## CNIS



# Low energy imprints of (Higgs) compositeness

Giacomo Cacciapaglia (IP2I Lyon)

Orsay 2023 W mass workshop

## Motivation

- · Composite models 'solve' the Hierarchy problem...
- with new scale in the multi-TeV!





multi-TeV mountain

#### What are we looking for?

- -> Precision EW + Higgs observables
- -> light composite scalars
- -> multi-TeV resonances (top partners, pNGBs, spin-1)

## Composite Higgs models 101



- · Symmetry broken by a condensate (of TC-fermions)
- Higgs and longitudinal Z/W emerge as mesons
   (pions)

Scales:

f : Higgs decay constant v : EW scale  $m_\rho \sim 4\pi f$ 

EWPTs + Higgs coupl. Limit:

 $f \gtrsim 4v \sim 1 \text{ TeV}$ 



## Composite Higgs models 101



- · Symmetry broken by a condensate (of TC-fermions)
- Higgs and longitudinal Z/W emerge as mesons
   (pions)

In the TC limit: f = vThe Higgs is a light scalar resonance (dilaton?)  $m_{\rho} \sim 4\pi f \sim 2 \text{ TeV}$ 



## A minimal case

T.Ryttov, F.Sannino 0809.0713 Galloway, Evans, Luty, Tacchi 1001.1361

	SU(2) <sub>TC</sub>	$SU(4)_{\psi}$	SU(2) <sub>L</sub>	<i>U</i> (1) <sub>Y</sub>
$\overline{\left(\begin{array}{c}\psi^1\\\psi^2\end{array}\right)}$			2	0
$\psi^3$			1	-1/2
$\psi^4$			1	1/2

Antisymmetric matrix  $\langle \psi^i \psi^j 
angle = \Sigma_0$ 

SU(4) -> Sp(4)

This theory is Asymptotic Free and confines in the IR!

Under the global symmetry SU(4):  $\Sigma_0 
ightarrow U \cdot \Sigma_0 \cdot U^T$ 

Higgs  $5_{\mathrm{Sp}(4)} 
ightarrow (2,2) \oplus (1,1)$ 



## WZW matters!



In QCD, coupling of the pions to EW gauge bosons are generated by (global) anomalies!

$$\mathcal{L}_{WZW} = \frac{d_{\psi}}{64\pi^2} \frac{\eta}{f} \left( g^2 W_{\mu\nu} \tilde{W}^{\mu\nu} - g'^2 B_{\mu\nu} \tilde{B}^{\mu\nu} \right)$$

o Predictive power!

Dimension of TC rep

 $d_{\psi} = 2$ 

Coupling to 2 photons vanishes!

## Composite Higgs models 101



How can light states emerge?



# The partial compositeness paradigm

Kaplan Nucl. Phys. B365 (1991) 259

 $\frac{1}{\Lambda_{\rm q}^{d-1}} \mathcal{O}_H q_L^c q_R \qquad \Delta m_H^2 \sim \left(\frac{4\pi f}{\Lambda_{\rm q}}\right)^{d-4} f^2 \qquad \text{Both irrelevant if}$ 

we assume:

 $d_H > 1$   $d_{H^2} > 4$ 

Let's postulate the existence of fermionic operators:

 $\frac{1}{\Lambda_{\rm fl.}^{d_F-5/2}} (\tilde{y}_L \ q_L \mathcal{F}_L + \tilde{y}_R \ q_R \mathcal{F}_R)$ 

This dimension is not related to the Higgs!

 $f(y_L \; q_L Q_L + y_R \; q_R Q_R)$  with  $y_{L/R} f \sim \left(rac{4\pi f}{\Lambda_{
m P}}
ight)^{d_F-5/2} 4\pi f$ 

Sequestering QCD in Partial compositeness G.Ferretti, D.Karateev  $\mathcal{G}_{\mathrm{TC}}$  : rep R' rep R 1312.5330, 1604.06467 Q $\chi$  $T' = QQ\chi$  or  $Q\chi\chi$ SM: EW colour + hypercharge global :  $\langle QQ \rangle \neq 0$ a)  $\langle \chi \chi \rangle \neq 0$ coloured pNGBs di-boson PNGB Higgs b)  $\langle \chi \chi \rangle = 0$ DM?

light top partners from 't Hooft anomaly conditions?

## Starting with misalignment models

Effect expressed in terms of oblique parameters: S &

Effects from H Loops



composite case: see 1502.04718



$$\Delta T = -\frac{3}{16\pi c_W^2} \sin^2\theta \ln \frac{\Lambda^2}{m_h^2}$$

$$\Delta S = \frac{1}{12\pi} \sin^2 \theta \ln \frac{\Lambda^2}{m_h^2} + S_{\rm UV} \sin^2 \theta$$

Arrows: naive contribution of top partner loops.

Contribution of top partners  
Top-partners are baryons of the new strong sector:  

$$EW gauge$$
Interactions  

$$\mathcal{L}_{composite} = tr \left[ \bar{\psi}_A \ i \ D \ \psi_A \right] - M_A \ tr \left[ \bar{\psi}_A \psi_A \right] + tr \left[ \bar{\psi}_1 \ i \ D \ \psi_1 \right] - M_1 \ tr \left[ \bar{\psi}_1 \psi_1 \right] + (2.12)$$

Seffect of mixing with the top/bottom fields usually considered:

$$\begin{split} \mathcal{L}_{\mathrm{mix}} &= y_{L1} f \ tr \left[ A_{L}^{\dagger} \gamma_{0} U_{\Pi} \psi_{5} U_{\Pi}^{T} \right] + y_{R1} f \ tr \left[ A_{R}^{\dagger} \gamma_{0} U_{\Pi} \psi_{5} U_{\Pi}^{T} \right] \\ &+ y_{L2} f \ tr \left[ A_{L}^{\dagger} \gamma_{0} U_{\Pi} \psi_{1} U_{\Pi}^{T} \right] + y_{R2} f \ tr \left[ A_{R}^{\dagger} \gamma_{0} U_{\Pi} \psi_{1} U_{\Pi}^{T} \right] \\ &+ y_{R}' f \ tr \left[ A_{R}^{(2)\dagger} \gamma_{0} U_{\Pi} \psi_{5} U_{\Pi}^{T} \right] + \mathrm{h.c.} \,, \end{split}$$

$$A_L = \begin{pmatrix} 0 & 0 & \frac{t_L}{\sqrt{2}} & 0 \\ 0 & 0 & \frac{b_L}{\sqrt{2}} & 0 \\ -\frac{t_L}{\sqrt{2}} & -\frac{b_L}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Mixing angles:  $\sin[\phi_L] \propto y_L$   $\sin[\phi_R] \propto y_R$   $Y_{top} \propto \sin[\phi_L] \sin[\phi_R] \frac{f}{M}$ 

Contribution of top partners  
Top-partners are baryons of the new strong sector:  

$$EW gauge$$
Interactions  

$$\mathcal{L}_{composite} = tr \left[ \bar{\psi}_A \ i \ D \ \psi_A \right] - M_A \ tr \left[ \bar{\psi}_A \psi_A \right] + tr \left[ \bar{\psi}_1 \ i \ D \ \psi_1 \right] - M_1 \ tr \left[ \bar{\psi}_1 \psi_1 \right] + \kappa' \ tr \left[ \bar{\psi}_A \ A \ \psi_A \right] + \kappa \left( tr \left[ \bar{\psi}_A \ A \ \psi_1 \right] + h.c. \right).$$
(2.12)

 Additional couplings in the strong sector: generate corrections to the W/Z coupling via misalignment

$$d_{\mu} = -\frac{\sqrt{2}}{f} \partial_{\mu} \left( hX^{4} + \sum_{a=5}^{\dim(\mathcal{G}/\mathcal{H})} \eta^{a} X^{a} \right) + \frac{\sin\alpha}{\sqrt{2}} \sum_{i=1}^{3} \left( g_{2} W_{\mu}^{i} - g_{1} B_{\mu} \delta^{i3} \right) X^{i} + \cdots$$

New contribution to S & T

H.Cai, G.C., 2208.04290

## Contribution of top partners SU(4)/Sp(4)

H.Cai, G.C., 2208.04290

Bi-doublet:

PDG

4

CDF

6

m<sub>T</sub> (TeV)

8

10

 $Sin[\phi_L]=0.3, Sin[\alpha]=0.15$ 



Top-partner masses in the multi-TeV scale, 0 compatible with lattice results

New effects crucial in some cases (bi-doublet) 0

#### "Technicolor" models

Compositeness scale f=v

The Higgs arises as a "light" resonance (dilaton-like?)

$$\mathcal{L} = \mathcal{L}_{\overline{\mathrm{SM}}} + \xi_V \left( 1 + 2\kappa_V \frac{h}{v} + \kappa_{2V} \frac{h^2}{v^2} \right) \frac{v^2}{4} \operatorname{Tr} D_\mu U^\dagger D^\mu U + \frac{1}{2} \partial_\mu h \partial^\mu h - \frac{\tilde{m}_h^2}{2} h^2 \left( 1 + V_{0,1} \lambda_{3h} \frac{h}{v} \right)$$

The couplings of the Higgs stem from the strong dynamics, and they could be larger than the SM ones!

$$T = -\frac{3}{16\pi c_W^2} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{m_h^2}, \qquad S = \frac{1}{12\pi} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{m_h^2} + \Delta S_{\rm UV},$$

 $\kappa_V > 1 \implies \Delta T > 0!!!$ 

#### "Technicolor" models

Compositeness scale f=v

$$T = -\frac{3}{16\pi c_W^2} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{m_h^2}, \qquad S = \frac{1}{12\pi} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{m_h^2} + \Delta S_{\rm UV}, \qquad \Delta S_{\rm UV} = \frac{n_d}{6\pi}$$

G.C., F.Sannino 2204.04514



 Possible to obtain sizeable shifts in the W mass, contribution from top partners (to be added) can further help!

#### "Technicolor" models

G.C., C.Cot, F.Sannino 2104.08818

The most intriguing "coincidence" stemming from the anomalous muon g-2 !

$$\Delta a_{\mu} = \frac{m_{\mu}^2}{\Lambda_{\rm TC}^2} \left( 1 + \frac{(y_L y_L)_{\mu\mu}}{g_{\rm TC}^2} \right)$$

$$\Delta a_{\mu} = 250 \cdot 10^{-11}$$
 for  $\Lambda_{\rm TC} = 2$  TeV ~  $4\pi v$ 

suppressed in misaligned models:

$$\Delta a_{\mu}(\text{CGH}) \approx \frac{g_{*}^{2}}{(4\pi)^{2}} \frac{m_{\mu}^{2}}{m_{*}^{2}} = \frac{v_{\text{SM}}^{2}}{f_{\text{CGH}}^{2}} \frac{m_{\mu}^{2}}{\Lambda_{\text{TC}}^{2}},$$

#### Old RK result:

