

# Improving the numerical description of high-pressure solid hydrogen phases

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# Introduction

- Under pressures of the megabar regime, a variety of hydrogen solid phases exist [1]
- Phase boundaries have been determined experimentally up to 300 GPa using Infrared and Raman spectroscopies [2]
- Structural information on these phases is however missing due to X-ray diffraction experiments being inconclusive on these systems [3]
- Growing interest of the scientific community for high-pressure solid hydrogen phases because they are potential high-temperature superconductors [1]

#### For this study:

• Use of the *Quantum Espresso* plane waves DFT code

• Use of the *TurboRVB* quantum Monte Carlo code

- Pressure range of interest: 100-350 GPa • Focus on the II-III phase transition
- Families of methods considered: Hartree-Fock (HF), Generalized Gradient Approximation (GGA) and hybrid Density Functional Theory (DFT), Random Phase Approximation (RPA), Coupled Cluster (CC), Diffusion Monte Carlo (DMC)
- Four different molecular solid phases of interest: C2/c-24 (monoclinic), P2<sub>1</sub>/c-24 (monoclinic), P6<sub>3</sub>/m-16 (hexagonal), Cmca-12 (orthorhombic)







**Motivations** 



- In this work, diverse methods are tested on various high-pressure solid hydrogen phases and compared to CCSD [4] and DMC [5]
- Apparent agreement of vdW-DF and PBEO with experimental II-III phase transition  $\succ$
- But more advanced CCSD and DMC methods predict the same phase transition to occur at much higher pressures

The large difference between CCSD and DMC correlated methods shows the

## **RPA and RPAx results**

- RPA and RPAx methods are constructed within the DFT framework and rely on the Adiabatic Connection Fluctuation Dissipation Theorem (ACFDT) [6]
- Both methods are based on the exact expression of the correlation energy:  $E_{c} = \frac{i}{2} \int_{0}^{1} d\lambda \int \frac{d\omega}{2\pi} Tr\{v(\boldsymbol{r} - \boldsymbol{r}')[\chi_{\lambda}(\boldsymbol{r}, \boldsymbol{r}', \omega) - \chi_{s}(\boldsymbol{r}, \boldsymbol{r}', \omega)]\}$
- Like with Time-Dependent DFT, the density response function  $\chi$  is needed [7]: •  $\chi_{\lambda} = \chi_{s} + \chi_{s} [\lambda \nu + f_{xc}^{\lambda}] \chi_{\lambda} = \chi_{s} + \chi_{s} \left[ \lambda \nu + \frac{\partial \nu_{xc}^{\lambda}}{\partial n} \right] \chi_{\lambda}$  $\circ \quad \underline{\text{For RPA}}: \quad f_{xc}^{\lambda} = 0 \qquad \circ \quad \underline{\text{For RPAx}}: \quad f_{xc}^{\lambda} = f_{x}^{\lambda}$
- The total RPA energy is notably composed of the RPA correlation energy and the exchange energy, determined using the EXact Exchange method (EXX):

$$E_{tot}^{RPA} = T_s + V_{ext} + E_H + E_x^{EXX} + E_c^{RPA}$$

On diverse systems, RPA and RPAx methods have shown a high level of accuracy for moderate computational costs [7,8]

• Not possible to relax structures using RPA or RPAx in *Quantum Espresso* yet: RPA and RPAx calculations done on PBE and vdW-DF geometries instead



- RPA@PBE and RPAx@PBE provide similar results on all phases of interest
- RPA@vdW-DF is in very good agreement with the DMC results we obtained using the *TurboRVB code* on the same set of vdW-DF geometries

The RPA method presents an excellent accuracy-cost ratio on high-pressure solid hydrogen phases

### **Hybrid PBE0 results**

- Standard PBEO hybrid functional admits  $\alpha$ =25% of exact exchange
- Difference in results between PBE025 and HF confirms the importance of optimizing the fraction of exact exchange for PBEO on these systems: done by calculating the RPA and RPAx energy on top of the equilibrium geometry and orbitals predicted by PBEO $\alpha$  on a single H<sub>2</sub> molecule
- Minimal energy obtained with 41% of exact exchange for RPA, 48% for RPAx

RPAx being a more accurate method than RPA, calculations will be performed with the PBE048 (=HYB<sub>ont</sub>) functional instead of PBE041





- Contrary to the PBE025 method, PBE048 hybrid calculations performed on vdW-DF geometries deliver improved results, in satisfying agreement with RPA
- The good agreement between RPA and HYB<sub>opt</sub> confirms that the exact exchange proportion to use for PBEO should be much larger than 25%
- The 48% value found on a single  $H_2$  molecule also appears to be  $\succ$ the suited proportion to study pure extended hydrogen systems

Interesting performance of HYB<sub>ont</sub> given its accuracy and it being much more affordable than RPA, RPAx and DMC in terms of computation costs

## **Role of the starting geometry**

- Despite the significant lack of accuracy, vdW-DF returns better enthalpy curves than PBE: vdW-DF geometries should be closer to the true geometries than PBE geometries do
- Performing hybrid DFT or RPA/RPAx calculations on a different set of geometries will yield different enthalpy curves
- The change of PBE geometries to vdW-DF geometries leads to a shift in pressures of about 20 GPa: impact rather small, the geometry is not the main source of error in the calculation of enthalpy curves

But the geometry still has a major impact: the use of better HYB<sub>opt</sub> geometries leads HYB<sub>ont</sub>, RPA and DMC to predict a phase never considered before, P-1-24, to be a promising candidate for the phase II structure [9]







#### P-1-24

*Triclinic phase close to C2/c-24* except that part of the H<sub>2</sub> molecules are rotated out of the (a,c)-plane

# The II-III phase transition

- Without the consideration of vibrations in the calculations, HYB<sub>ont</sub>, RPA and DMC find that phase III corresponds to C2/c-24 while phase II corresponds to  $P2_1/c-24$  [9]
- When vibrations are considered, with the improved HYB<sub>ont</sub> geometries, the phase P-1-24 is predicted to become more stable than any of the other structures suggested [9]
- The calculated Raman spectrum shows that the evolution of the frequency of the vibron agrees with experiments [10]. This is true only when moving from C2/c-24 to P-1-24. The shift is much larger when transitioning from C2/c-24 to P2<sub>1</sub>/c-24 [9]
- The decrease in the frequency of the vibron when transitioning from C2/c-24 to P-1-24 can be explained by the strengthening of the intramolecular bond of the H<sub>2</sub> molecules that are rotated out-of-plane, since they lose their inplane stability, leading to a shortening of their H-H bond distance [11]



### **Summary**

PBE048, RPA and RPAx methods offer satisfying accuracy-cost ratios on high-pressure solid hydrogen phases [9]

C2/c-24

Monoclinic phase where  $H_2$ 

molecules are all arranged

within the (a,c)-plane

- All three methods predict phase III to consist of the C2c-24 phase, while the prediction for phase II varies from P2<sub>1</sub>/c-24 and P-1-24, depending on the additional effects being considered during the calculations
- Along with DMC, they predict the II-III phase transition to occur close to 250 GPa, almost 100 GPa higher than experimental results [1]
- The large discrepancy noted with experiments is likely due to numerous factors: incorrect geometries, non-consideration of vibrations and quantum nuclear effects in the calculations,... [9,12]
- Future works should focus on evaluating the impact of each factor on the numerical results obtained so far

#### References

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