

LHCb m_W measurement with 2016 data

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on behalf of the LHCb Collaboration

W mass workshop, Paris, 23/02/23 - 24/02/23

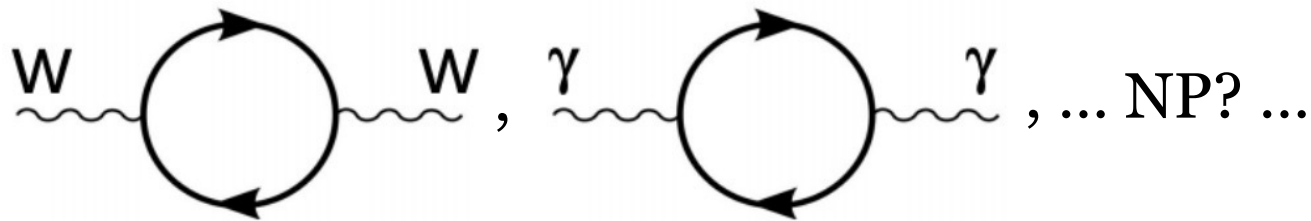


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Scientific Context

In the Standard Model:

$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2} \right) = \frac{\pi\alpha}{\sqrt{2}G_F} (1 + \Delta)$$



- Can indirectly predict m_W in global EW fits with inputs from rest of the SM parameters.
- Comparing with direct m_W measurements constrains new physics.

Status of the field

- 2021 global EW fit:

$$\Delta m_W^{EW fit} = 6 \text{ MeV},$$

- Recent direct measurements:

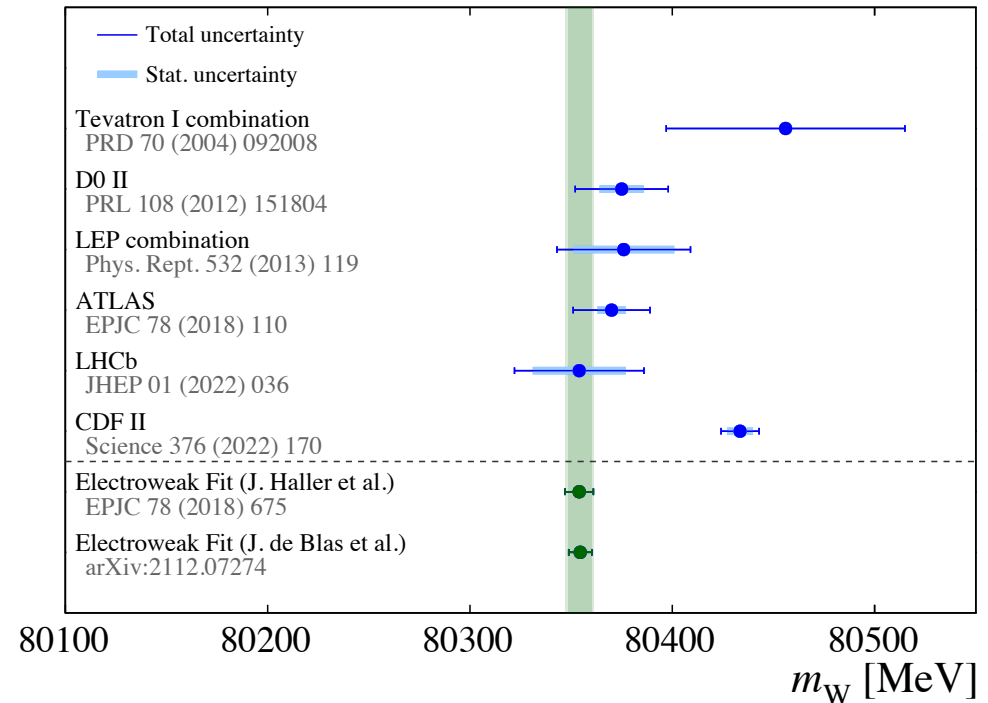
$$\Delta m_W^{ATLAS} = 19 \text{ MeV},$$

$$\Delta m_W^{CDF II} = 9 \text{ MeV},$$

$$m_W^{CDF II} - m_W^{EW fit} \approx 7\sigma,$$

$$m_W^{CDF II} - m_W \text{ (LEP, D0 + LHC)} \approx 4\sigma.$$

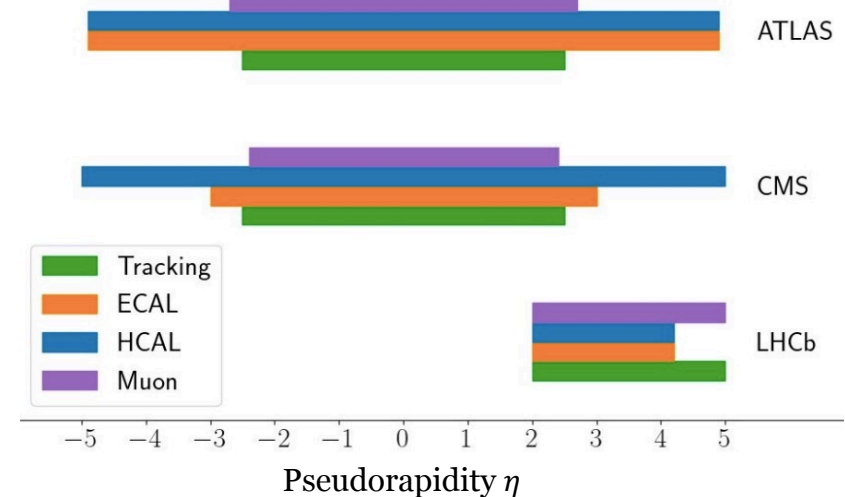
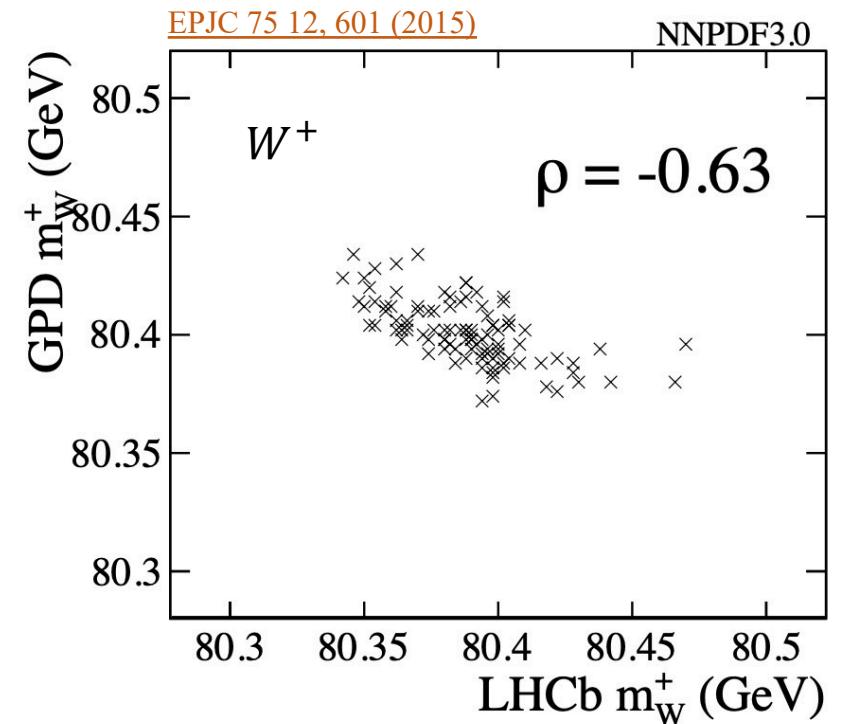
LHCb-FIGURE-2022-003



More m_W measurements from the LHC are necessary!

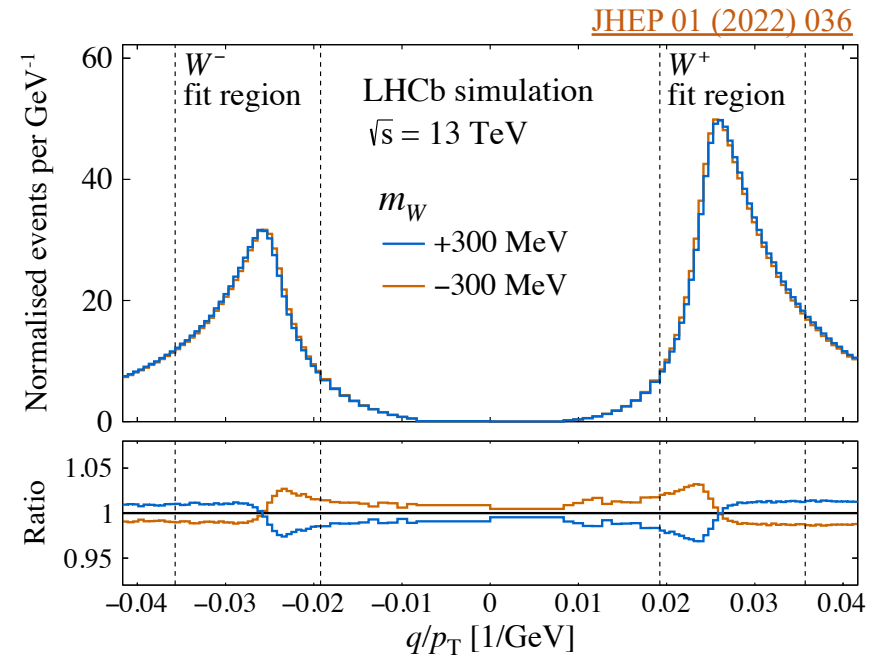
Why LHCb?

- LHCb Run-2 data: $O(10)$ MeV statistical uncertainty on m_W ($O(10^7)$ $W \rightarrow \mu\nu$ candidates),
- Historically-limiting PDF uncertainties expected to anti-correlate in a GPD-LHCb combination.
- Mostly designed for flavour physics, but with a strong programme of probing vector boson production.
 - Full list of LHCb EW papers [here](#).



How we measure m_W

- $W \rightarrow \mu\nu$ gives a single, high- p_T , isolated muon (LHCb doesn't reconstruct missing energy).
- m_W sensitivity from p_T^μ , which peaks at $\sim m_W/2$, therefore we extract m_W in a **template** fit to the muon q/p_T distribution.
- Need supreme understanding of important factors that affect the p_T^μ shape.



“Experimental” modelling

e.g. muon momentum scale & calibration, detector misalignment, reconstruction & selection efficiencies etc.

“Theoretical” modelling

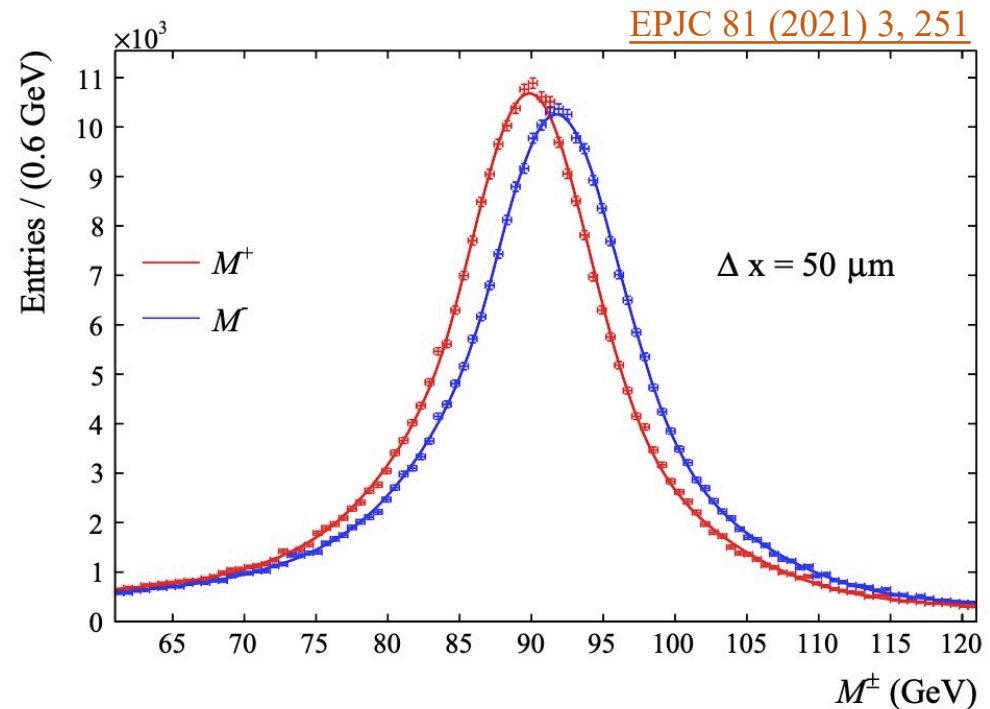
e.g. W cross-section predictions (unpolarised and angular distribution), QED FSR, PDFs etc.

Experimental modelling

Detector alignment corrections

- Biases in p_T^μ can originate from detector misalignments. Fix with:
 - 1) Custom alignment for high- p_T muons,
 - 2) Finer, analysis-level curvature (q/p) corrections from the “pseudomass” method on $Z \rightarrow \mu\mu$.
- Differences in M^+ and M^- allow for mapped curvature bias corrections across the detector.

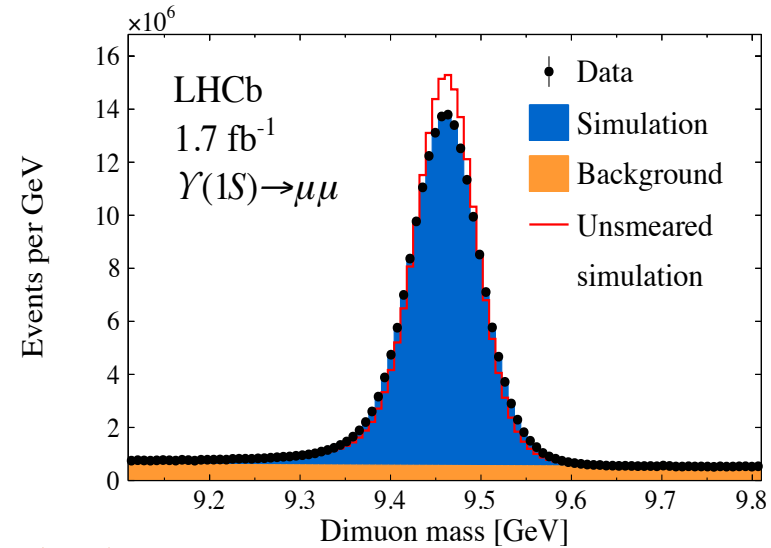
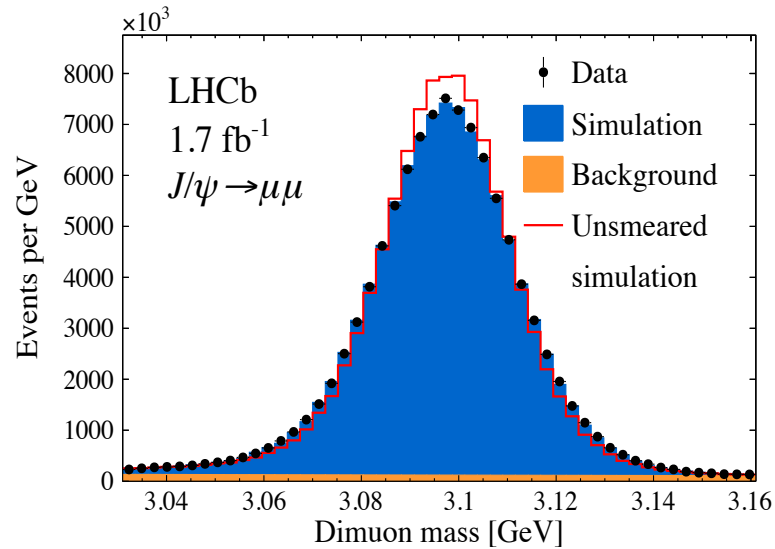
$$M^\pm = \sqrt{2p^\pm p_T^\pm \frac{p^\mp}{p_T^\mp} (1 - \cos \theta)},$$



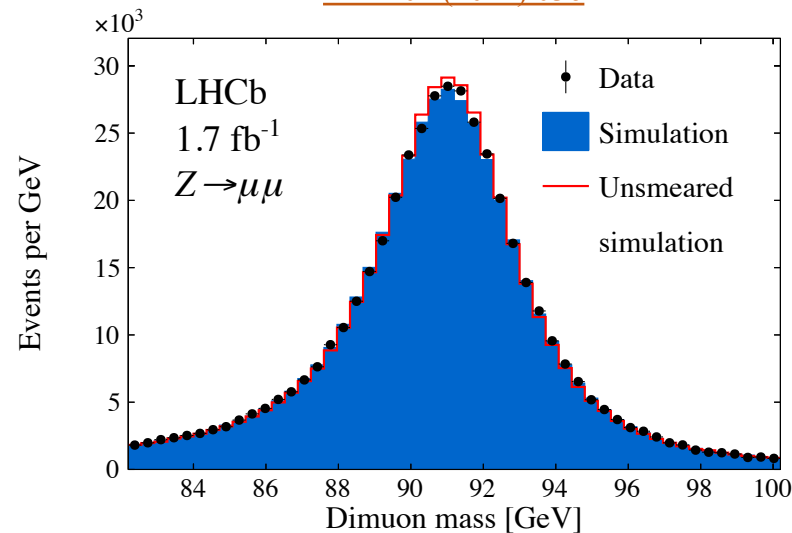
Inspired by [PRD 91, 072002 \(2015\)](#)

Momentum scale calibration

Further smearing of the simulation is then needed:



[JHEP 01 \(2022\) 036](#)



Simultaneous fit to the
 $J/\psi, \Upsilon(1S)$ and $Z m_{\mu\mu}$

(binned in different detector regions,
polarities...)

Smearing model applied to
all templates.

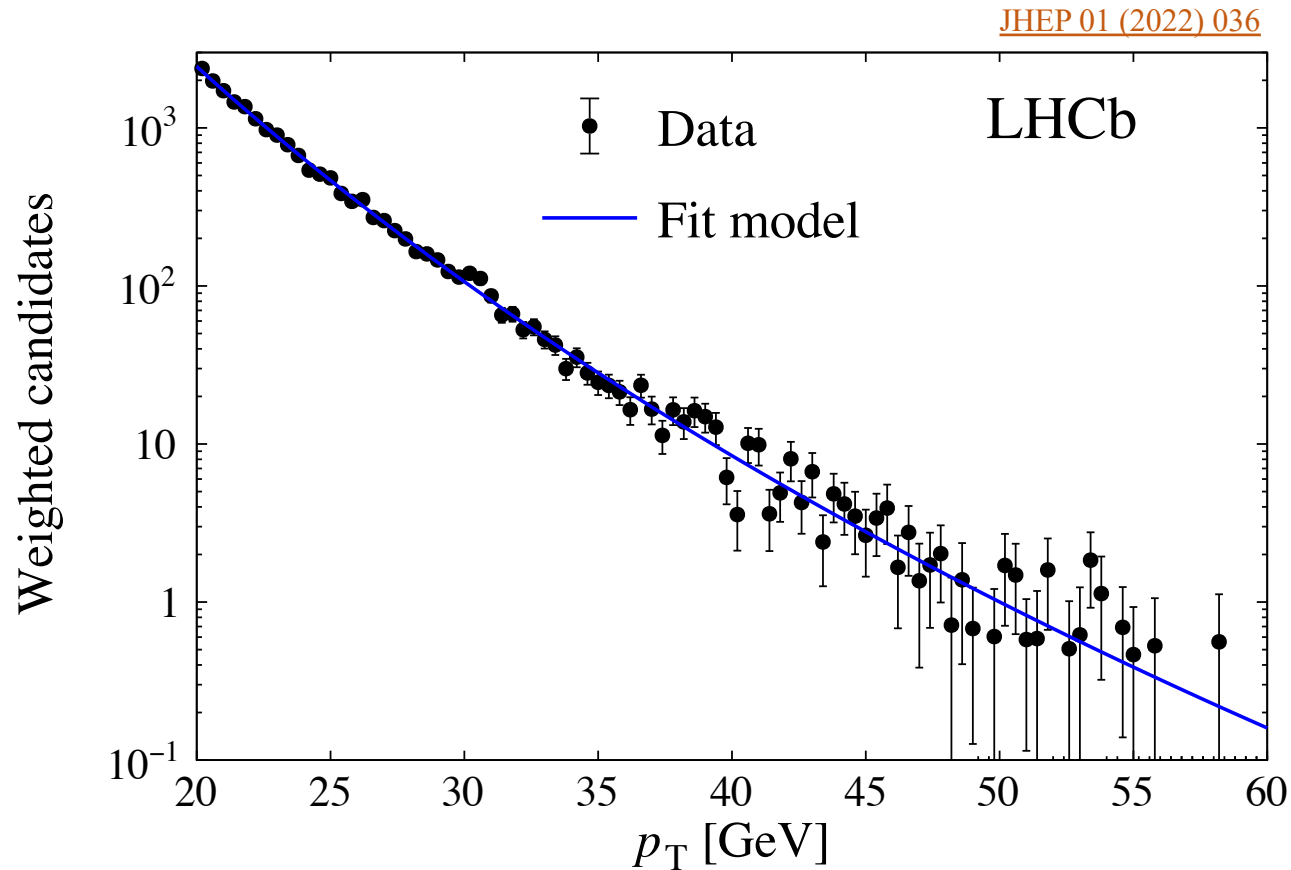
Signal selection

- Veto events with second high- p_T^μ muon in acceptance ($p_T^\mu > 25$ GeV); rejects $Z \rightarrow \mu\mu$,
- Signal muon is well-reconstructed, muon ID-ed and required to fire high- p_T single muon triggers,
- Muon candidate is isolated; rejects heavy flavour & decay-in-flight backgrounds.

This selects 2.4M events in the fit window $28 < p_T^\mu < 52$ GeV, $2.2 < \eta < 4.4$.

Treatment of backgrounds

- Electroweak backgrounds constrained with $Z \rightarrow \mu\mu$.
- Remaining decay-in-flight hadronic background (10x heavy flavour) modelled with a parametric shape, trained on a hadron-enriched data sample:



Theoretical modelling

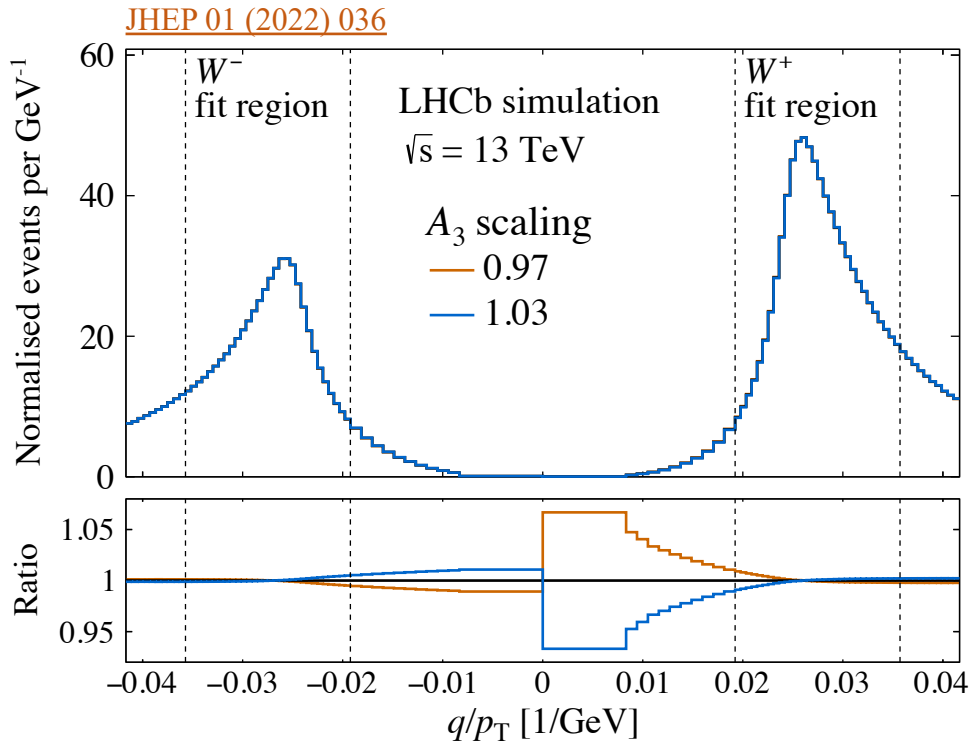
Angular coefficients

At the Born level
(before QED FSR):

$$\frac{d\sigma}{dp_T^W dy dM d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{\text{unpol}}}{dp_T^V dy dM} \left. \vphantom{\frac{d\sigma}{dp_T^W dy dM d\cos\theta d\phi}} \right\} \text{Unpolarised cross-section}$$

$$\left. \vphantom{\frac{d\sigma}{dp_T^W dy dM d\cos\theta d\phi}} \right\} \text{Angular terms } (A_i = \text{angular coefficients})$$

$$\left\{ \begin{aligned} & (1 + \cos^2\theta) + A_0 \frac{1}{2} (1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \\ & + A_2 \frac{1}{2} \sin^2\theta \cos 2\phi + \boxed{A_3 \sin\theta \cos\phi} + A_4 \cos\theta \\ & + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \end{aligned} \right\}$$

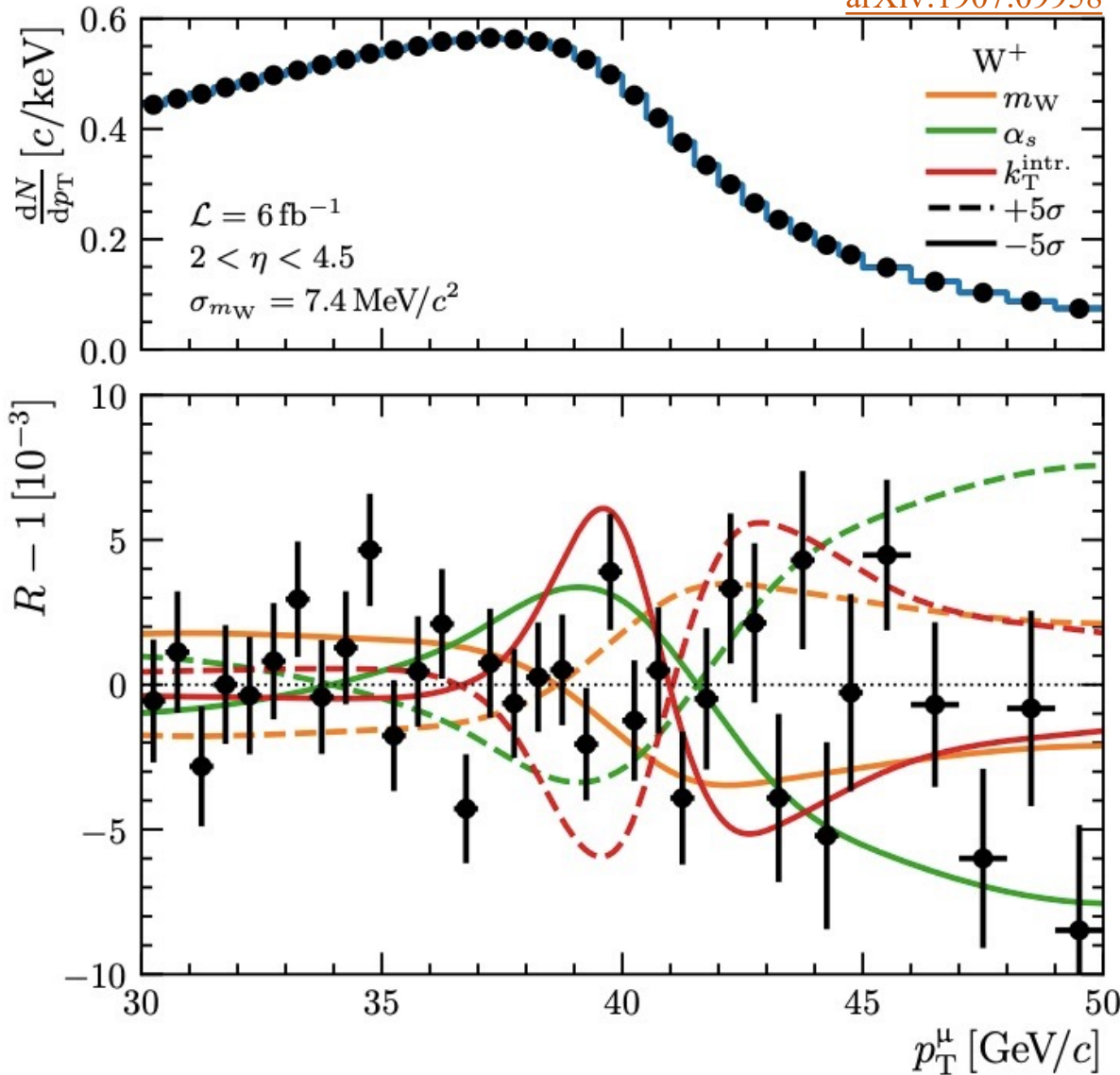


2016 measurement strategy:

- A_i predictions from DYTurbo at $\mathcal{O}(\alpha_S^2)$.
- Floated a scale factor in the fit to absorb the (dominating) uncertainty on the A_3 prediction.
- Conservative uncertainty treatment (from [JHEP 11\(2017\) 003](#)) with uncorrelated scale variations $\rightarrow 10 \text{ MeV}$.

Physics modelling: σ^{unpol}

arXiv:1907.09958

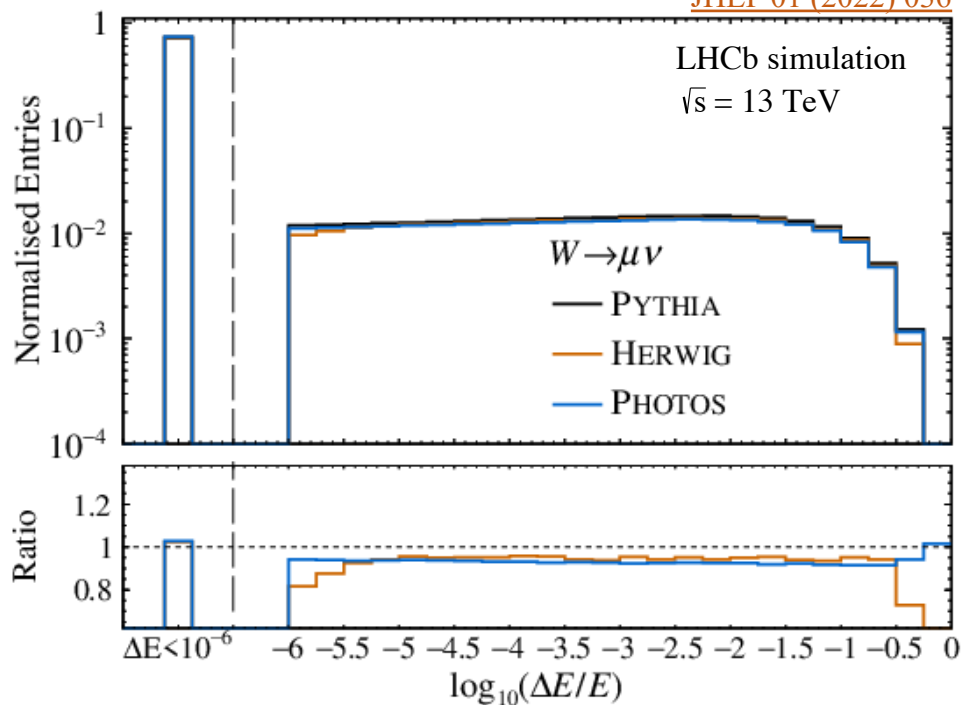
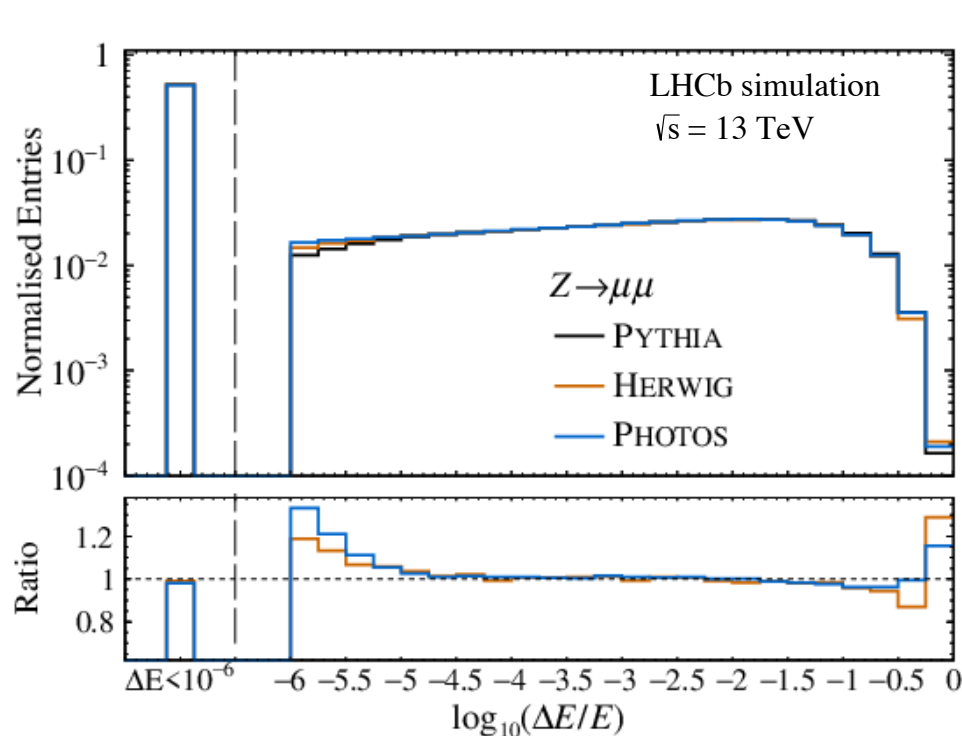


$$\left. \frac{d\sigma^{unpol}}{dp_T^V dy dM} \right\} \text{Unpolarised cross-section}$$

- POWHEG-Box + Pythia8 was our central model.
 - Previous m_W measurements rely on tuning to p_T^Z . Does this tune hold for p_T^W ?
 - Variations in α_s and $k_T^{intr.}$ affect p_T^μ differently to variations in m_W .
- ⇒ Floated these QCD parameters in a simultaneous fit to $W q/p_T^\mu$ and $Z \phi^*$.

QED Final State Radiation

- Different FSR predictions mimicked by reweighting in $\Delta E/E$,
- No preference between predictions from Pythia, Herwig and Photos,



- Uncertainty from the envelope of fits with each.

PDFs and their uncertainties

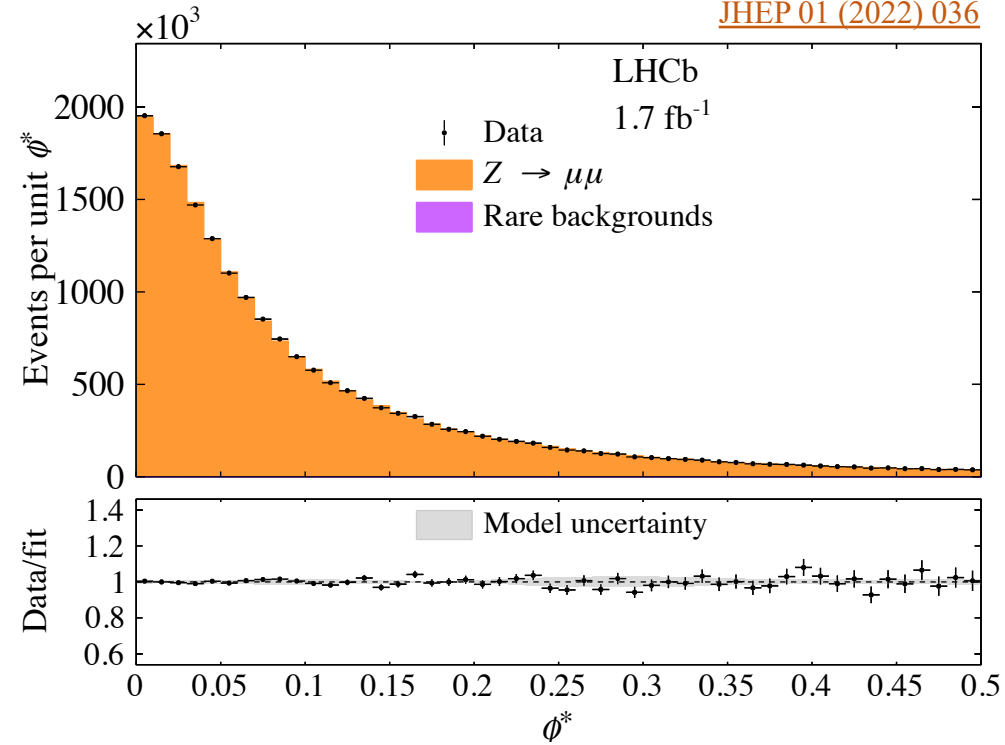
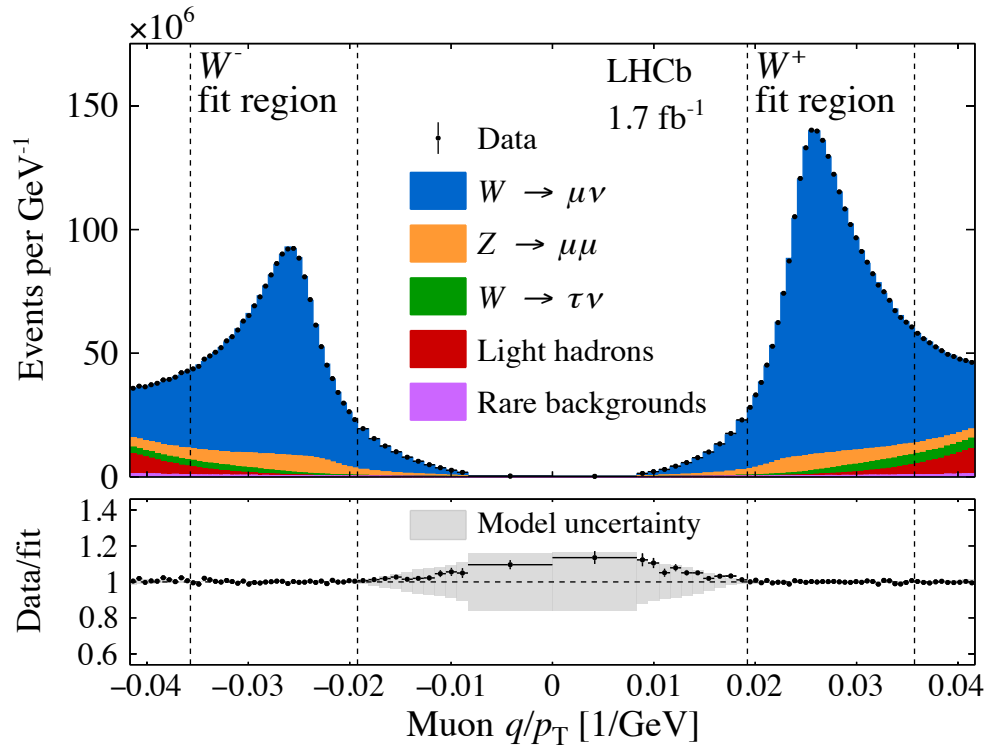
- Treated PDFs from [NNPDF3.1](#), [CT18](#) and [MSHT20](#) equally, and their uncertainties as fully-correlated:

$$m_W = \frac{1}{3} [m_W(\text{NNPDF}) + m_W(\text{CTEQ}) + m_W(\text{MSHT})],$$
$$\Delta m_W(\text{PDF}) = \frac{1}{3} [\Delta m_W(\text{NNPDF}) + \Delta m_W(\text{CTEQ}) + \Delta m_W(\text{MSHT})].$$

Set	$\sigma_{\text{PDF,base}}$ [MeV]	$\sigma_{\text{PDF},\alpha_s}$ [MeV]	σ_{PDF} [MeV]
NNPDF3.1	8.3	2.4	8.6
CT18	11.5	1.4	11.6
MSHT20	6.5	2.1	6.8

The 2016 fit result

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Parameter	Value
Fraction of $W^+ \rightarrow \mu^+\nu$	0.5288 ± 0.0006
Fraction of $W^- \rightarrow \mu^-\nu$	0.3508 ± 0.0005
Fraction of hadron background	0.0146 ± 0.0007
α_s^Z	0.1243 ± 0.0004
α_s^W	0.1263 ± 0.0003
k_T^{intr}	$1.57 \pm 0.14 \text{ GeV}$
A_3 scaling	0.975 ± 0.026

$$\chi^2/ndf = 105/102$$

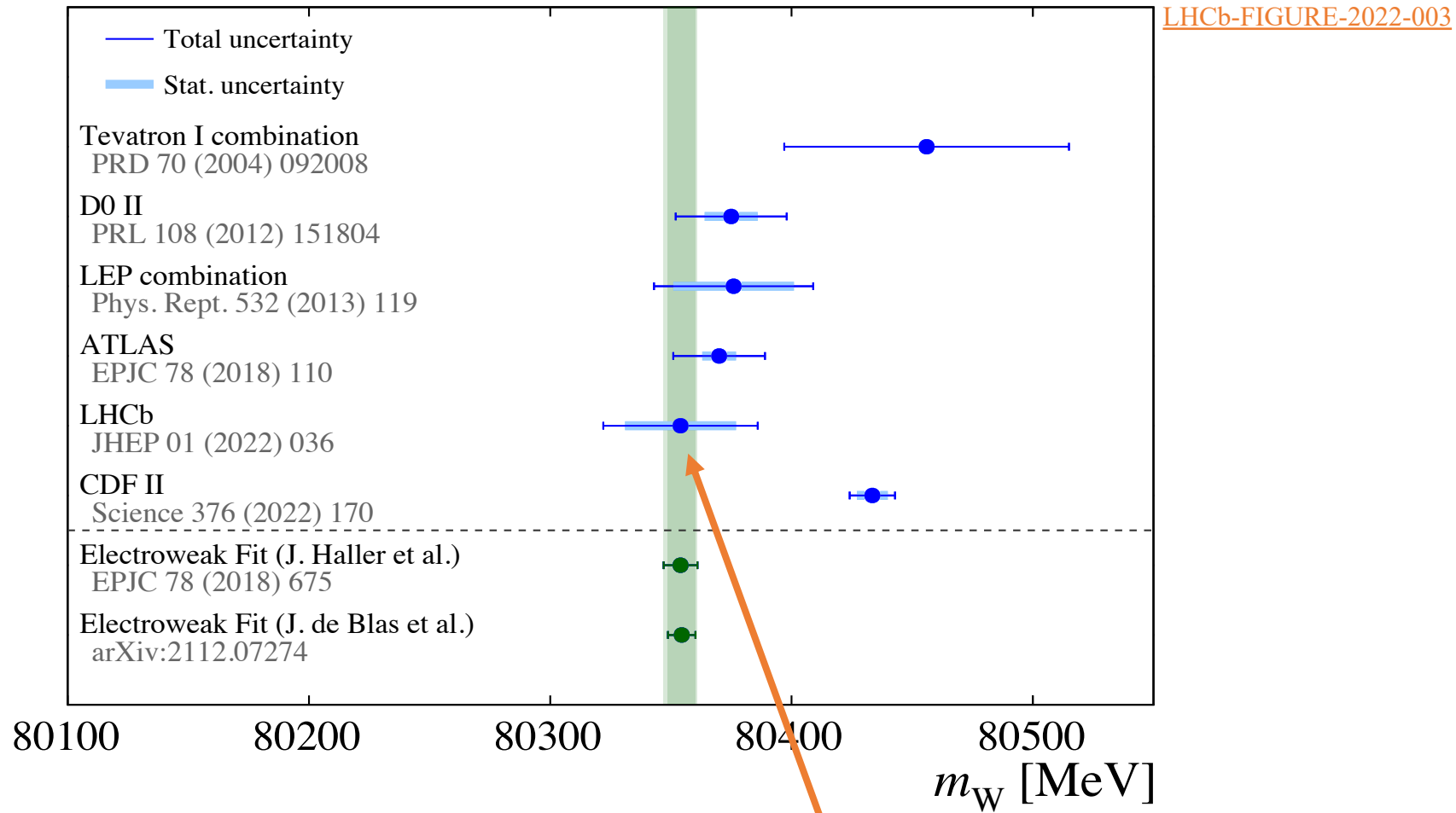
$$\sigma_{\text{stat}} = 23 \text{ MeV}$$

Uncertainty breakdown

Source	Size [MeV]	
Parton distribution functions	9.0	Average of NNPDF31, CT18, MSHT20
Theory (excl. PDFs) Total	17	
Transverse momentum model	11	Envelope from five different models
Angular Coefficients	10	Uncorrelated scale variation
QED FSR model	7	Envelope of Pythia8, Photos and Herwig7
Additional electroweak corrections	5	Tested with POWHEGW
Experimental Total	10	
Momentum scale and resolution modelling	7	Includes statistical uncertainties, details of the methods (e.g. binning, smoothing) and dependence on external inputs.
Muon ID, trigger and tracking efficiency	6	
Isolation efficiency	4	
QCD background	2	
Statistical	23	
Total	32	

The 2016 result

- Taking the arithmetic average of results with [NNPDF31](#), [CT18](#) and [MSHT20](#):



$$m_W = 80354 \pm 23_{stat} \pm 10_{exp} \pm 17_{theory} \pm 9_{PDF} \text{ MeV} = 80354 \pm 32 \text{ MeV}$$

Conclusions and outlook

- First measurement of the W boson mass at LHCb achieved a precision of ~ 32 MeV, using roughly 1/3 of the Run 2 dataset,
- Result was published at beginning of 2022 as [JHEP 01 \(2022\) 036](#),
- Largest uncertainties were theoretical (angular coefficients, PDFs and boson p_T model),
- Future prospects from LHCb \rightarrow talk on Friday by Miguel Ramos Pernas,
- LHCb measurement(s) expected to provide significant impact on a LHC-wide average due to potential anti-correlation of PDF uncertainties \rightarrow next talk by Maarten Bonnekamp.

Thank you for your attention.
Any questions?

Backup

Reconstruction & selection efficiencies

Each muon is **well-reconstructed & identified**, **fires relevant triggers** and is **isolated**.

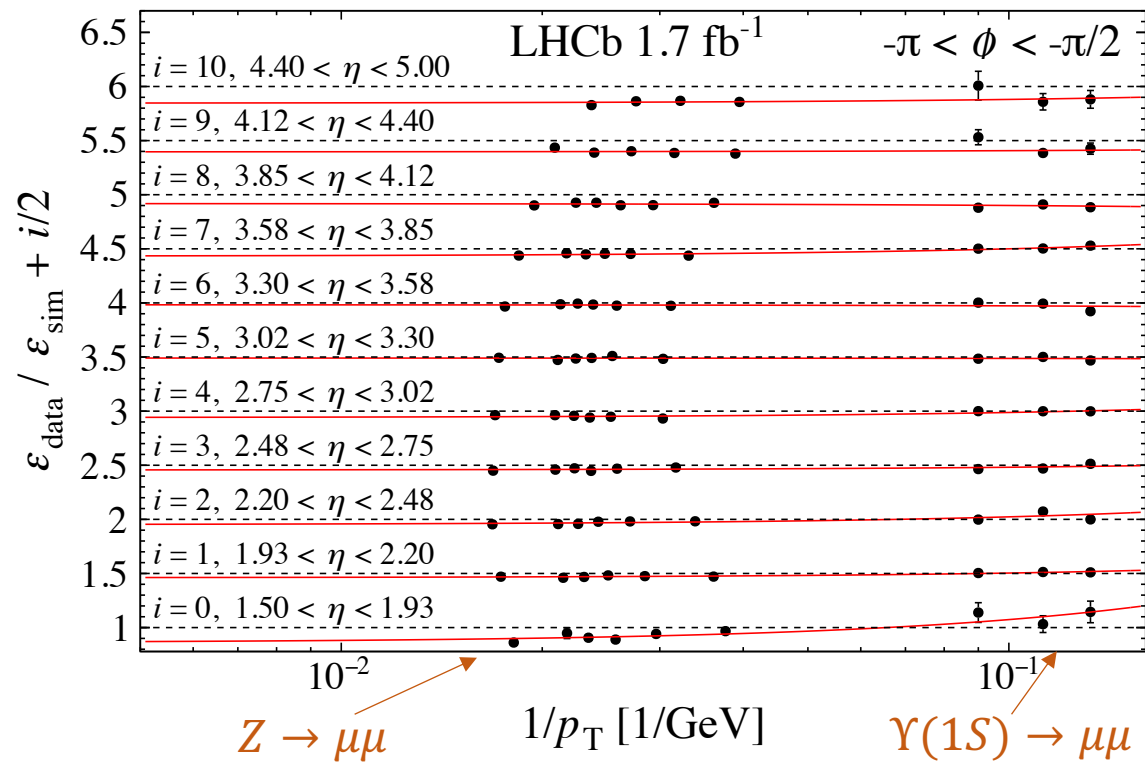
$$\epsilon_{sim}(p_T, \eta, \phi, \dots) = \epsilon_{data}(p_T, \eta, \phi, \dots) ?$$

Simulation corrected with event weights $w(p_T, \eta, \phi, \dots) = \epsilon_{data} / \epsilon_{sim}(p_T, \eta, \phi, \dots)$

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Reco, ID & trigger efficiencies:

- Tag & probe method with $Z \rightarrow \mu\mu$ and $\Upsilon(1S) \rightarrow \mu\mu$ gives ϵ_{sim} & ϵ_{data} .
- Weights from fit to efficiency ratio as function of p_T^μ , binned in η and ϕ .



Reconstruction & selection efficiencies

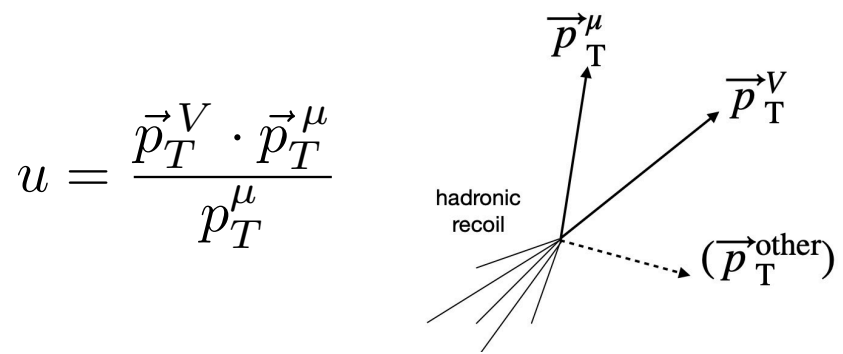
Each muon is well-reconstructed & identified, fires relevant triggers and is **isolated**.

$\varepsilon_{sim}(p_T, \eta, \phi, \dots) = \varepsilon_{data}(p_T, \eta, \phi, \dots) ?$

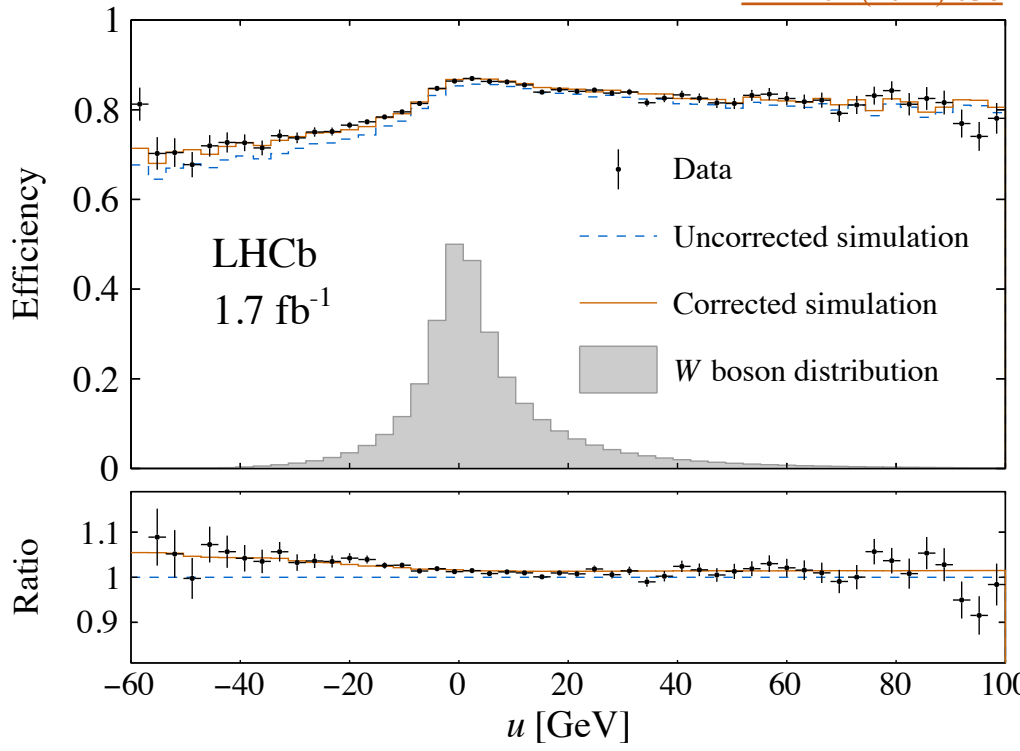
Simulation corrected with event weights $w(p_T, \eta, \phi, \dots) = \varepsilon_{data} / \varepsilon_{sim}(p_T, \eta, \phi, \dots)$

Isolation efficiencies:

- Tag & probe method with $Z \rightarrow \mu\mu$ gives ε_{sim} & ε_{data} .
- Weights from efficiency ratios binned in recoil projection u and η .



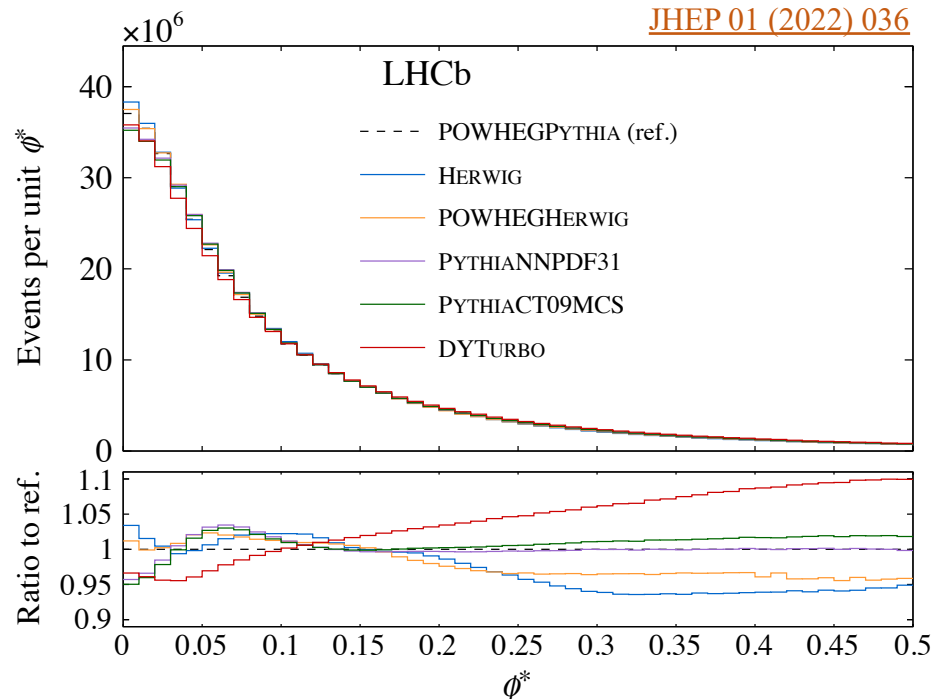
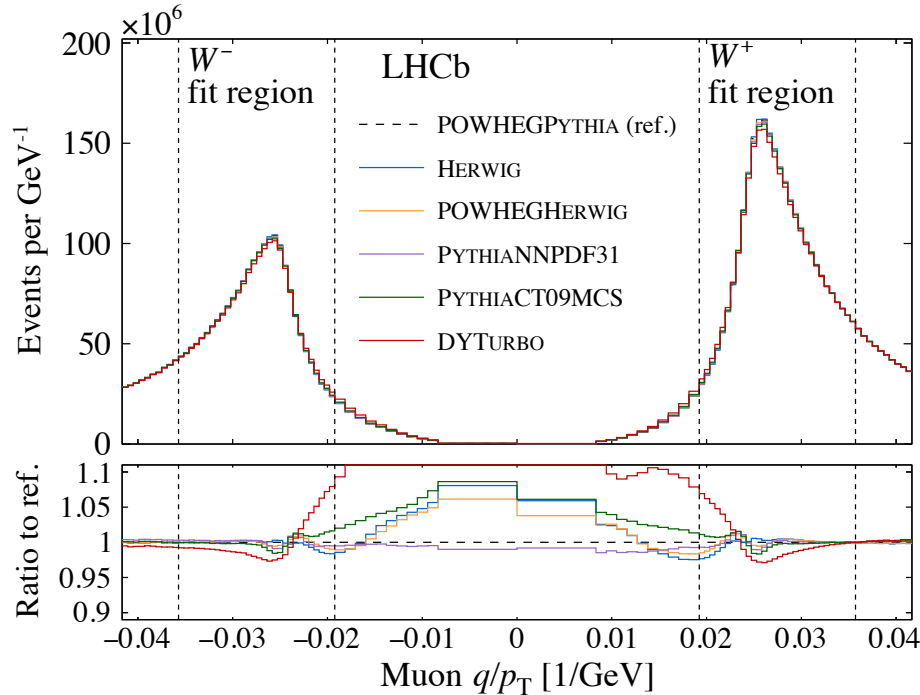
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Cross-checks

1. **Orthogonal splits:** Five ~50:50 splits of the data (polarity, charge \times polarity, etc...) all result in $[m_W]$ differences within 2σ .
2. **Fit range:** The result is stable w.r.t. variations in the upper/lower limits.
3. **Fit freedom:** The result is stable w.r.t. variations in the model freedom (e.g. 3 independent α_s values instead of 2, etc...)
4. **W-like fit of the Z mass:** Measurements with μ^+ and μ^- agree to better than 1σ and their average agrees with the PDG value to better than 1σ .
5. **δm_W fit:** Alternative fit with the difference between the W^+ and W^- masses as another floating parameter: this parameter is consistent with zero within 1σ .
6. **Additional tests** with NNLO PDFs instead of NLO PDFs, variations in the charm quark mass, etc... affect m_W at the ≈ 1 MeV level.

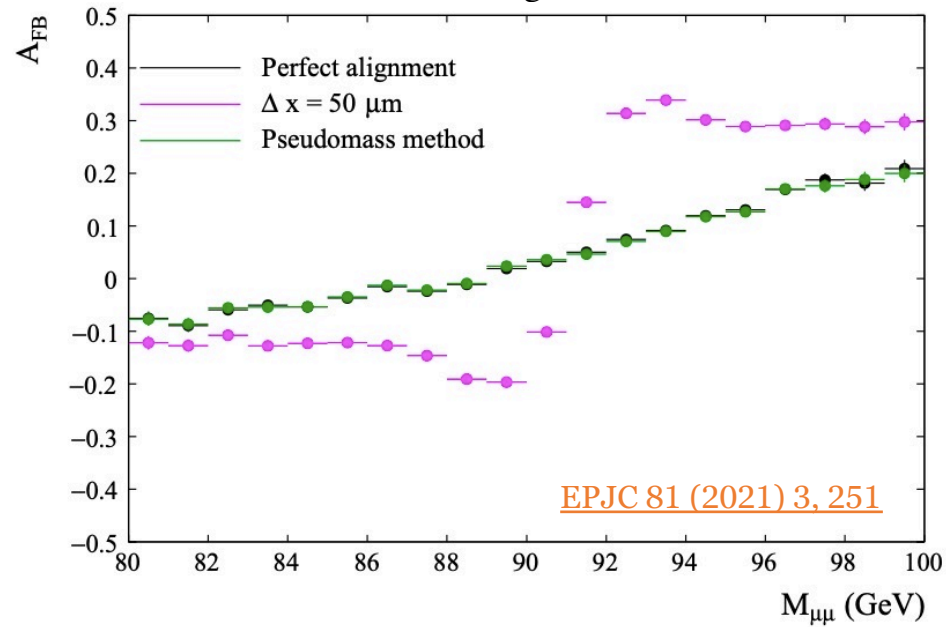
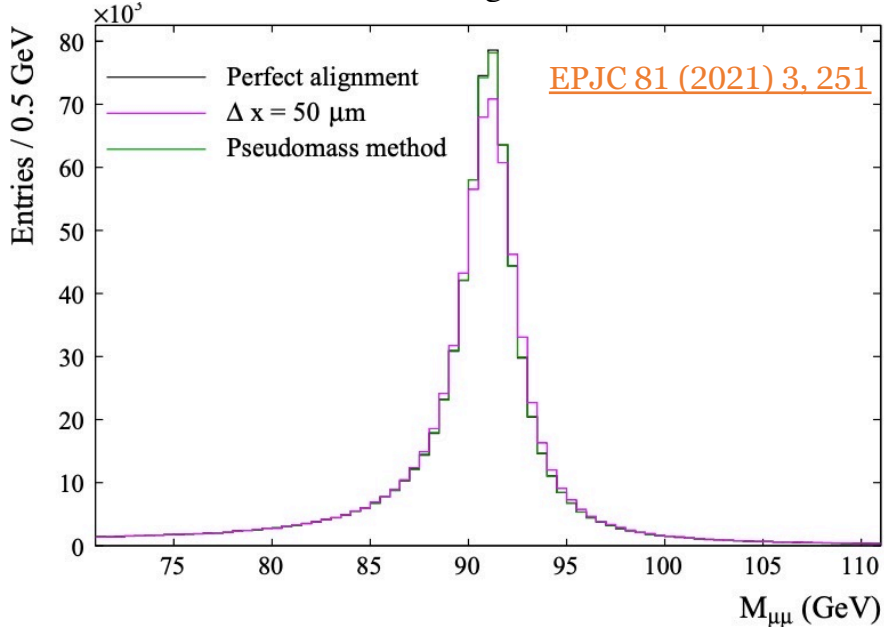
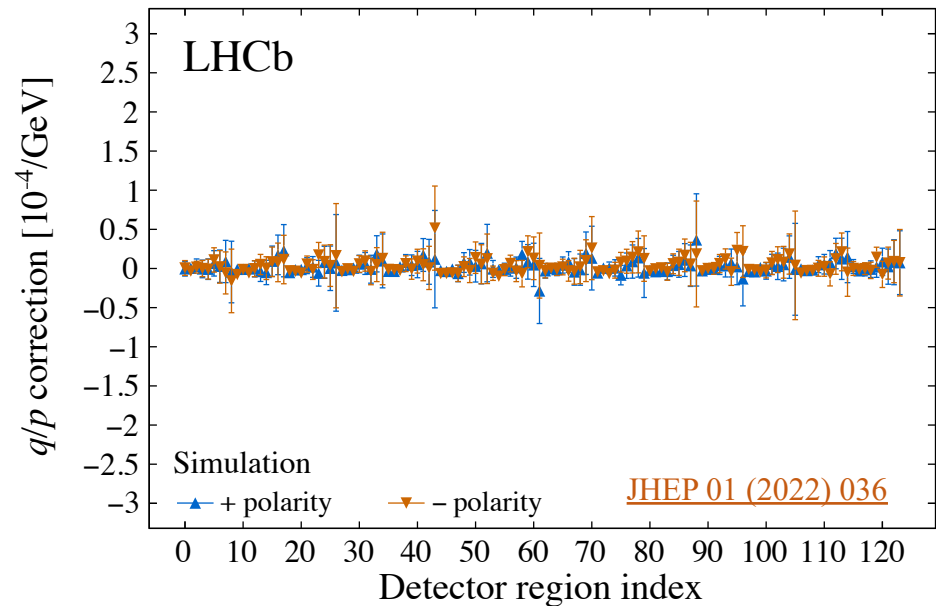
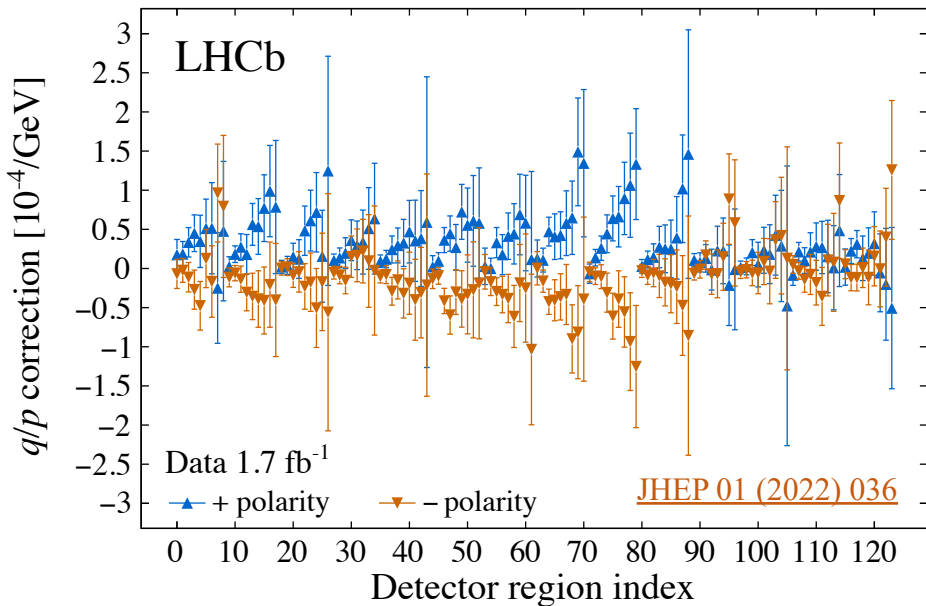
Model validation: [Pseudo]data challenges



- Using our central model to fit pseudodata generated from different models (e.g. HerwigNLO) gives a similar spread as using those different models to fit the real data.

Data config.	χ_W^2	χ_Z^2	δm_W [MeV]
POWHEGPYTHIA	64.8	34.2	—
HERWIG	71.9	600.4	1.6
POWHEGHERWIG	64.0	118.6	2.7
PYTHIA, CT09MCS	71.0	215.8	-2.4
PYTHIA, NNPDF31	66.9	156.2	-10.4
DYTURBO	83.0	428.5	4.3

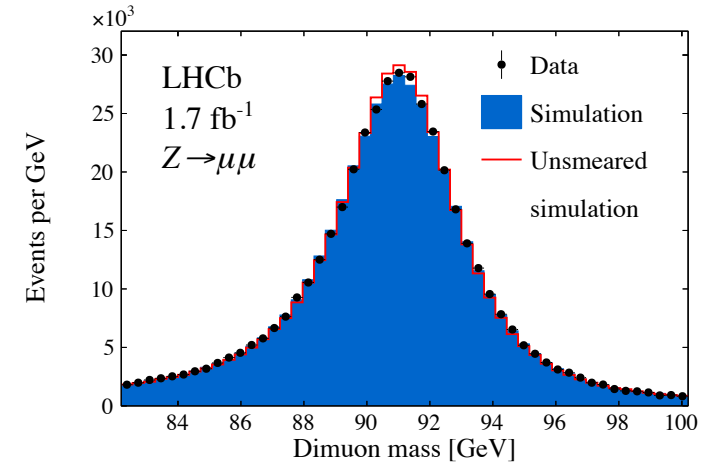
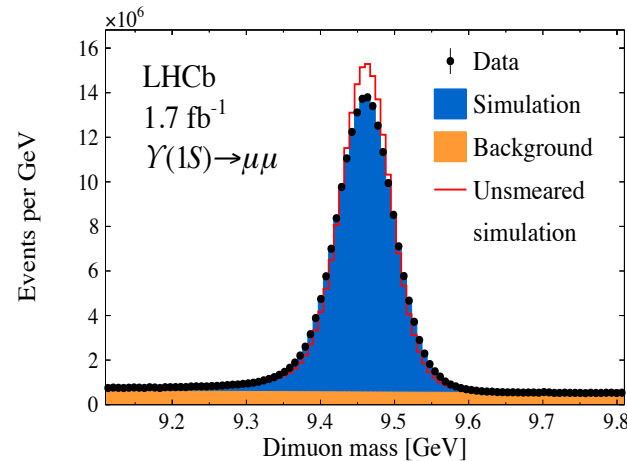
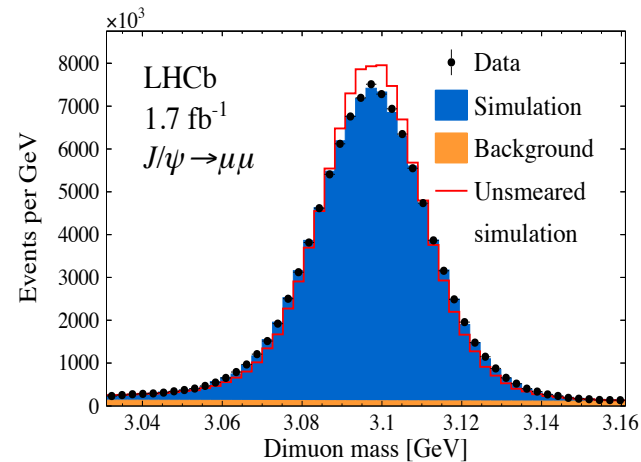
Charge-dependent curvature biases



Momentum smearing function

3) Additional smearing of the simulation to better model the data:

$$\frac{q}{p} \rightarrow \frac{q}{p \cdot \mathcal{N}(1 + \alpha, \sigma_{MS})} + \mathcal{N}\left(\delta, \frac{\sigma_\delta}{\cosh \eta}\right),$$

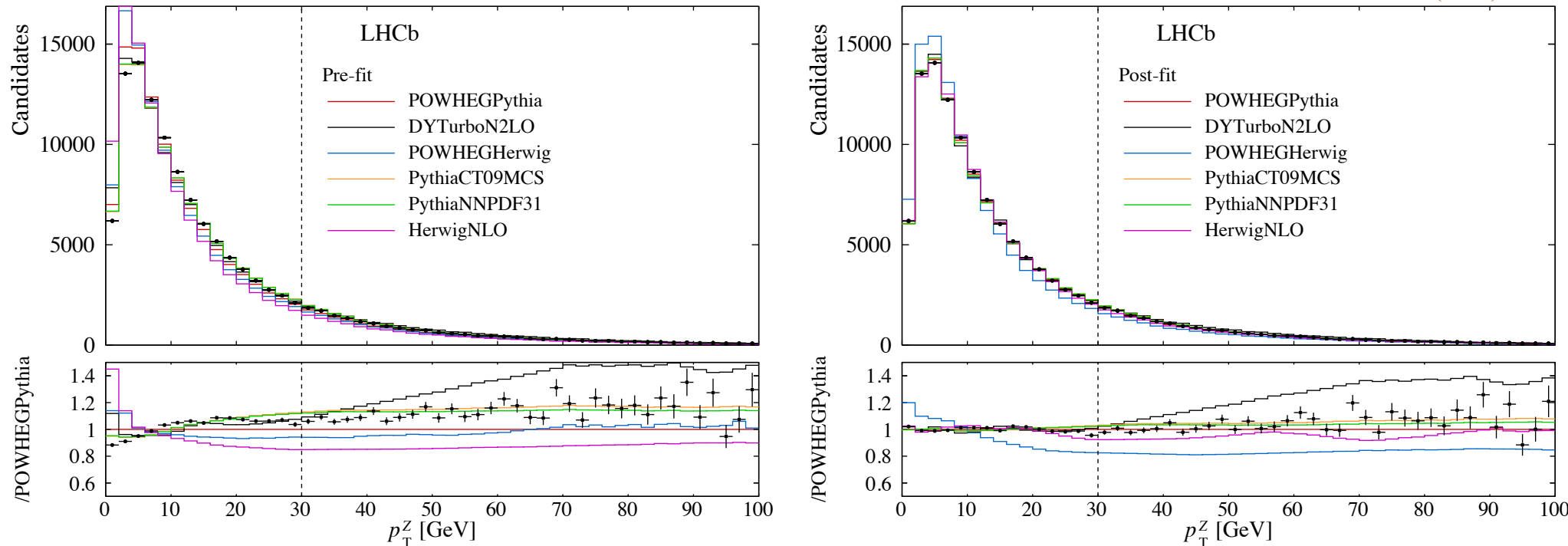


Effects modelled are curvature bias (δ), momentum *scale* ($1 + \alpha$), momentum-independent (σ_{MS}) and momentum-dependent (σ_δ) smearing.

POWHEG+Pythia as central model

Tuning of α_s and k_T^{intr}

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- POWHEG-Box + Pythia8 = best description of p_T^Z -> our central model.
- Other models (POWHEG-HERWIG, PYTHIACT09MCS, PYTHIANNPDF31 and HERWIG NLO) are used to evaluate systematic uncertainty (12 MeV).