



Quantum Simulation of Hamiltonian Using Variational Approach Involving Linear Combination of Unitary Operators

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- Introduction
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Introduction

- Why Quantum simulations?
 - Certain problems can be solved in polynomial time on the quantum computer, which cannot be solved in polynomial time on a classical computer.
 - Classical simulations face difficulty in simulating any quantum system due to exponential growth in resource requirements.
- Similar complexities occur in nuclear many-body problems.
- We are trying to solve the smaller problem on the quantum computer of the present NISQ era.
- Also, trying to give proof of concept that problems can be solved on the quantum computer once they are a practical reality in the near future.



P.C.: IBM Quantum



- In the last decade, quantum computers have been used to simulate various molecules, and a wide variety of chemistry problems are being simulated.
- In recent years researchers have drawn their attention to nuclear physics problems which can be solved on a quantum computer.

E. Dumitrescu et al. PRL **120** (2018) 210501.
PS and PA, PRC **104** (2021) 235501.

Motivation for current work.

A. Roggero et al. PRC 100 (2019) 034610.



Fig. 1: Deuteron binding energy as function of number of basis states.







- In Childs *et al. QIC* **12** (2012), 901-924, a new approach for simulating Hamiltonian dynamics based on implementing linear combinations of unitary (LCU) operations is proposed.
- LCU method is used in PS and PA, PRC **105** (2022) *064318,* to calculate the expectation value of operators.
- In TB and SK, JCP **154** (2021) 194107, proposed a technique to simulate molecular resonances.
- In A. Roggero et al. PRC **102** (2020) 064624, used LCU for creation of excited states.
- Motivated by these, we are trying to modify the variational algorithm and implement it for complex energy calculations.

Encodings and quantum algorithms

- We transform the nuclear problem in such a way that it can be solved on a quantum computer.

$$H \,=\, \sum_{ij}\, h_{ij}a_i^\dagger a_j$$
 Where, $h_{ij}\,=\,\, < i \Big| \hat{T} + \hat{V} \Big| j > 0$

- We map many-body basis states and operators to qubit basis states and operations respectively.
- This mapping is done via various encodings, among which we will primarily use **Jordan-Wigner transformation** (JWT) and **Gray code encoding** (GC).

$$a_j^\dagger \,=\, rac{1}{2}(X_j-iY)\otimes \prod_{k < j} Z_k, \qquad a_j \,=\, rac{1}{2}(X_j+iY)\otimes \prod_{k < j} Z_k.$$

- After mapping, the second quantized hamiltonian transforms to Pauli strings which can be applied to quantum circuits.
- We have various algorithms like VQE, Quantum phase estimation (QPE) and LCU etc., which can be used to simulate the problem.



- For simplicity, we choose the problem of deuteron BE calculation.
- Hamiltonian for N=2 , N=3 and N= 4 are:
- $H_2 \,=\, 5.9067091 I_0 6.34329 Z_0 4.28661 X_0$
- $egin{array}{rll} H_3 \ = \ 7.765855 I_0 I_1 7.984145 Z_0 1.859145 I_0 Z_1 + 1.640855 Z_0 Z_1 \ \ 2.143305 (X_0 I_1 + X_0 Z_1) 3.91312 (I_0 X_1 Z_0 X_1) \end{array}$
- $egin{array}{rll} H_4 \ = \ 14.328I 7.814X_0 3.913X_1 + 3.913Z_0X_1 1.422Z_0 \ \ 8.422Z_1 + 3.527X_0Z_1 4.922Z_0Z_1 \end{array}$
 - O. D. Matteo et al. PRA **103** (2021), 042405.
- PS and PA, PRC **104 (**2021) *235501*.

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Variational Quantum Eigensolver (VQE)



Fig. 2: Architecture of the variational-quantum eigensolver.

|0| $Y(\theta$ (a) $Y(\theta_1$ $|0\rangle$ $Y(\theta_2$ $-\theta_{2}$ (b) $Y(\theta_2)$ $Y(\theta_1)$ $Y(\theta_3$ (c)

Fig. 3: Parameterized ansatz for N=2, 3 and 4 for GC encoding.

PS and PA, PRC 104 (2021) 235501

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A. Peruzzo and J. McClean et al. Nature Communications 5, 4213 (2014)

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Table 1: Energy calculated with QASM simulator for EFT potential

Fig. 4: Energy as a function of number of states for EFT potential for GC encoding

Expectation value through LCU



• The method for calculating expectation value based on LCU is discussed in PS and PA, PRC **105** (2022) 064318 and A. Roggero et al. PRC 102 (2020) 064624.

$$\mathcal{O} = \sum_{i=0}^{k-1} \beta_i U_i$$

$$V_P |0\rangle^{\otimes n_a} = \frac{1}{\sqrt{\Lambda}} \sum_i \sqrt{\beta_i} |i\rangle,$$

$$Where, \Lambda = \sum_i \beta_i \text{ and } W = V_P^{\dagger} V_S V_P$$

$$W |0\rangle^{\otimes n_a} |\psi\rangle = \frac{1}{\Lambda} |0\rangle^{\otimes n_a} \mathcal{O} |\psi\rangle + |\Psi^{\perp}\rangle.$$

$$|0\rangle$$

Fig. 5: General circuit for LCU

Continue ...



• Once we get the $\mathscr{O} |\psi\rangle$ we can calculate expectation value of operator by taking its overlap with $|\psi\rangle$



- J.C. Garcia-Escartin et al. PRA **87** (2013), 052330.
- PS and PA, PRC **105** (2022), 064318

Modified algorithm





Fig. 7: Modified variational algorithm for expectation value calculation

- We have applied this algorithm for binding energy of deuteron under EFT potential, and Hamiltonian is encoded via Gray Code encoding.
- Classical optimizer used for our calculation is SPSA optimizer.
- Executed the algorithm for 100000 shots and 100 independent runs.
- For error bar we used MAD.

Basis size (N)	True Value	Modified variational algorithm
2	-1.74916	-1.7459993 ± 0.025
3	-2.04567	-2.0464696 ± 0.044
4	-2.145	-2.0992413 ± 0.078

Table 2: Energy calculated with QASM simulator for EFT potential through modified algorithm



Results





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- We modified the variational algorithm such that it can be applied to system having non hermitian Hamiltonian.
- We calculated the BE of deuteron with our method and results are in agreement with the published VQE results.
- Even though this method is more prone to error due to larger gate depth, this method can be applied to the non-hermitian matrix, which occurs in the case of resonances.
- This shows a validity of method and further it can be extended for the Hamiltonian having complex energies.

