasser number $\Lambda = \int B^2 / \rho \mu \Omega \eta dV$ versus Ro for $Ek = 10^{-5}$, $N/\Omega = 1.24$, $Pr = 0.1$

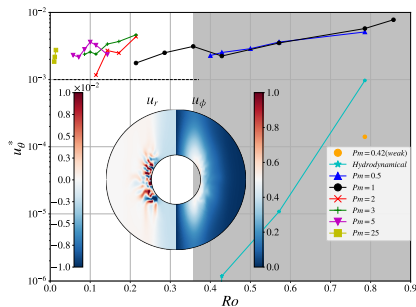


Figure 5 – Bifurcation diagram of the fluctuations of the velocity field u_θ^* versus Ro for $Ek = 10^{-5}$, $N/\Omega = 1.24$, $Pr = 0.1$

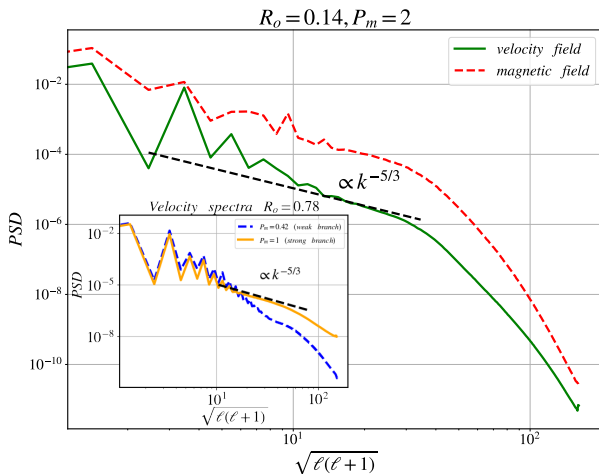


Figure 6 – *Main*: Power Spectrum Density for the run $R_o = 0.14$, $P_m = 2$, both for the velocity and magnetic field, computed in the active dynamo region ($0.4 < r < 0.6$) and time-averaged over the saturated phase. *Inset*: PSD for two runs at the same $R_o = 0.78$, but for different P_m (strong and weak branches).

Effect of the turbulence on Angular momentum transport

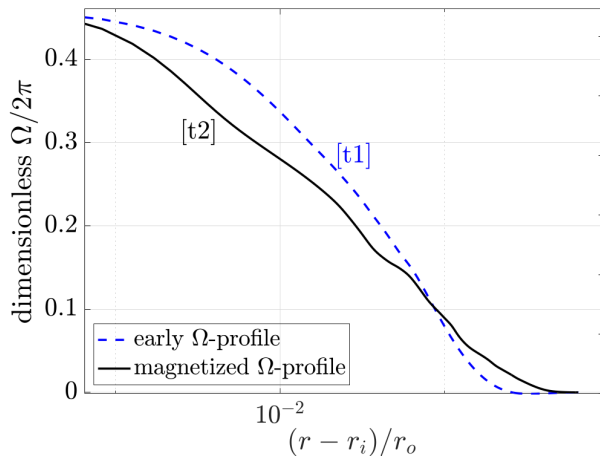


Figure 7 – Angular velocity profile Ω before the turbulence sets in (t1) and right after (t2) (*Petitdemange et al. 2023*)

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Reduced model involving (only) 3 modes :

- $m = 0$ Magnetic mode B (Dominant part of B)
- $m \neq 0$ Magnetic mode P
- $m \neq 0$ Hydrodynamic mode V (coupling via induction term)

Symmetries

- $(B, P) \rightarrow (-B, -P)$

Rotational invariance

- $V \rightarrow V e^{im\phi}, P \rightarrow P e^{im\phi}$

V grows

$$\begin{cases} \dot{B} = (-\mu &)B \\ \dot{P} = (-\nu &)P \\ \dot{V} = (\lambda &)V \end{cases} \quad (3)$$

V grows, B grows

$$\begin{cases} \dot{B} = (-\mu + |V|^2)B \\ \dot{P} = (-\nu \quad \quad)P \\ \dot{V} = (\quad \lambda \quad \quad)V \end{cases} \quad (4)$$

V grows, B grows, P grows

$$\begin{cases} \dot{B} = (-\mu + |V|^2)B \\ \dot{P} = (-\nu + |B|^2)P \\ \dot{V} = (\lambda \quad \quad) V \end{cases} \quad (5)$$

V grows, B grows, P grows

$$\begin{cases} \dot{B} = (-\mu + |V|^2)B \\ \dot{P} = (-\nu + |B|^2)P + VB \\ \dot{V} = (\lambda \quad \quad) V + \gamma BP \end{cases} \quad (6)$$

V grows, B grows, P grows, saturation

$$\begin{cases} \dot{B} = (-\mu + |V|^2)B - |B|^2B \\ \dot{P} = (-\nu + |B|^2)P + VB - |P|^2P \\ \dot{V} = (\lambda \quad \quad) V + \gamma BP - |V|^2V \end{cases} \quad (7)$$

$$\begin{cases} \dot{B} = (-\mu + c_1|V|^2)B - \alpha_1|B|^2B \\ \dot{P} = (-\nu + \beta_2|B|^2)P + c_2VB - \beta_1|P|^2P \\ \dot{V} = \lambda V + \gamma BP - c_3|V|^2V \end{cases} \quad (8)$$

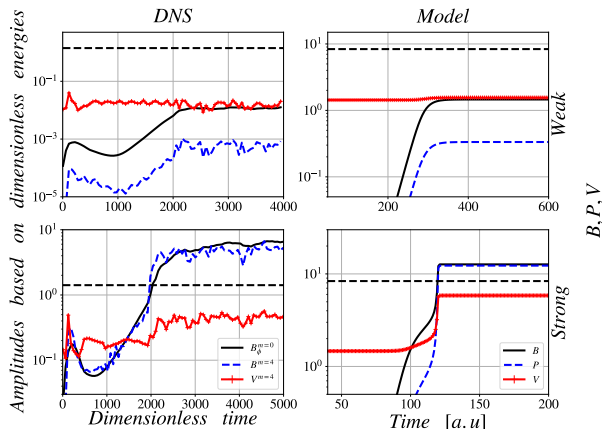


Figure 8 – Timeseries from our DNS (left) compared to the ones from the model derived here (right) for the weak (top) and the strong (bottom) dynamo branches. The horizontal dashed lines corresponds to Taylor instability threshold.

Bifurcation

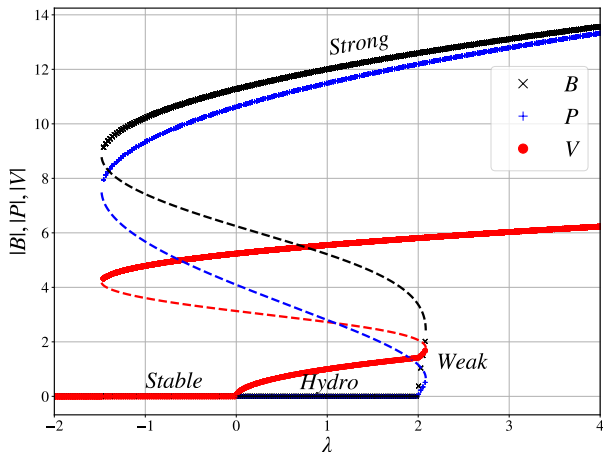
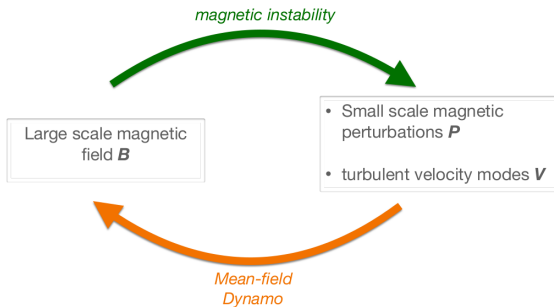


Figure 9 – Bifurcation diagram for the amplitude of each mode as λ varies. The other parameters are fixed.

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Take-home message



- Spruit H, *Dynamo action by differential rotation in a stably stratified stellar interior* AA 2002
- Petitdemange L, Marcotte F, Gissinger, *Spin-down by dynamo action in simulated radiative stellar layers* Science 2023
- Daniel F, Petitdemange L, Gissinger C, *Subcritical transition to turbulence triggered by a magnetic dynamo*, Under consideration for publication in PRF
- Dormy E, Cardin P, Jault D, *MHD flow in a slightly differentially rotating spherical shell, with conducting inner core, in a dipolar magnetic field*, EPSL 1998
- Aubert J, Aurnou J, Wicht J, *The magnetic structure of convection-driven numerical dynamos* GJI, 2007