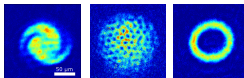


A fast rotating superfluid on a curved surface



Romain Dubessy & BEC group

Laboratoire de physique des lasers, CNRS UMR 7538
Université Sorbonne Paris Nord, Villetaneuse, France

26^{ème} Congrès de la SFP - 3–7 Juillet 2023
MC16 – Fluides classiques et quantiques hors équilibre

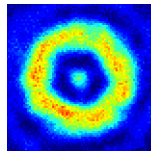


Dynamics of quantum gases: superfluidity

Specific dynamical properties

Quantum gases with weak repulsive interactions are **superfluid**. Superfluidity is a **dynamic property** with subtle effects.

- absence of **viscosity**, leading to **persistent currents** in a circular guide

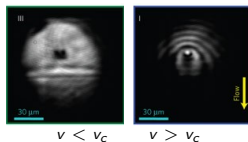
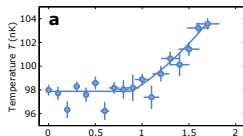


Dynamics of quantum gases: superfluidity

Specific dynamical properties

Quantum gases with weak repulsive interactions are **superfluid**.
Superfluidity is a **dynamic property** with subtle effects.

- absence of viscosity, leading to **persistent currents** in a circular guide
- **critical velocity** $v_c = c$ for excitations (Landau criterion) [data: 2D gas, Desbuquois, Nat. Phys. 2012
picture: polaritons, Amo Nat. Phys. 2009]

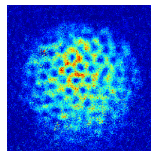


Dynamics of quantum gases: superfluidity

Specific dynamical properties

Quantum gases with weak repulsive interactions are **superfluid**. Superfluidity is a **dynamic property** with subtle effects.

- **absence of viscosity**, leading to **persistent currents** in a circular guide
- **critical velocity** $v_c = c$ for excitations (Landau criterion) [data: 2D gas, Desbuquois, Nat. Phys. 2012
picture: polaritons, Amo Nat. Phys. 2009]
- **hydrodynamic behaviour** with irrotational flow: quantized circulation and **quantized vortices** in a rotating gas



Dynamics of quantum gases: superfluidity

Specific dynamical properties

Quantum gases with weak repulsive interactions are **superfluid**. Superfluidity is a **dynamic property** with subtle effects.

- **absence of viscosity**, leading to **persistent currents** in a circular guide
- **critical velocity** $v_c = c$ for excitations (Landau criterion) [data: 2D gas, Desbuquois, Nat. Phys. 2012
picture: polaritons, Amo Nat. Phys. 2009]
- **hydrodynamic behaviour** with irrotational flow: quantized circulation and **quantized vortices** in a rotating gas

Interest of cold atoms for studying superfluidity: great **flexibility** in the control of the **confinement geometry** (harmonic traps, rings, lattices, low dimensions. . .)

Dynamics of quantum gases: superfluidity

Specific dynamical properties

Quantum gases with weak repulsive interactions are **superfluid**. Superfluidity is a **dynamic property** with subtle effects.

- **absence of viscosity**, leading to **persistent currents** in a circular guide
- **critical velocity** $v_c = c$ for excitations (Landau criterion) [data: 2D gas, Desbuquois, Nat. Phys. 2012
picture: polaritons, Amo Nat. Phys. 2009]
- **hydrodynamic behaviour** with irrotational flow: quantized circulation and **quantized vortices** in a rotating gas

Interest of cold atoms for studying superfluidity: great **flexibility** in the control of the **confinement geometry** (harmonic traps, rings, lattices, low dimensions. . .)

In this talk: a superfluid rotating in a **bubble** trap.

Physics in a bubble

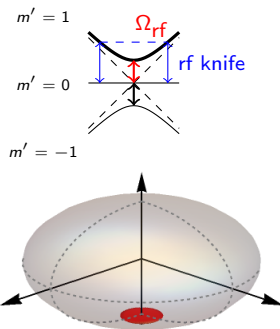
rf-induced adiabatic potentials — the dressed quadrupole trap

Adiabatic potentials for rf-dressed atoms

Ingredients: inhomogeneous B field + strong rf field, coupling Ω_{rf}

Here: quadrupole field, magnetic gradient b'

- local B and rf fields: atomic spin follows **adiabatically** a local eigenstate
- local eigenenergy acts as a **potential**
- atoms are **strongly confined** to a **resonant isomagnetic surface**
 $\mu B(\mathbf{r}) = \hbar\omega$
- **smooth** surface potentials
- cooling with an **rf knife**.



For a quadrupole field: ellipsoidal **isomagnetic surface**

$$x^2 + y^2 + 4z^2 = r_0^2$$

$$\text{with } r_0 \propto \omega/b'.$$

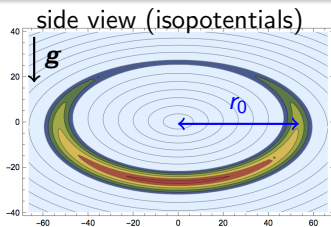
[reviews Garraway/Perrin: JPB 2016 and Adv.At.Mol.Opt.Phys. 2017]

Trapping atoms on a surface

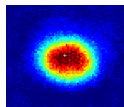
A smooth two-dimensional trap

$$\omega_z \propto \frac{b'}{\sqrt{\Omega_{\text{rf}}}} \sim 1 \text{ kHz}$$

$$\omega_x, \omega_y \propto \sqrt{\frac{g}{r_0}} \sim 20\text{-}50 \text{ Hz}$$

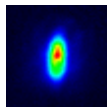


top view



in situ

side view



tof

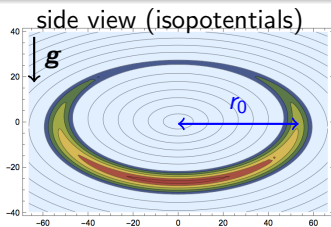
- very flat $\omega_z \gg \omega_{x,y}$
- in-plane anisotropy $\eta = \frac{\omega_x}{\omega_y}$ controlled through rf polarization:
- **rotationally invariant** ($\eta = 1$) for a σ^+ polarization along z
- **anisotropic** ($\eta \neq 1$) for **linear** horizontal polarization

Trapping atoms on a surface

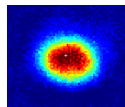
A smooth two-dimensional trap

$$\omega_z \propto \frac{b'}{\sqrt{\Omega_{\text{rf}}}} \sim 1 \text{ kHz}$$

$$\omega_x, \omega_y \propto \sqrt{\frac{g}{r_0}} \sim 20\text{-}50 \text{ Hz}$$

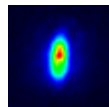


top view



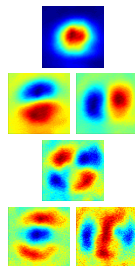
in situ

side view



tof

- very flat $\omega_z \gg \omega_{x,y}$
- in-plane anisotropy $\eta = \frac{\omega_x}{\omega_y}$ controlled through rf polarization:
- **rotationally invariant** ($\eta = 1$) for a σ^+ polarization along z
- **anisotropic** ($\eta \neq 1$) for **linear** horizontal polarization
- geometry can be modified **dynamically**
- ideal for the study of the 2D trapped gas **dynamics**



[Dubessy NJP 2014]

Why rotations ?

Quantum Hall effect with atoms

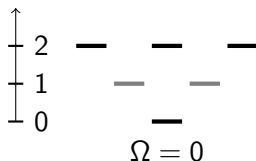
ideal 2D trapped
rotating atomic gas

\leftrightarrow

2D electron gas with a
uniform magnetic field

$$H = \hbar\omega_r \left(\hat{a}_x^\dagger \hat{a}_x + \hat{a}_y^\dagger \hat{a}_y \right) - \Omega \hat{L}_z$$

$$E = n\hbar\omega_r$$



Why rotations ?

Quantum Hall effect with atoms

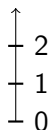
ideal 2D trapped
rotating atomic gas

\leftrightarrow

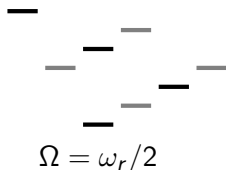
2D electron gas with a
uniform magnetic field

$$H = \hbar\omega_r \left(\hat{a}_x^\dagger \hat{a}_x + \hat{a}_y^\dagger \hat{a}_y \right) - \Omega \hat{L}_z$$

$$E = n\hbar\omega_r$$



$$\Omega = 0$$



$$\Omega = \omega_r/2$$

Why rotations ?

Quantum Hall effect with atoms

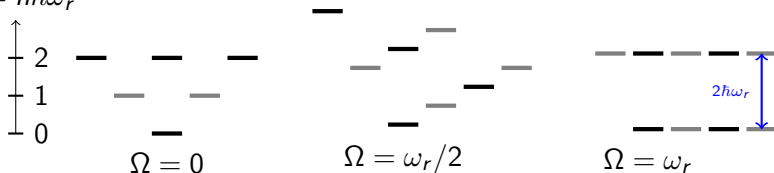
ideal 2D trapped
rotating atomic gas

\leftrightarrow

2D electron gas with a
uniform magnetic field

$$H = \hbar\omega_r \left(\hat{a}_x^\dagger \hat{a}_x + \hat{a}_y^\dagger \hat{a}_y \right) - \Omega \hat{L}_z$$

$$E = n\hbar\omega_r$$



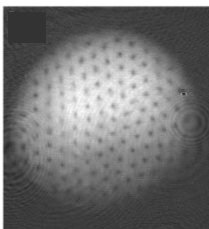
- Landau level structure \Rightarrow highly degenerate groundstate

[Fetter RMP 2009, Fletcher Science 2021 & gauge fields: Chalopin Nat. Phys. 2020]

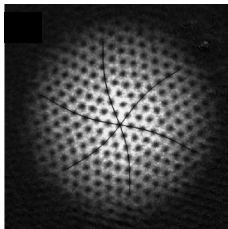
- Small **energy gap** \Rightarrow increased **role of temperature**

Vortex crystals

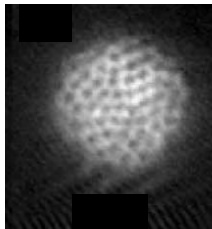
Low energy modes & melting transition



[Abo-Shaeer Science 2001]



[Coddington PRL 2003]

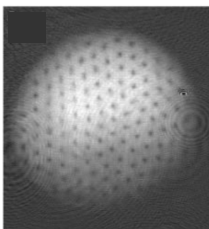


[Bretin PRL 2004]

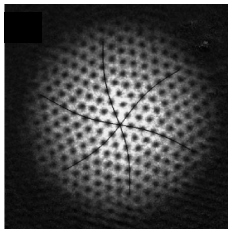
- for $\Omega \leq \omega_r$ groundstate is a large **Abrikosov lattice**
- with well defined modes:
 - longitudinal (Kelvin) [Pitaevskii Sov. Phys. JETP 1961]
 - in plane: elasticity of the lattice [Tkachenko Sov. Phys. JETP 1966]
- **Thermal population** of the modes can melt the lattice...
never observed in BEC !

Vortex crystals

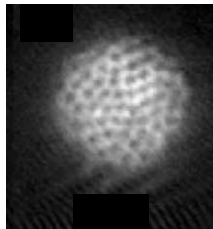
Low energy modes & melting transition



[Abo-Shaeer Science 2001]
 $\omega_r/\omega_z = 4.2$



[Coddington PRL 2003]
1.6

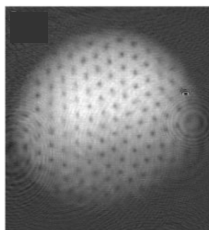


[Bretin PRL 2004]
5.89

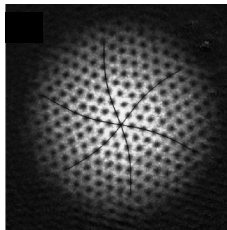
- for $\Omega \leq \omega_r$ groundstate is a large **Abrikosov lattice**
- with well defined modes:
 - longitudinal (Kelvin) [Pitaevskii Sov. Phys. JETP 1961]
 - in plane: elasticity of the lattice [Tkachenko Sov. Phys. JETP 1966]
- **Thermal population** of the modes can melt the lattice...
never observed in BEC !

Vortex crystals

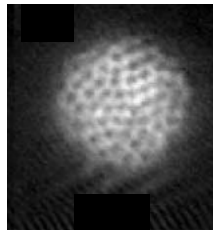
Low energy modes & melting transition



[Abo-Shaeer Science 2001]
 $\omega_r/\omega_z = 4.2$



[Coddington PRL 2003]
1.6 $T_m/T_c \sim 0.78$



[Bretin PRL 2004]
5.89

- for $\Omega \leq \omega_r$ groundstate is a large **Abrikosov lattice**
- with well defined modes:
 - longitudinal (Kelvin) [Pitaevskii Sov. Phys. JETP 1961]
 - in plane: elasticity of the lattice [Tkachenko Sov. Phys. JETP 1966]
- **Thermal population** of the modes can melt the lattice...
never observed in BEC !

To lower T_m/T_c : decrease $\omega_r/\omega_z \Rightarrow$ **go 2D** ! **bubble**: $\omega_r/\omega_z \leq 0.1$

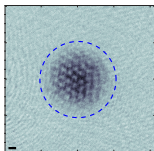
Vortex lattice in fast rotating bubble trap

A first quick look

- Start from a a degenerate cloud at rest at the bottom of the bubble, $\omega_r = 2\pi \times 34$ Hz
- Induce an **in-plane elliptic deformation** rf polarization

$$V(r) = M\omega_r^2/2 \times [(1 - \epsilon)x'^2 + (1 + \epsilon)y'^2] + \dots$$
- **Rotate** the trap main axes x' , y' at frequency Ω_{rot}
- Restore the rotationally invariant trap

Increasing $\Omega_{\text{rot}} \dots$



$\Omega_{\text{rot}}/(2\pi) = 20$ Hz

Blue dashed circle: Thomas-Fermi radius after 27 ms time-of-flight

$\epsilon = 0.18$

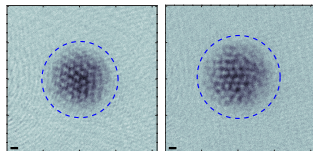
Vortex lattice in fast rotating bubble trap

A first quick look

- Start from a degenerate cloud at rest at the bottom of the bubble, $\omega_r = 2\pi \times 34$ Hz
- Induce an **in-plane elliptic deformation** rf polarization

$$V(r) = M\omega_r^2/2 \times [(1 - \epsilon)x'^2 + (1 + \epsilon)y'^2] + \dots$$
- **Rotate** the trap main axes x' , y' at frequency Ω_{rot}
- Restore the rotationally invariant trap

Increasing $\Omega_{\text{rot}} \dots$



$\Omega_{\text{rot}}/(2\pi) = 20$ Hz

21 Hz

Blue dashed circle: Thomas-Fermi radius after 27 ms time-of-flight

$\epsilon = 0.18$

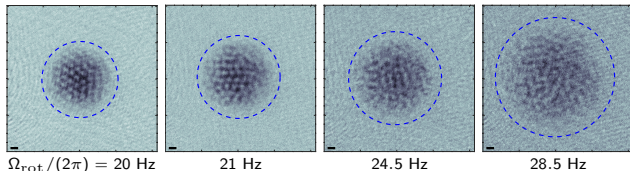
Vortex lattice in fast rotating bubble trap

A first quick look

- Start from a degenerate cloud at rest at the bottom of the bubble, $\omega_r = 2\pi \times 34$ Hz
- Induce an **in-plane elliptic deformation** rf polarization

$$V(r) = M\omega_r^2/2 \times [(1 - \epsilon)x'^2 + (1 + \epsilon)y'^2] + \dots$$
- **Rotate** the trap main axes x' , y' at frequency Ω_{rot}
- Restore the rotationally invariant trap

Increasing Ω_{rot} ... **disordered lattice**...



Blue dashed circle: Thomas-Fermi radius after 27 ms time-of-flight

$\epsilon = 0.18$

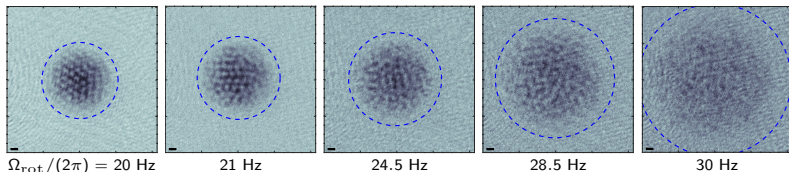
Vortex lattice in fast rotating bubble trap

A first quick look

- Start from a degenerate cloud at rest at the bottom of the bubble, $\omega_r = 2\pi \times 34$ Hz
- Induce an **in-plane elliptic deformation** rf polarization

$$V(r) = M\omega_r^2/2 \times [(1 - \epsilon)x'^2 + (1 + \epsilon)y'^2] + \dots$$
- **Rotate** the trap main axes x' , y' at frequency Ω_{rot}
- Restore the rotationally invariant trap

Increasing Ω_{rot} ... **disordered lattice**... **melting** ?



Blue dashed circle: Thomas-Fermi radius after 27 ms time-of-flight

$\epsilon = 0.18$

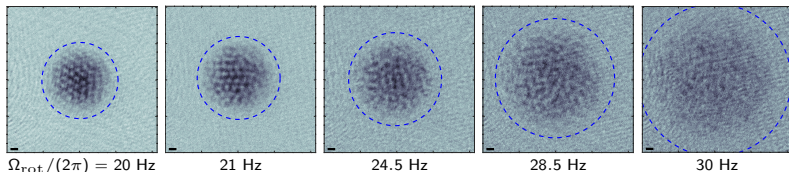
Vortex lattice in fast rotating bubble trap

A first quick look

- Start from a degenerate cloud at rest at the bottom of the bubble, $\omega_r = 2\pi \times 34$ Hz
- Induce an **in-plane elliptic deformation** rf polarization

$$V(r) = M\omega_r^2/2 \times [(1 - \epsilon)x'^2 + (1 + \epsilon)y'^2] + \dots$$
- **Rotate** the trap main axes x' , y' at frequency Ω_{rot}
- Restore the rotationally invariant trap

Increasing Ω_{rot} ... **disordered lattice**... **melting** ?



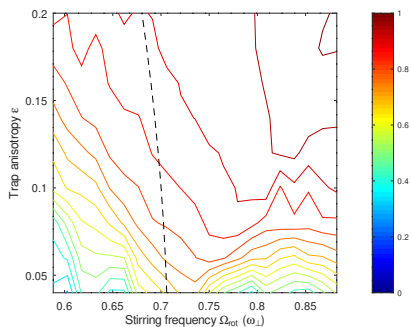
Blue dashed circle: Thomas-Fermi radius after 27 ms time-of-flight

$\epsilon = 0.18$

We need to **measure** N , Ω , T and **compare to** $T_m(N, \Omega)$...

Rotation control

Stirring and evaporating



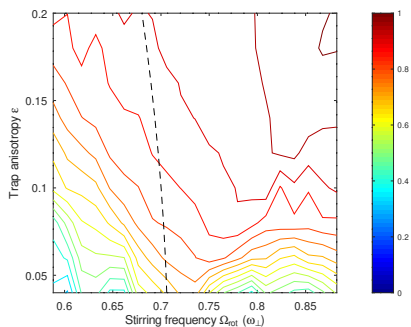
Effective rotation Ω/ω_r

with post-stirring rf ramp 80 \rightarrow 60 kHz

- allows to reach $\Omega \sim \omega_r$
- at constant $T \sim 18$ nK

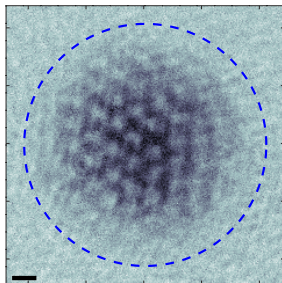
Rotation control

Stirring and evaporating



Counting vortices

- top view, 27 ms tof



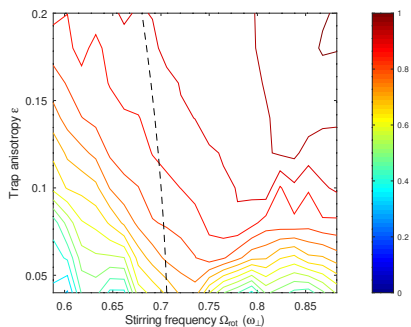
Effective rotation Ω/ω_r

with post-stirring rf ramp 80 \rightarrow 60 kHz

- allows to reach $\Omega \sim \omega_r$
- at constant $T \sim 18$ nK

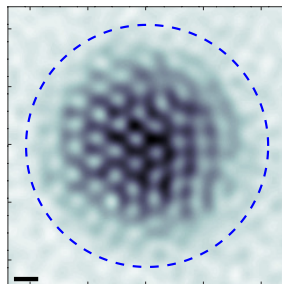
Rotation control

Stirring and evaporating



Counting vortices

- top view, 27 ms tof
- enhance visibility



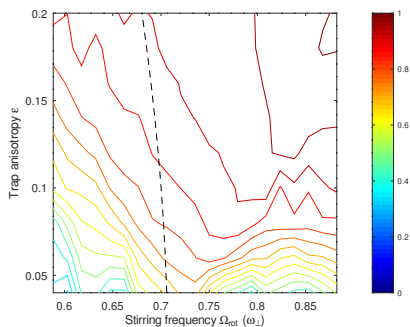
Effective rotation Ω/ω_r

with post-stirring rf ramp 80 \rightarrow 60 kHz

- allows to reach $\Omega \sim \omega_r$
- at constant $T \sim 18$ nK

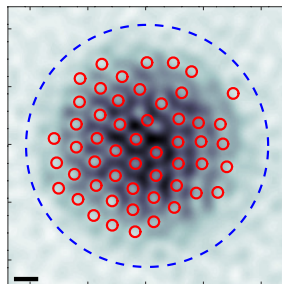
Rotation control

Stirring and evaporating



Counting vortices

- top view, 27 ms tof
- enhance visibility
- detect positive curvature



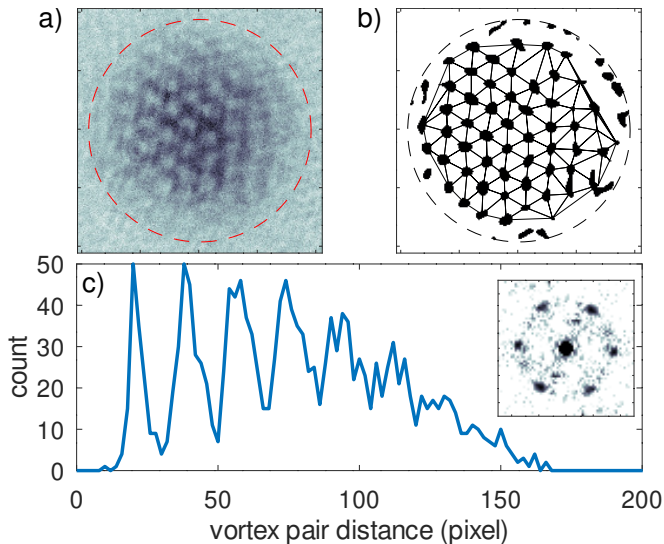
Effective rotation Ω/ω_r

with post-stirring rf ramp 80 \rightarrow 60 kHz

- allows to reach $\Omega \sim \omega_r$
- at constant $T \sim 18$ nK

Quantitative study of the vortex lattice

Vortex-vortex correlations



Thermal melting of the vortex lattice ?

Studying a quasi-2D crystal

quasi 2D rotating Bose gas $T_m \leq 0.23 T_{BKT}$ [Gifford PRA 2008]
upper bound on melting temperature
(computed using low energy modes of the crystal lattice & KTHNY theory)
studied in many systems: superconductors, colloids, ... [Gasser ChemPhysChem 2010]

Thermal melting of the vortex lattice ?

Studying a quasi-2D crystal

quasi 2D rotating Bose gas $T_m \leq 0.23 T_{BKT}$

[Gifford PRA 2008]

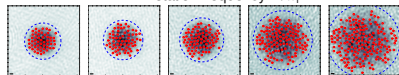
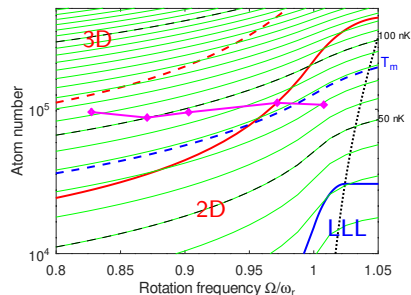
upper bound on melting temperature

(computed using low energy modes of the crystal lattice & KTHNY theory)

studied in many systems: superconductors, colloids, ... [Gasser ChemPhysChem 2010]

What is T_{BKT} ?

semi-classical + LDA + \mathcal{D}_c from QMC



Thermal melting of the vortex lattice ?

Studying a quasi-2D crystal

quasi 2D rotating Bose gas $T_m \leq 0.23 T_{BKT}$

[Gifford PRA 2008]

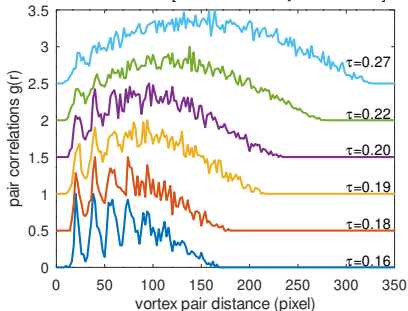
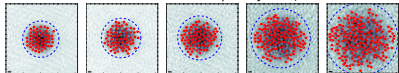
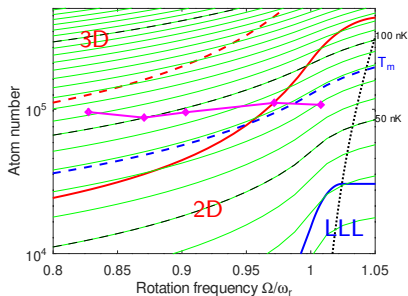
upper bound on melting temperature

(computed using low energy modes of the crystal lattice & KTHNY theory)

studied in many systems: superconductors, colloids, ... [Gasser ChemPhysChem 2010]

What is T_{BKT} ?

semi-classical + LDA + \mathcal{D}_c from QMC



- Remember: rf-knife sets the temperature $T \simeq 18$ nK,
 $\tau = T/T_{BKT}$

- Qualitative agreement: change in $g(r)$ as T/T_m varies.

Can we rotate even faster ?

Fighting the centrifugal force

To restore the trapping potential, add a quartic term to $V(r)$:

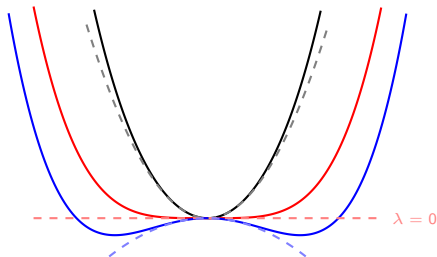
$$V_{\text{eff}}(r) = \frac{M}{2}(\omega_r^2 - \Omega^2)r^2 + \lambda r^4.$$

[Bretin PRL 2004]

$$\Omega = 0$$

$$\Omega = \omega_r$$

$$\Omega = 1.15 \omega_r$$



\Rightarrow the bubble trap has higher order terms.

Theoretical predictions

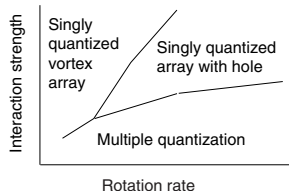
Rotating beyond the trapping frequency

Giant vortex in a harmonic + quartic trap:

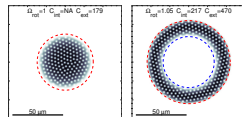
vortex
lattice



dynamical
ring



giant vortex

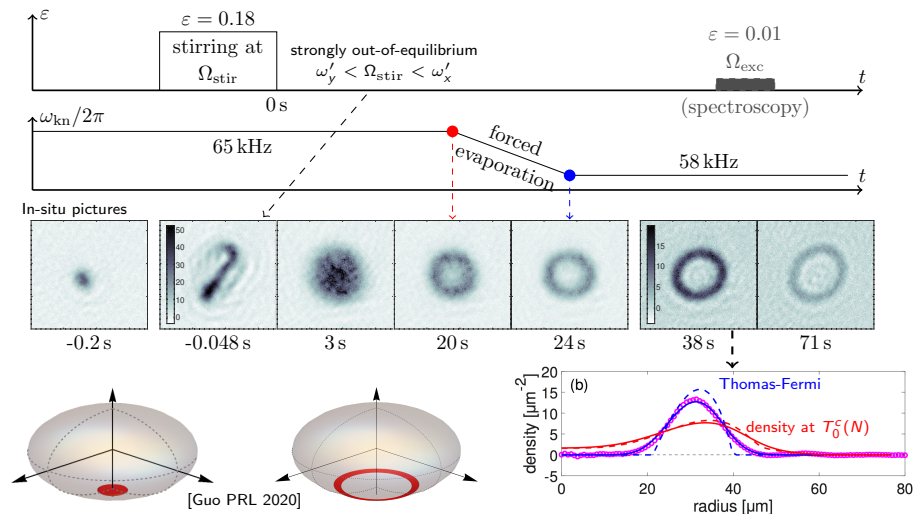


[Kavoulakis NJP 2003, Fetter PRA 2005]

GP simulation for the bubble
(quartic approximation)

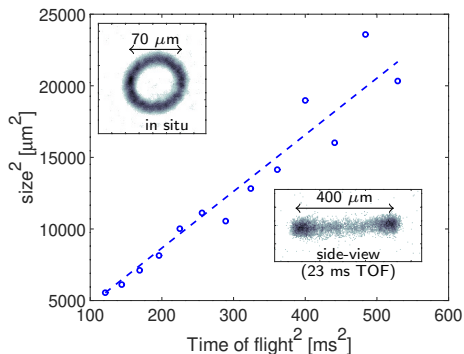
Creating a dynamical ring

Using the spin-up evaporation



A supersonic flow

Measuring the rotation from time-of-flight expansion



- size² scales as t_{TOF}^2
(ballistic expansion)
- fit gives: $\Omega \sim 1.05\omega_r$,
i.e. $v = 7.4$ mm/s
- peak density
 $n_0 \sim 15 \mu\text{m}^{-2}$
 $\Rightarrow c_0 = 0.4$ mm/s

A degenerate gaz flowing at **Mach 18** !

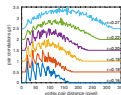
[see also Pandey Nature 2019]

Summary & prospects

Fast rotations on a shell

A very smooth and tunable shell trap to study **fast rotations**

- Fine control of the **effective rotation**
- Vortex **lattice melting** for $\Omega \sim \omega_r$
- Formation of a **long-lived dynamical ring** flowing at Mach 18 for tens of second for $\Omega > \omega_r$



- ⇒ investigate the decay mechanisms
- ⇒ test the melting scenario (KTHNY)
- ⇒ play with the curvature in the rotating frame

obstacle in a supersonic flow

loss of orientational order

β effect

Thanks for your attention !

The BEC group at Villetaneuse



front row: **R. Sharma**, M. Nouama, S. Thomas, **H. Perrin**, L. Longchambon, M. Ballu, A. Perrin
 behind: R. Dubessy, T. Badr, S. Cuk, **D. Rey**, K. Lamraoui

ANR funded postdoc position available [2 years]

Former PhDs:

Collaborations (on going)



M. de Goër
de Herve



Y. Guo



V. Bagnato



M. Olshanii



A. Minguzzi



S. Nazarenko

www-lpl.univ-paris13.fr/bec