

Freezing of nonlinear waves over an uneven bottom

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in collaboration with:

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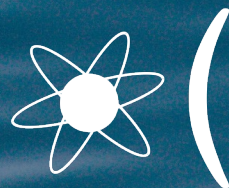


Amin Chabchoub

Stefano Trillo



Huber Branger & Christopher Luneau



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Congrès Général SFP 2023
Paris 03.07.2023

Nonlinear Schrödinger equation (NLS)

Focusing,
1D cubic

$$i\psi_z + \frac{1}{2}\psi_{tt} + |\psi|^2\psi = 0$$

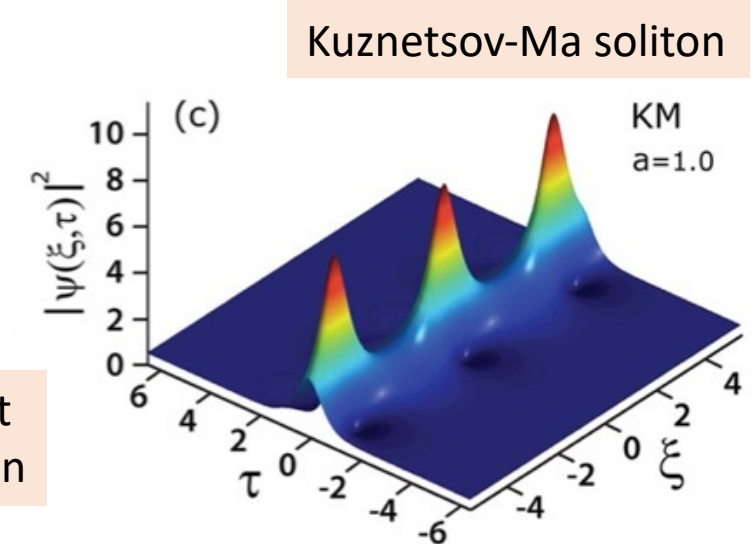
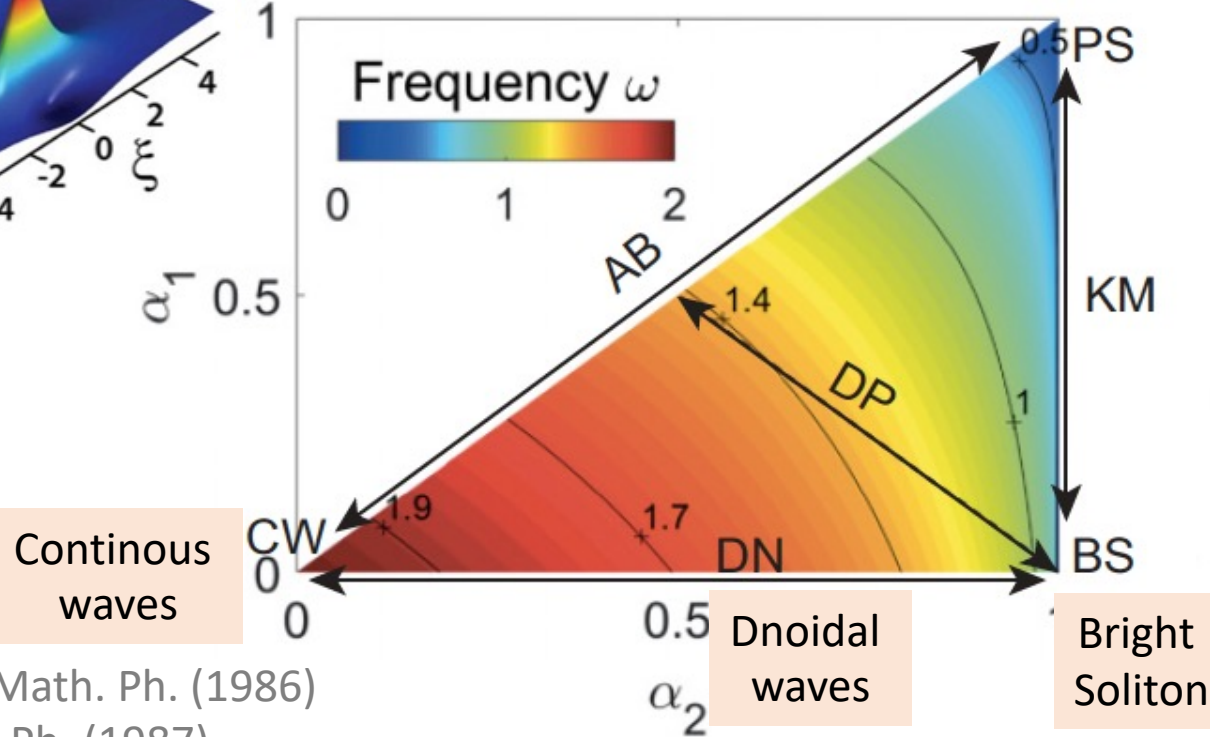
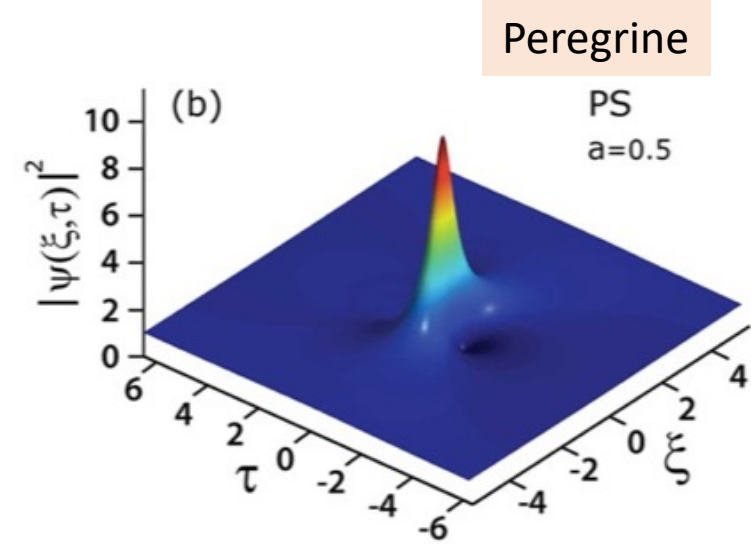
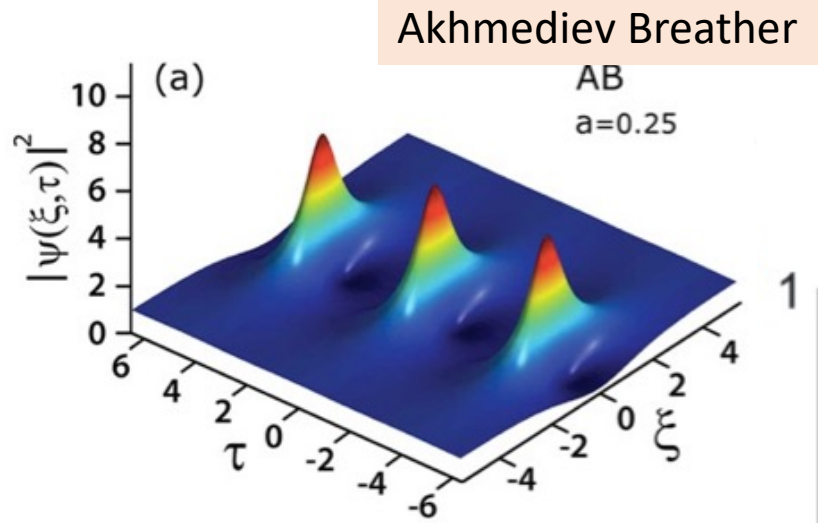
Dispersion

Nonlinearity

APPLICATIONS:
Optics
Hydrodynamics
Plasma
Cold atoms
Galaxies...



Nonlinear Schrödinger equation (NLS): three-parameter family of solutions (arbitrary amplitude and periods)

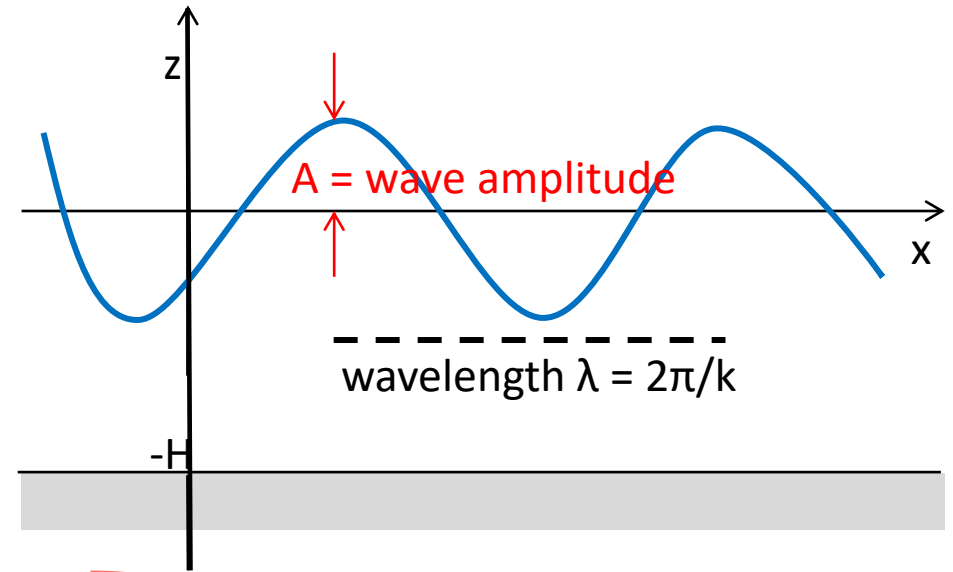


Akhmediev & Korneev *Theor. Math. Ph.* (1986)
 Akhmediev *et al.* *Theor. Math. Ph.* (1987)
 Toenger *et al.* *Sc. Rep.* (2015)
 Conforti *et al.* *PRA* (2020)

From Navier-Stokes eq. to NLS

$$\frac{D\vec{v}}{dt} = -2\vec{\Omega} \times \vec{v} - \frac{\vec{\nabla}P}{\rho} + \vec{g} + \vec{F}$$

Coriolis term
Gravity
Viscosity term



- Incompressible fluid: constant density
- Large Rossby number: Coriolis term can be neglected
- Viscosity effect neglected
- Constant bathymetry
- No wind forcing
- Velocity can be expressed in terms of a scalar potential function
- 1-dimensional propagation
- $kA \ll 1$: weak nonlinearities
- $kH \gg 1$: deep-water limit

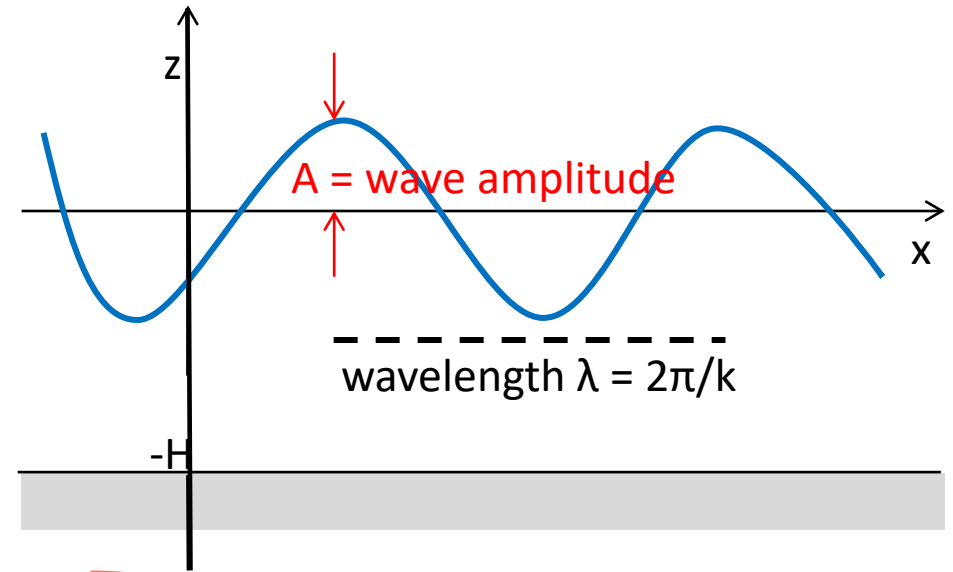
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$$i\psi_z + \frac{1}{2}\psi_{tt} + |\psi|^2\psi = 0$$



Outline

- **Bathymetry effect**

- an abrupt bathymetry change can freeze an Akhmediev Breather at max amplitude
- simplified model with only 3 waves
- experiment in a water tank

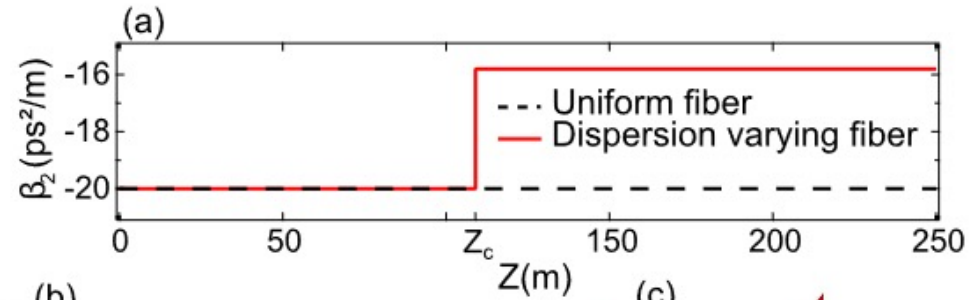
- **Wind effect**

- change between types of solutions
- experiment in a water tank with wind turbine

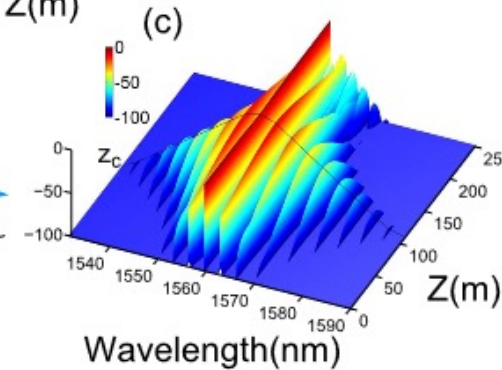
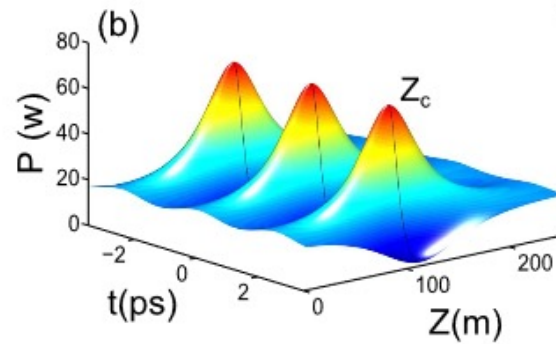


Starting idea: experiment in an optical fiber with varying dispersion

Dispersion coefficient

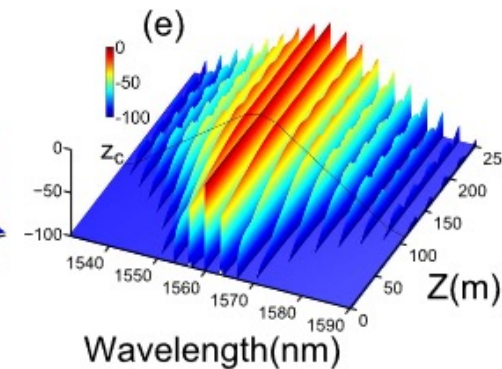
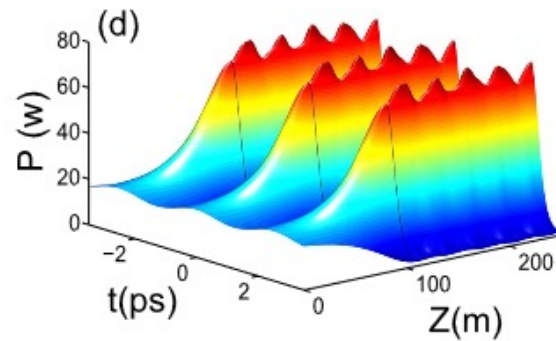


Uniform fiber



Akhmediev breather

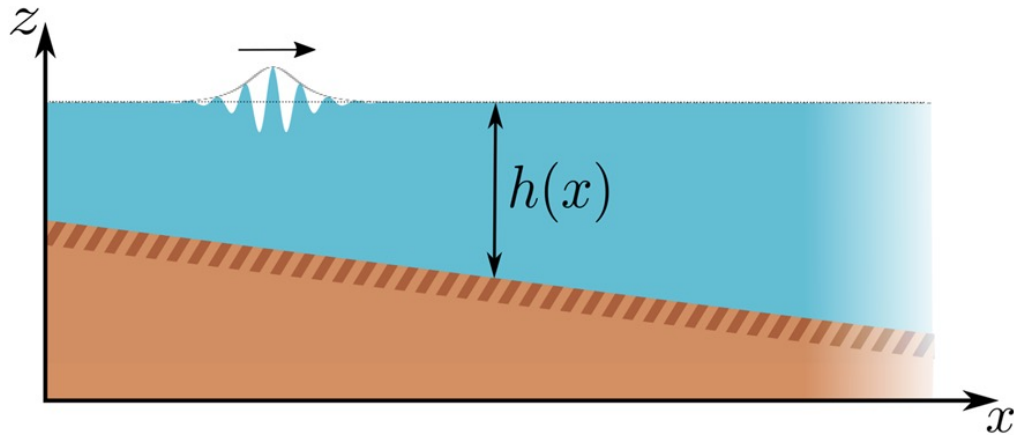
Dispersion varying fiber



Evolution can be quasi-stabilized in space at maximal compression

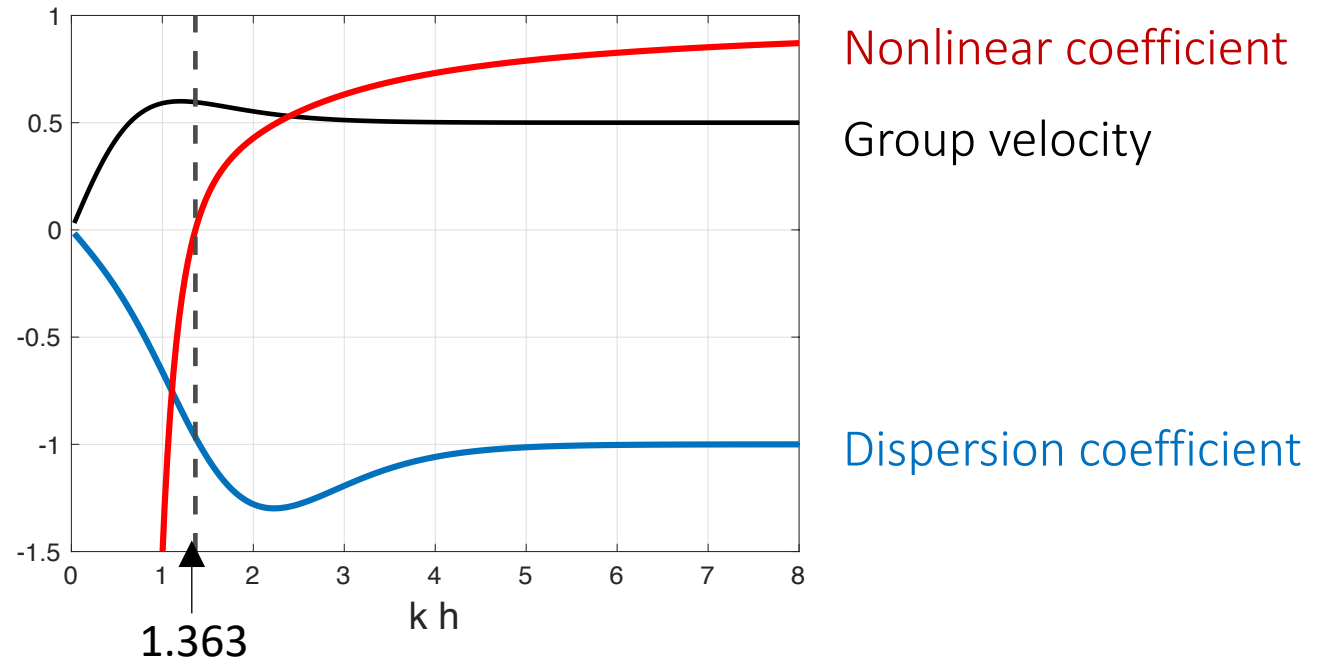


Can we reproduce the stabilization in surface water waves?



Uneven bottom

$$i\partial_{\xi}U + \underbrace{\alpha(kh)\partial_{\tau}^2U}_{\text{Dispersion}} - \underbrace{\beta(kh)|U|^2U}_{\text{Non linearity}} = -\underbrace{i\mu_0\partial_{\xi}(kh)U}_{\text{Shoaling}} - \underbrace{i\nu U}_{\text{Loss}}$$

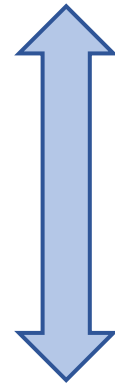


Derivation of a NLSE-like equation

$$i\partial_\xi U + \underbrace{\alpha(kh)\partial_\tau^2 U}_{\text{Dispersion}} - \underbrace{\beta(kh)|U|^2 U}_{\text{Non linearity}} = - \underbrace{i\mu_0\partial_\xi(kh)U}_{\text{Shoaling}} - \underbrace{i\nu U}_{\mu^{\text{loss}}}$$

$$U = V \exp \left[- \int_0^\xi \mu(y) dy - \nu\xi \right]$$

$$\tilde{\beta} = \beta(\xi) \frac{c_g(0)}{c_g(\xi)} \exp(-2\nu\xi)$$



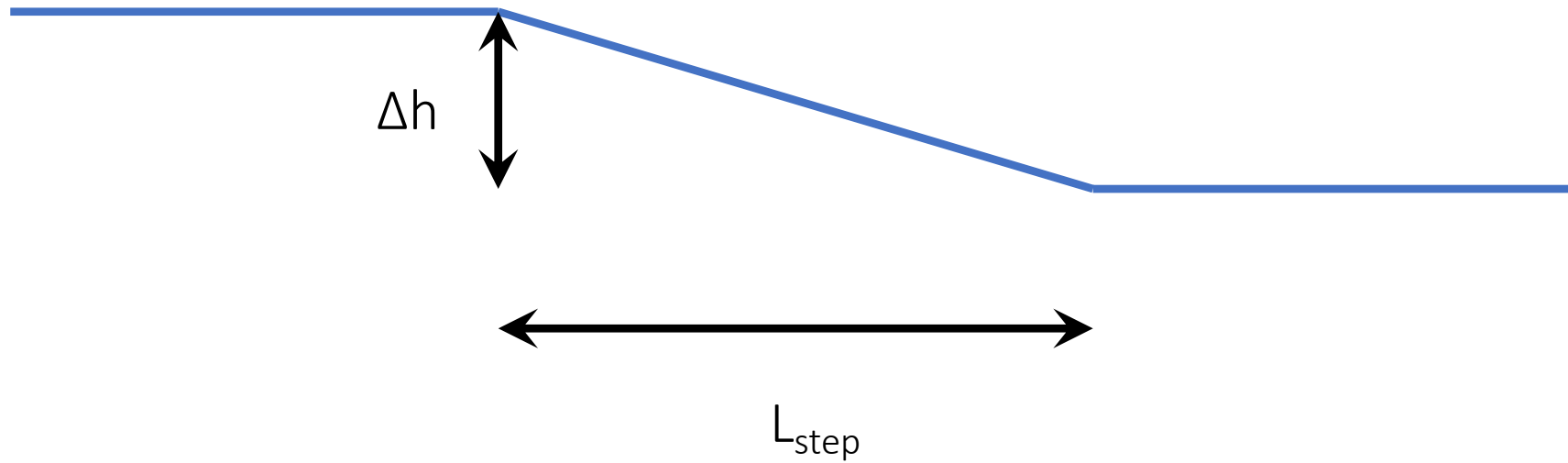
$$i\partial_\xi V + \underbrace{\alpha(kh)\partial_\tau^2 V}_{\text{Dispersion}} - \underbrace{\tilde{\beta}(kh)|V|^2 V}_{\text{Non linearity}} = 0$$



Uneven bottom

Akhmediev breather
before the depth change

Dnoidal
after the depth change

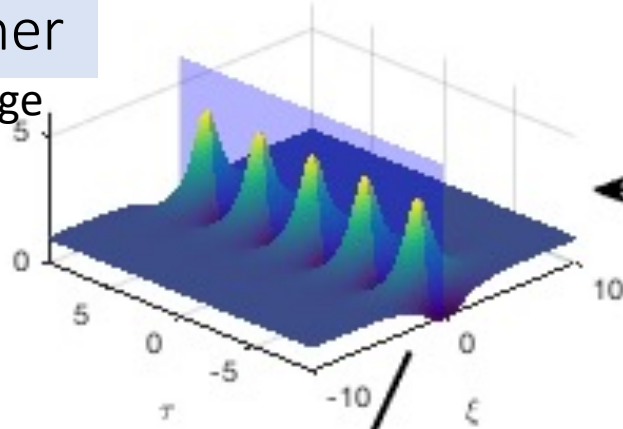


$$\Delta h \ll L_{\text{step}} \ll L_{\text{nonlinear}}$$

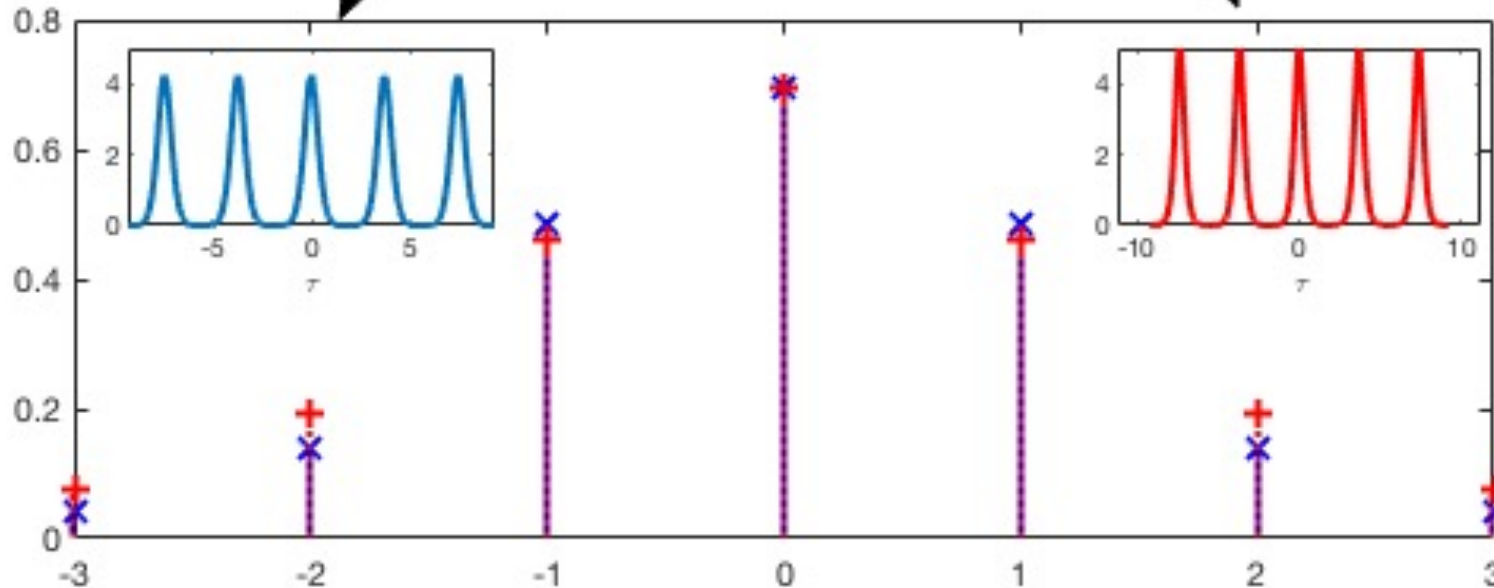
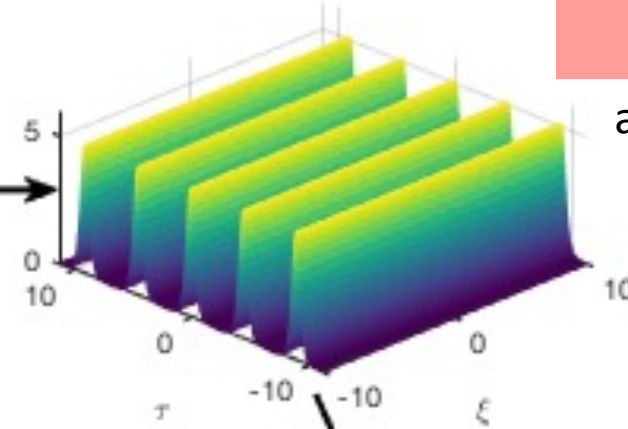


Matching NLS solutions

Akhmediev breather
before the depth change



Dnoidal
after the depth change



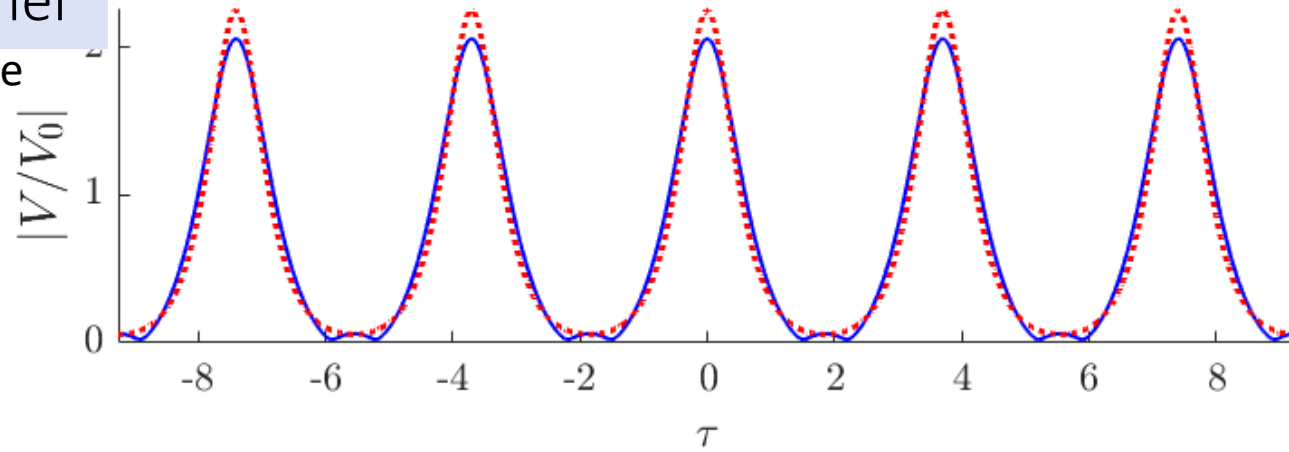
Spectrum

Gomel et al. PRL (2021)



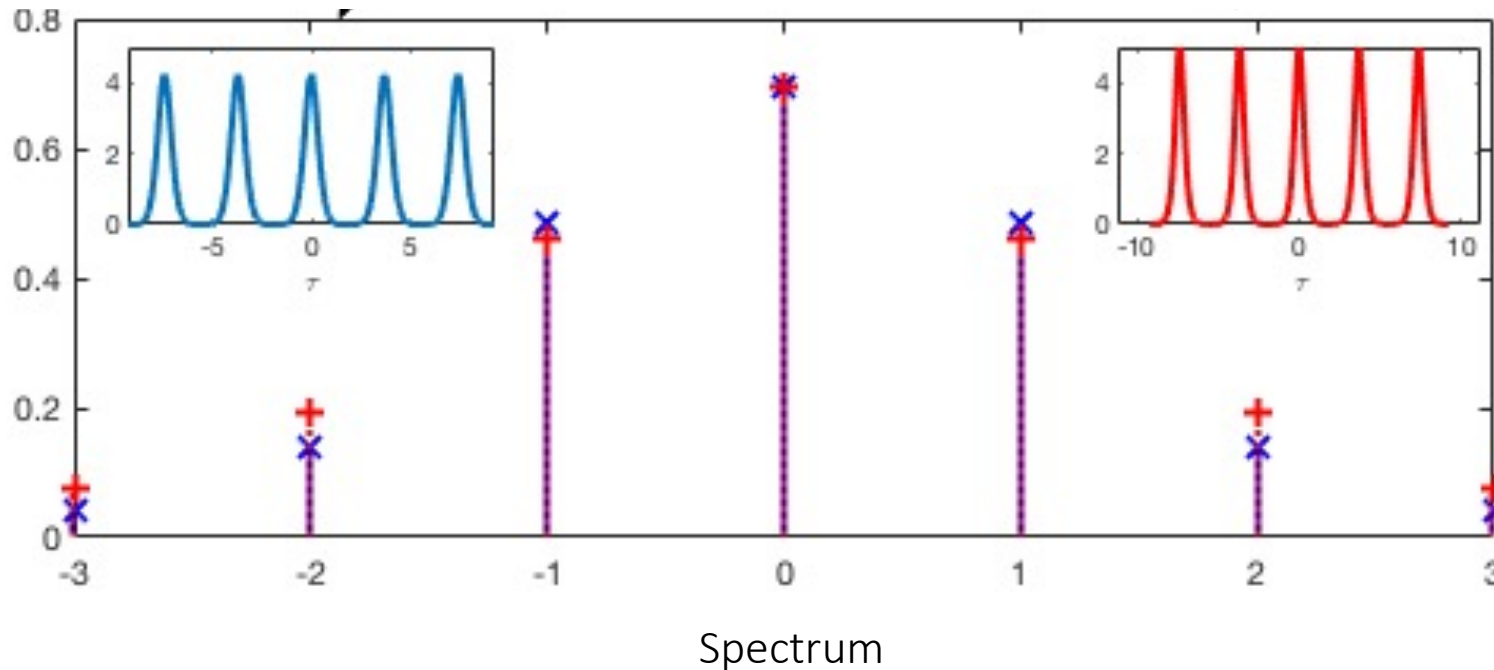
Matching NLS solutions

Akhmediev breather
before the depth change



Dnoidal

after the depth change



Simplified model: three-wave truncation

$$V(\xi, \tau) = A_0(\xi) + A_1(\xi)e^{i\Omega\tau} + A_{-1}(\xi)e^{-i\Omega\tau}$$

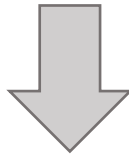
Conversion rate
to sidebands:

$$\eta = \frac{|A_1|^2 + |A_{-1}|^2}{E}$$

$$E = |A_0|^2 + |A_1|^2 + |A_{-1}|^2$$

Relative phase:

$$\psi = \frac{\phi_1 + \phi_{-1}}{2} - \phi_0$$



$$H(\psi, \eta) = \eta(\eta - 1) \cos(2\psi) + \frac{3\eta^2}{4} + \gamma\eta$$

$$\text{with: } \gamma = - \left[\frac{\alpha\Omega^2}{\tilde{\beta}E} + 1 \right]$$



Simplified model: three-wave truncation

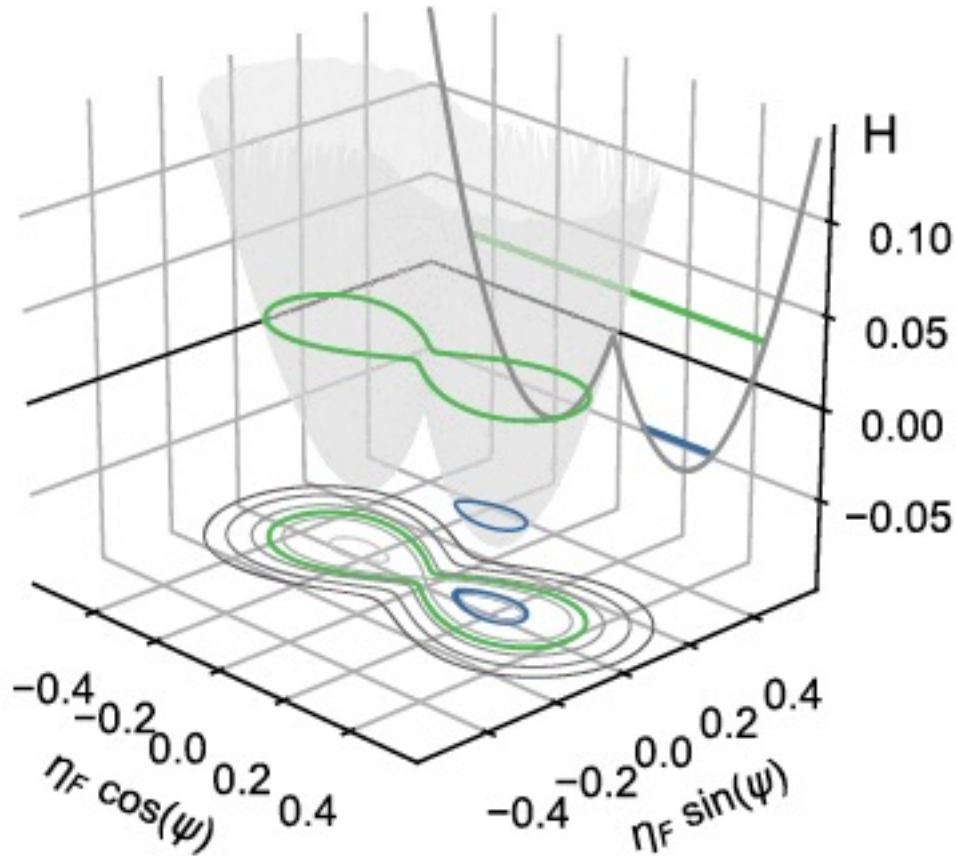
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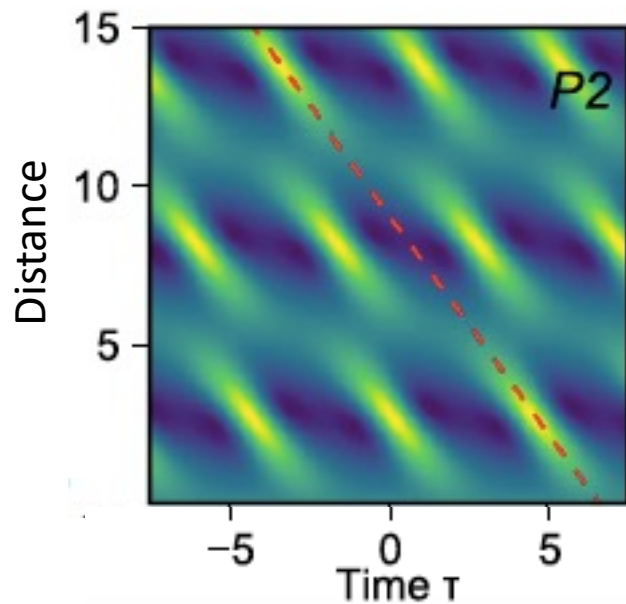
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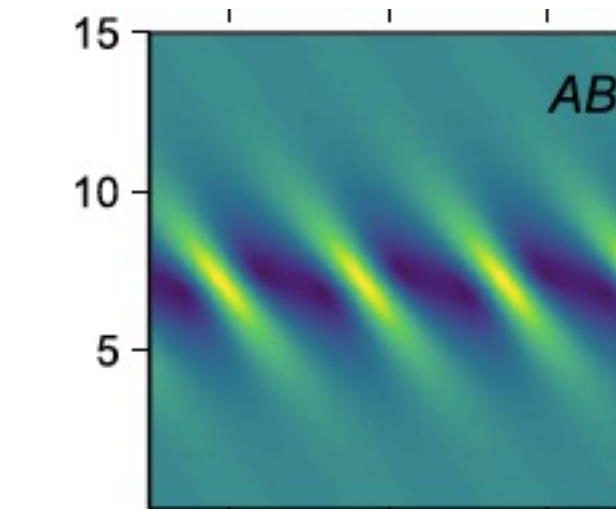
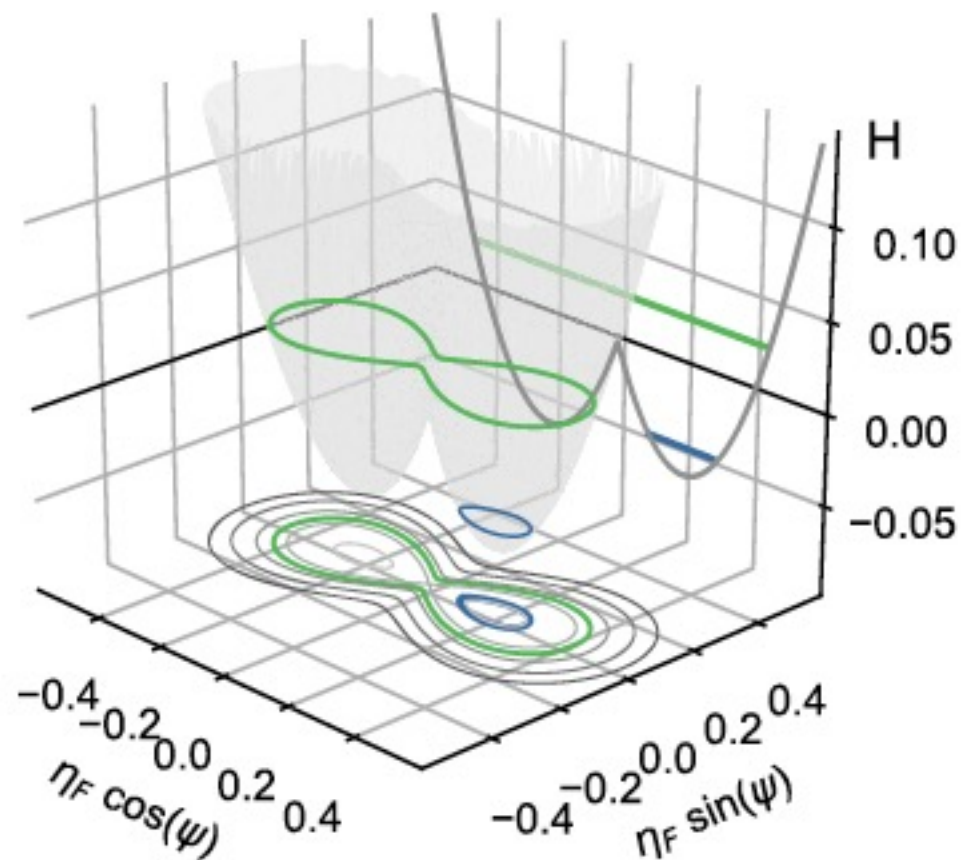
Trillo & Wabnitz *Opt. Lett.* (1991)
Armaroli, Brunetti, Kasparian *PRE* (2017)



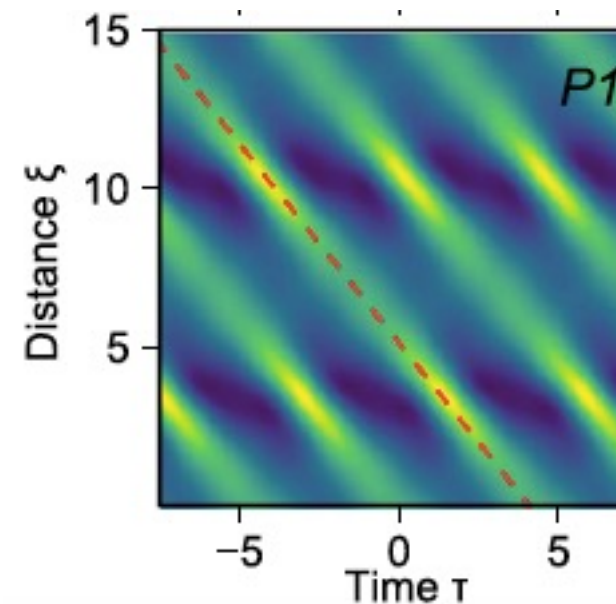
Simplified model: three-wave truncation



Outside the separatrix:
P2 orbits



Separatrix: Akhmediev Breather



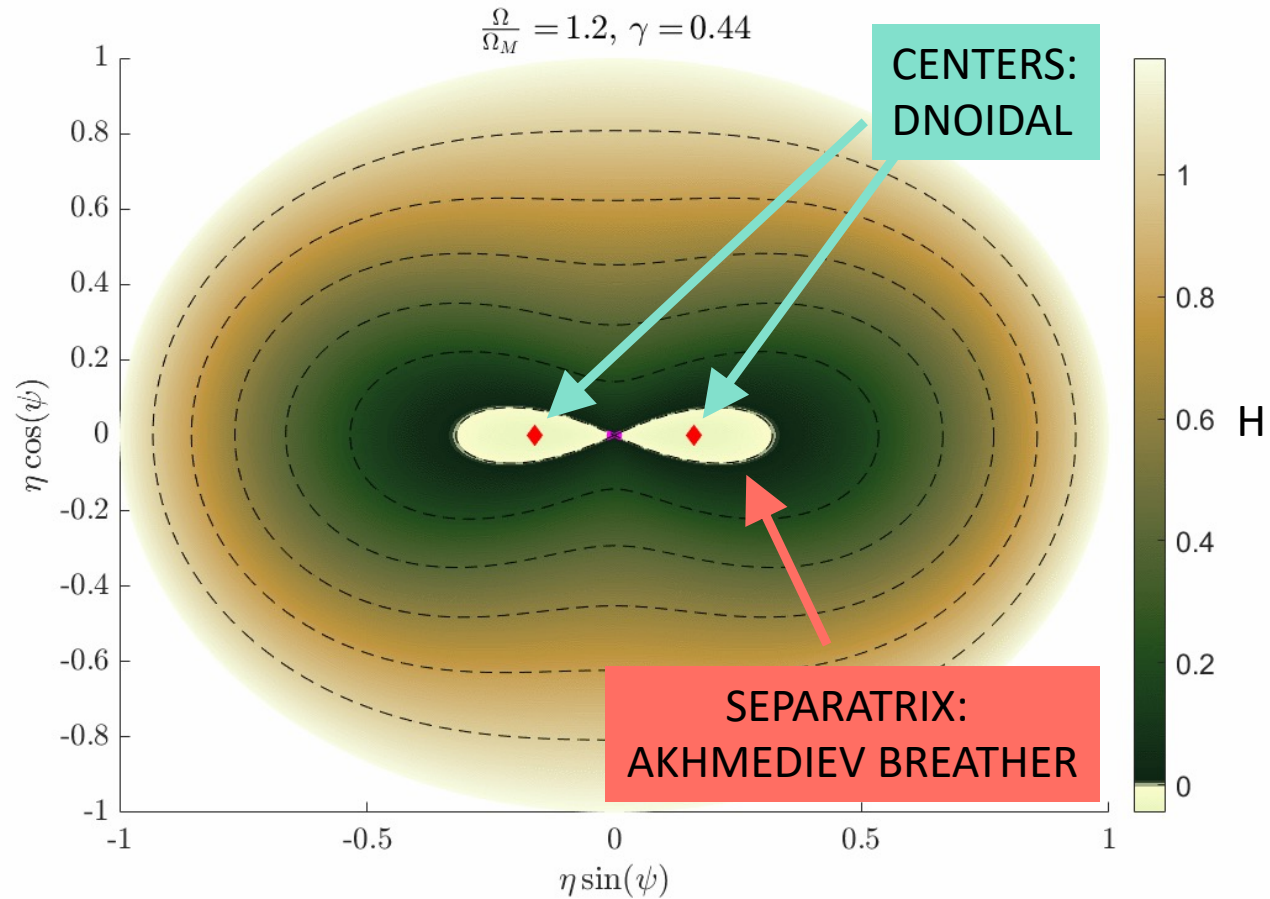
Inside the separatrix: P1 orbits

Trillo & Wabnitz Opt. Lett. (1991)

Armaroli, Brunetti, Kasparian PRE (2017)

Eeltink et al. Nonlin. Dyn. (2020)

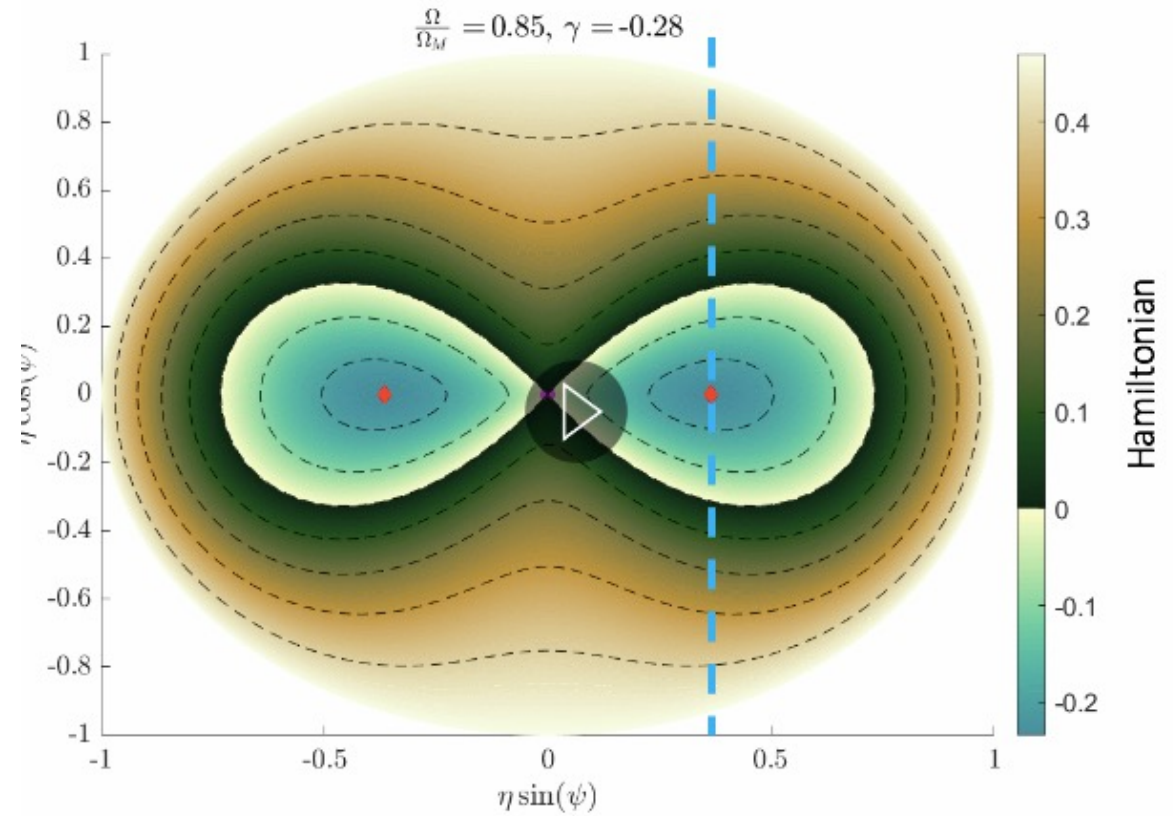
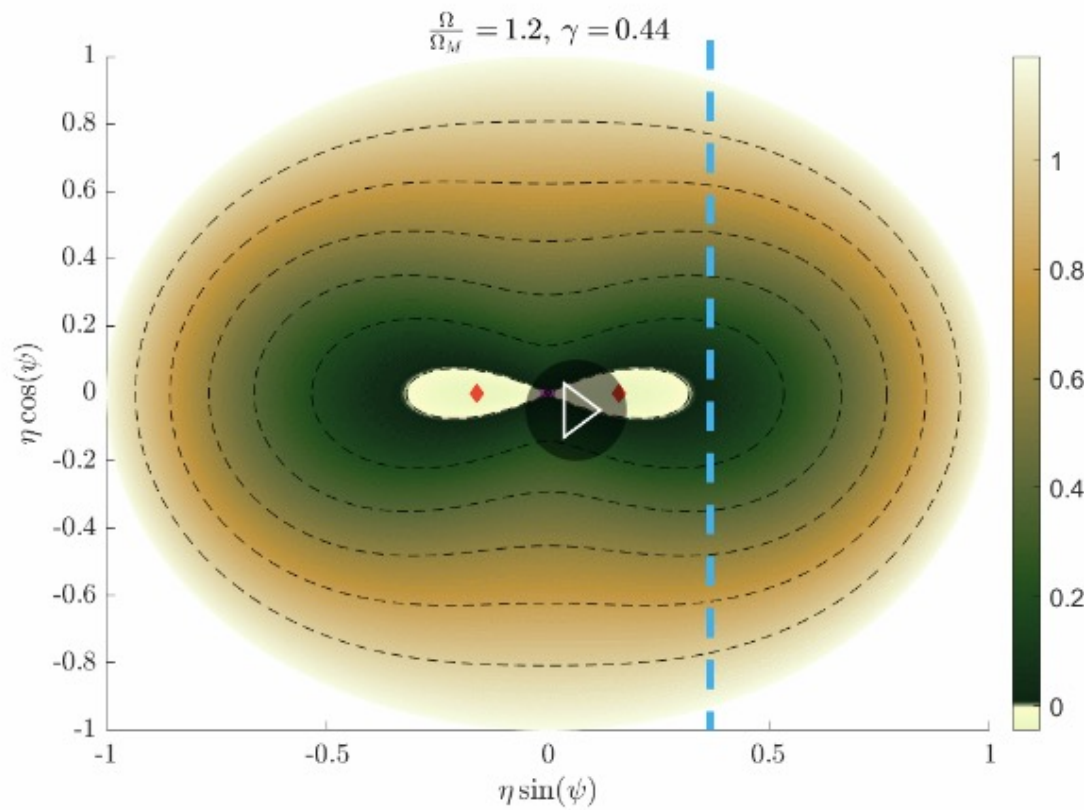
Simplified model: three-wave truncation



Trillo & Wabnitz *Opt. Lett.* (1991)
Armaroli, Brunetti, Kasparian *PRE* (2017)



Phase space evolution

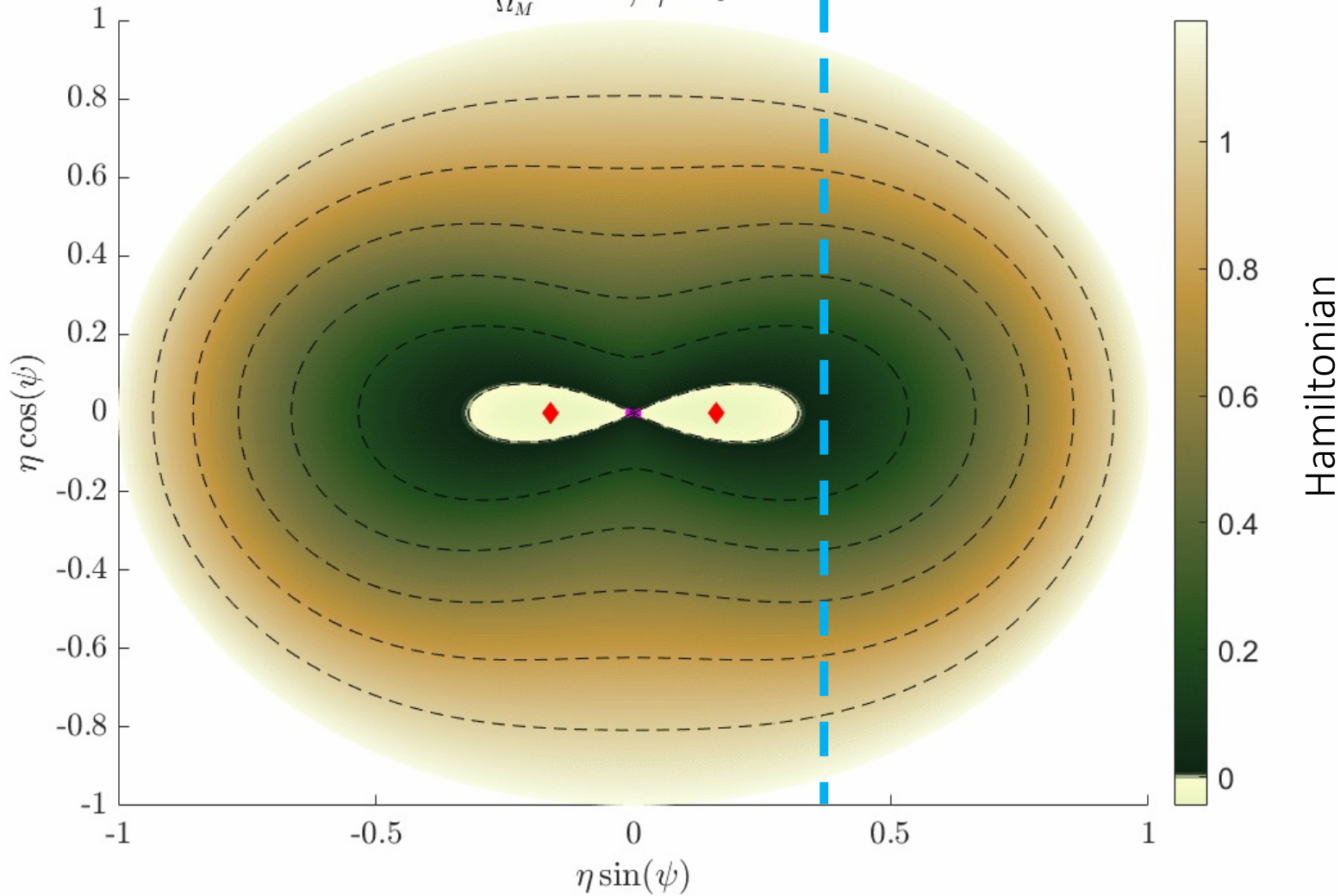


Matching condition between AB and dnoidal detuning: $\Omega_{dn} = 2(\Omega_{AB} - 1)$



Phase space evolution

$$\frac{\Omega}{\Omega_M} = 1.2, \gamma = 0.44$$



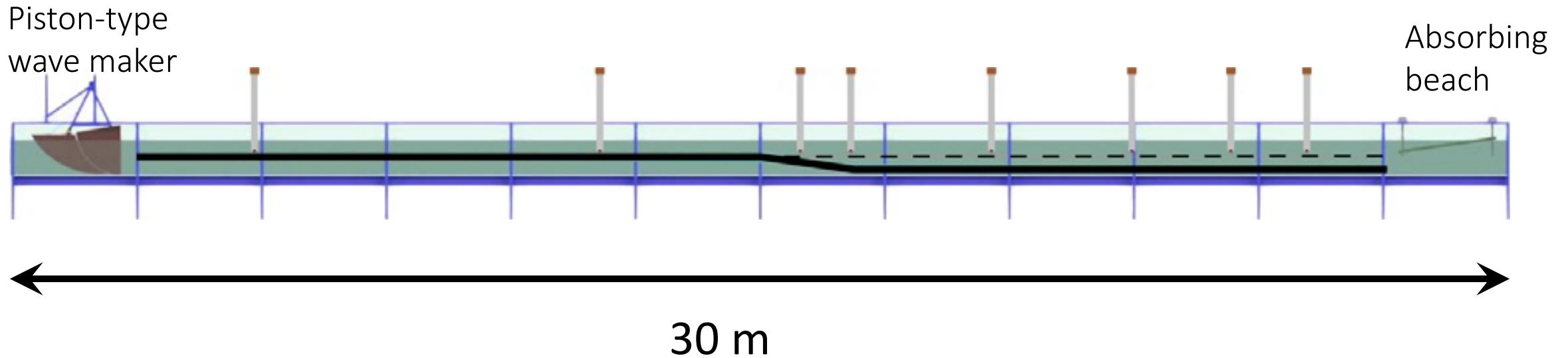
$$\gamma = - \left[\frac{\alpha \Omega^2}{\tilde{\beta} E} + 1 \right] = \left(\frac{\Omega}{\Omega_M} \right)^2 - 1$$

Matching condition between AB and dnoidal detuning:

$$\Omega_{dn} = 2(\Omega_{AB} - 1)$$



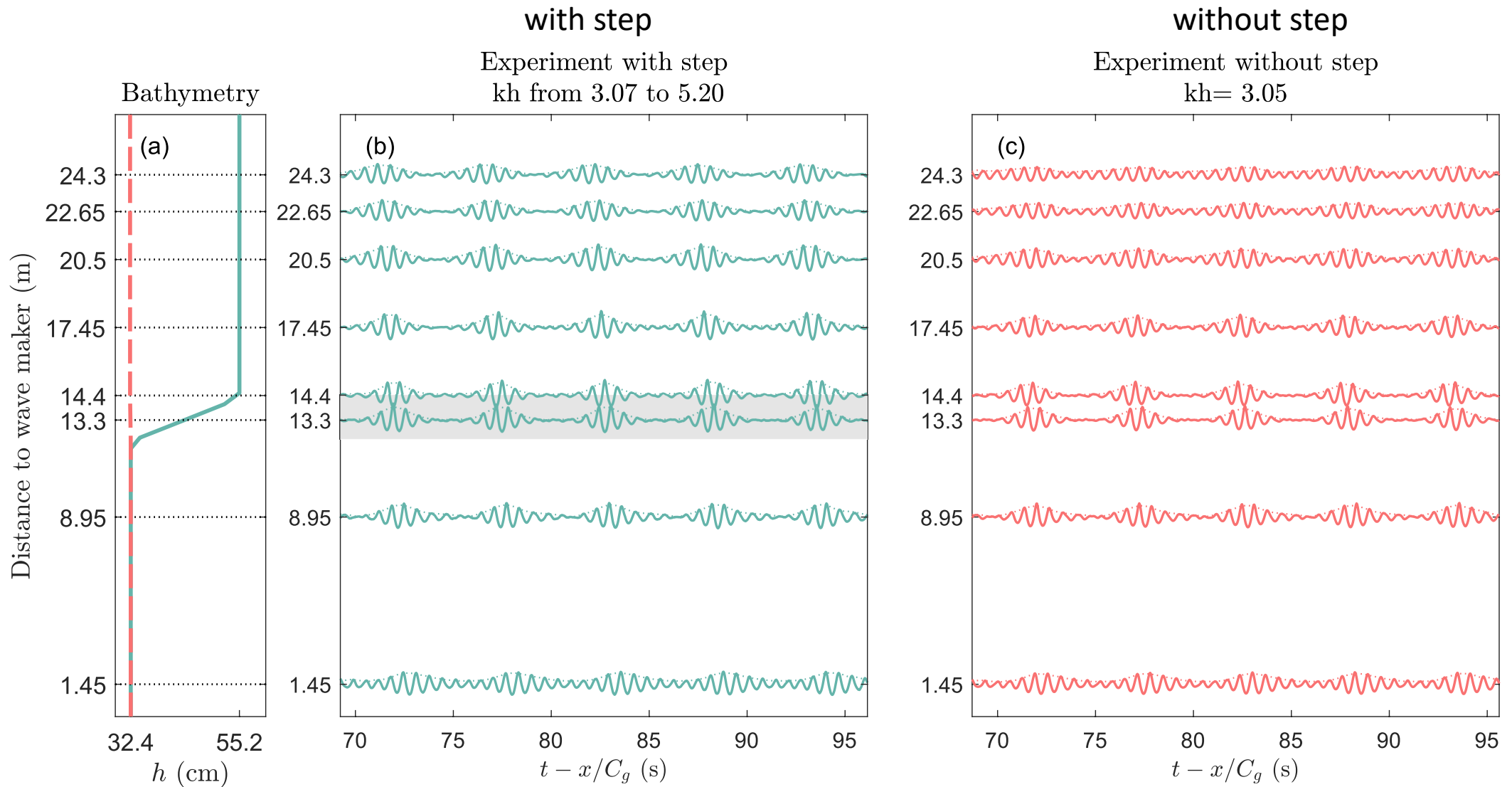
Experiments in the water tank at The University of Sydney by Alexis Gomel & Amin Chabchoub



- Water wave flume with artificial floor setup
- 8 wave gauges
- Water depth varied from $h = 32.4$ cm to 55.2 cm
- $\Delta h \ll L_{\text{step}} \ll L_{\text{nonlinear}}$ to prevent spurious reflections

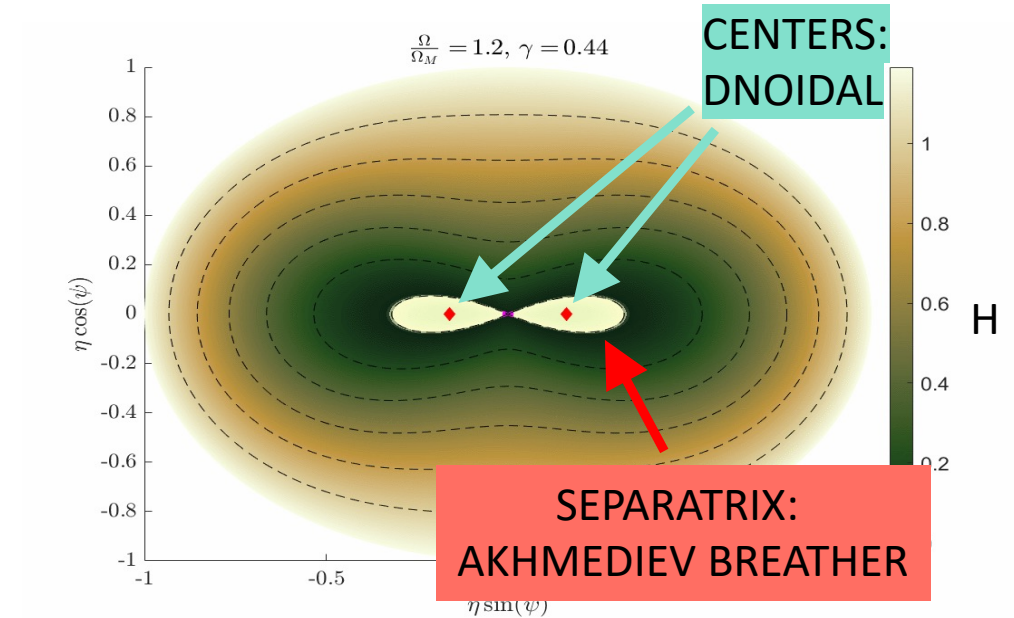
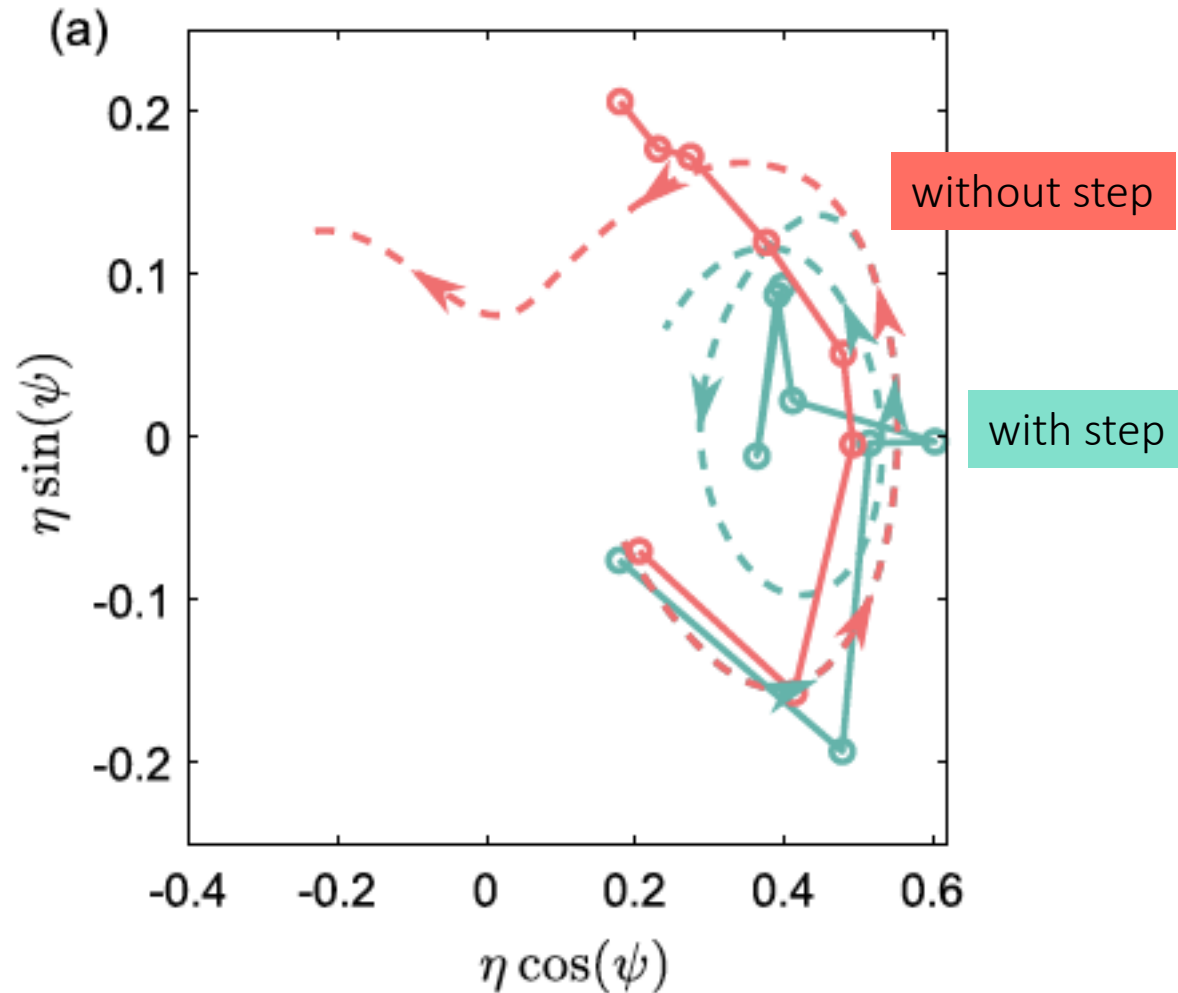


Experiments in the water tank



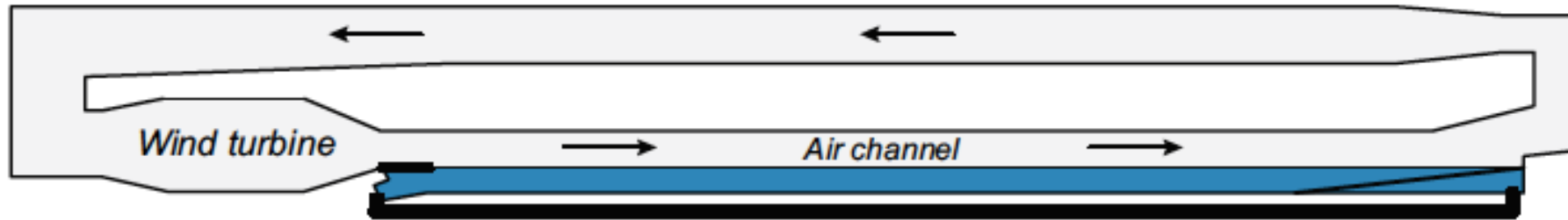
Experiments in Sydney vs numerical simulations

Evolution in 3-wave phase space

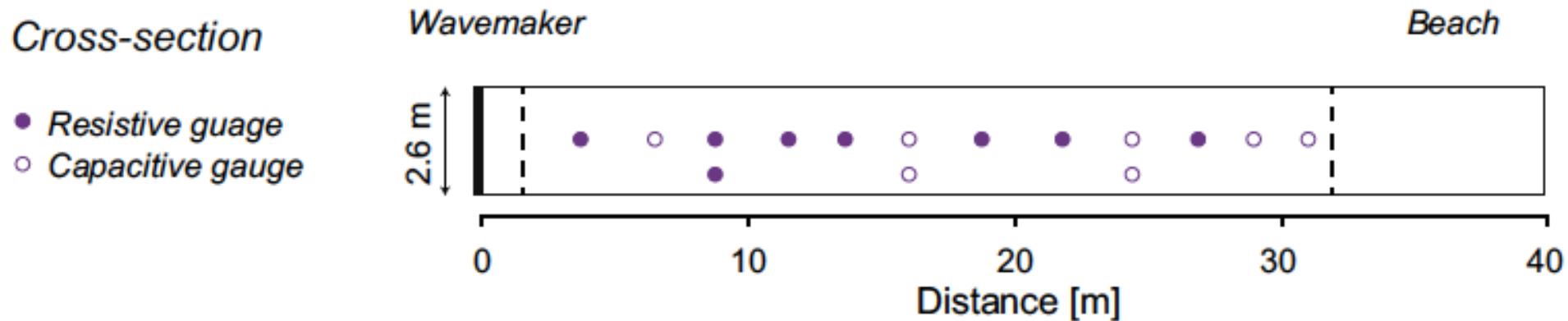


Experiments in the wind-wave facility at
IRPHE/PYTHEAS (Luminy) Aix Marseille University
by Debbie Eeltink, Huber Branger & Christopher Luneau

Longitudinal-section



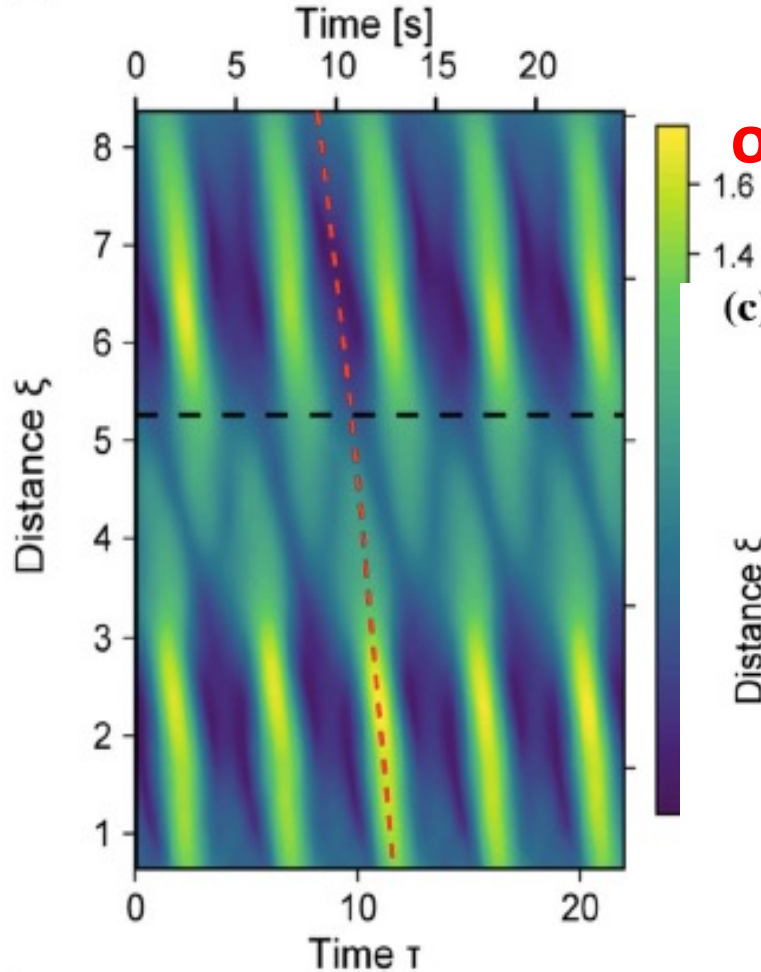
Cross-section



Wind effect: separatrix crossing

(a)

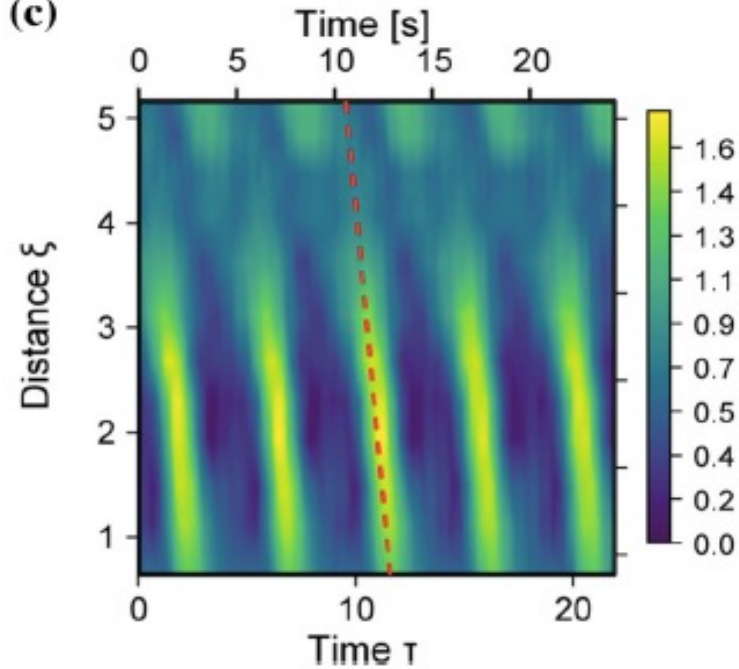
No wind



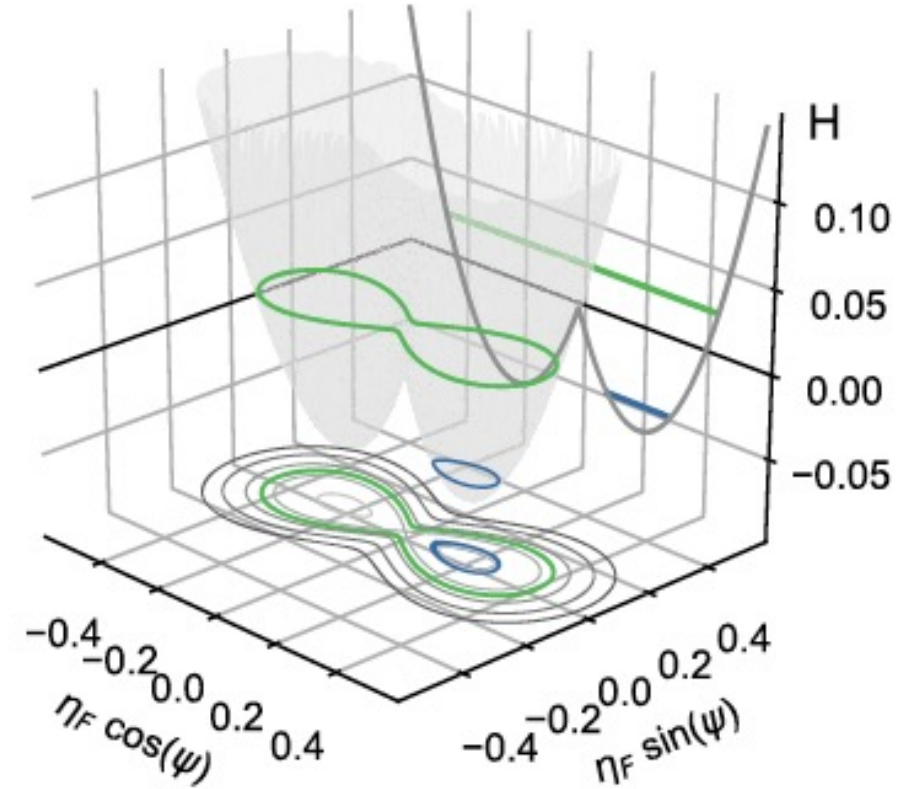
SIMULATION

WIND OFF:
outside the separatrix
P2 orbit

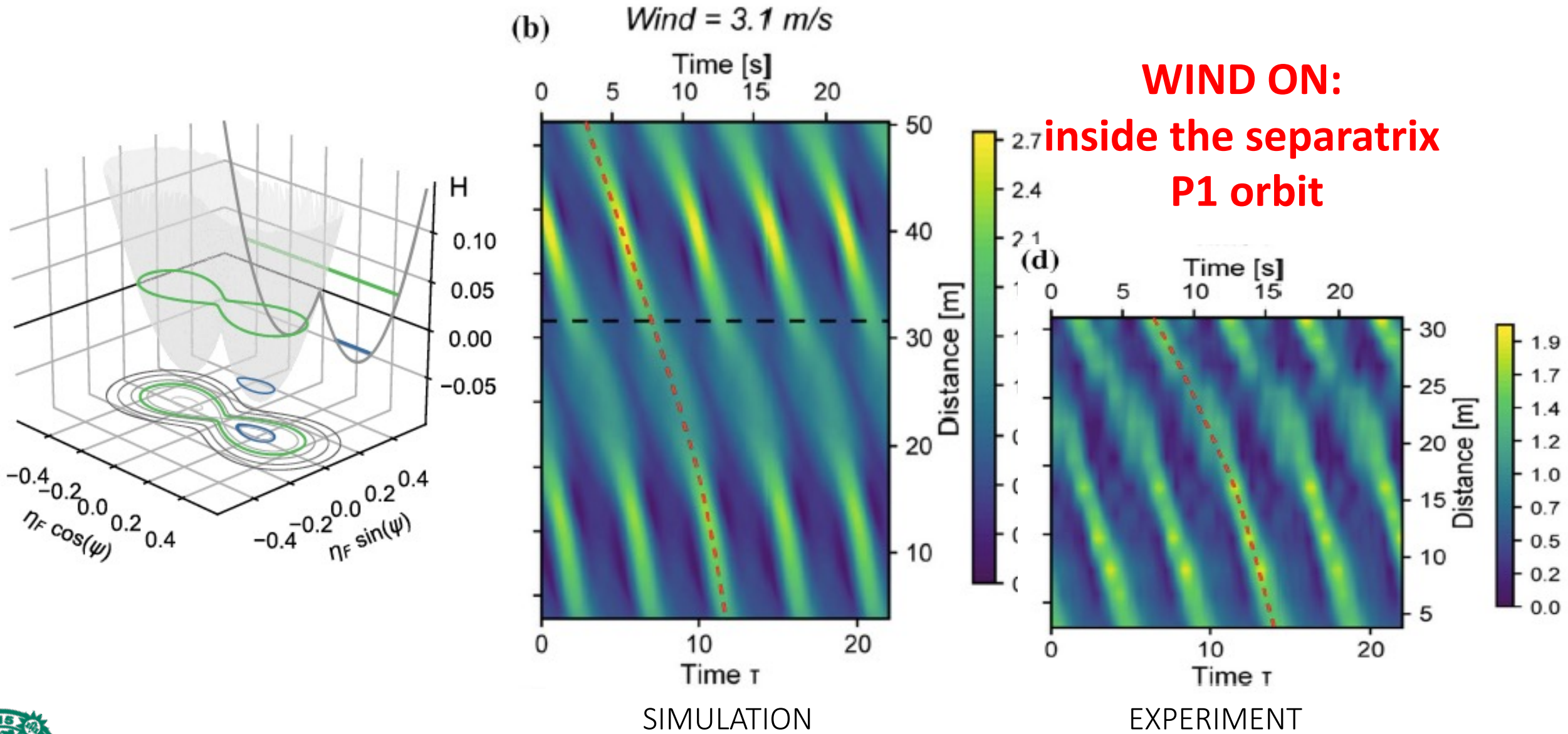
(c)



EXPERIMENT



Wind effect: separatrix crossing

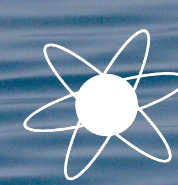


Conclusions

1. **Phase space manipulation:**
bathymetry is changed to obtain a shift toward a stable state
2. The experiment is described sufficiently well by the **NLS framework:**
 - dissipation & high-order terms are small in one recurrence cycle
 - reflections can be controlled by choosing a mild slope
$$\Delta h \ll L_{\text{step}} \ll L_{\text{nonlinear}}$$
3. **Separatrix crossing:** example of wind forcing
4. Main physical processes are described by **NLS** and **three-wave approximation:**
integrable equation, construction of whole families of solutions,
many possibilities for matching
5. The same procedure can be applied to other systems
(for example KdV, see Binder, Vanden-Broeck & Dias, Chaos (2005))



Thanks!



Faculté
des Sciences
Aix*Marseille Université

Gomel, Chabchoub, Brunetti, Trillo, Kasparian, Armaroli
Stabilization of Unsteady Nonlinear Waves by Phase-Space Manipulation
Physical Review Letters 126, 174501 (2021)

Armaroli, Gomel, Chabchoub, Brunetti, Kasparian
Stabilization of uni-directional water wave trains over an uneven bottom
Nonlinear Dynamics 101, 1131-1145 (2021)

Eeltink, Armaroli, Luneau, Branger, Brunetti, Kasparian
Separatrix crossing and symmetry breaking in NLSE-like systems due to forcing and damping
Nonlinear Dynamics 102, 2385-2398 (2020)