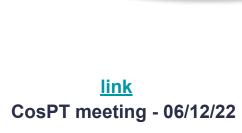






Adrien La Posta IJClab

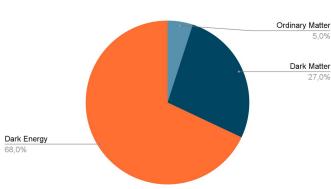


$$ds^{2} = -c^{2}dt^{2} + a^{2}(t)\left(\frac{dr^{2}}{1 - Kr^{2}} + r^{2}d\Omega^{2}\right)$$

#### Friedmann equation

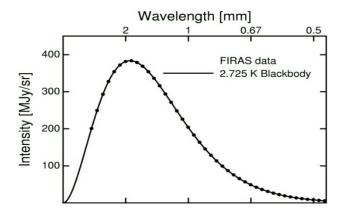
$$H^2(z) = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \left[ (\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_r^0(1+z)^4 + \rho_\Lambda \right]$$
 baryon radiation dark energy

CDM





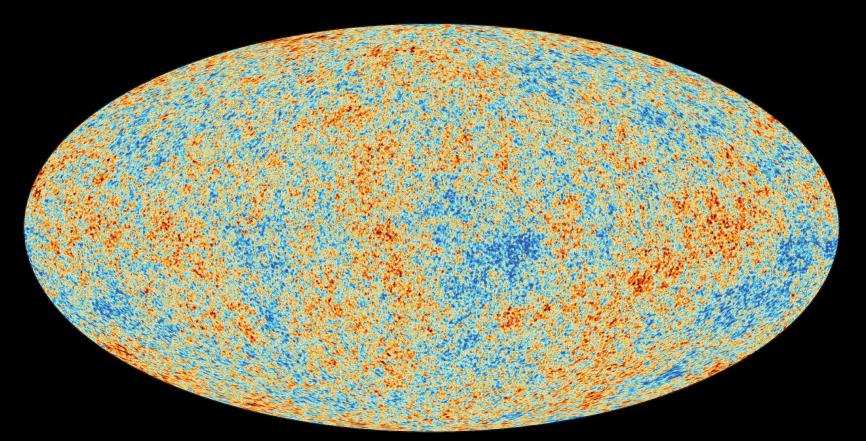




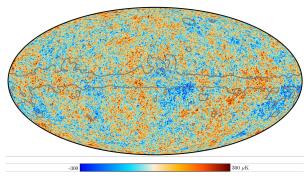
#### **Nearly isotropic blackbody spectrum at T = 2.725 K**

$$\frac{\delta T}{T} \sim 10^{-5}$$

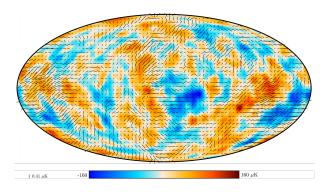
# CMB temperature as measured by the Planck satellite



#### **Temperature**



#### **Polarization E-modes**

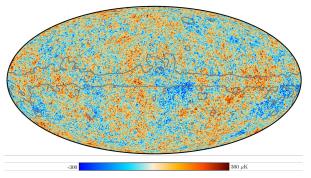


#### **Spherical harmonics**

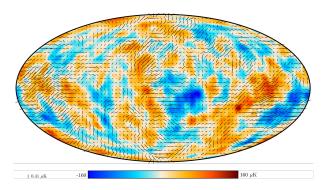
$$\delta T(\hat{n}) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} a_{\ell m}^T Y_{\ell}^m(\theta, \phi)$$

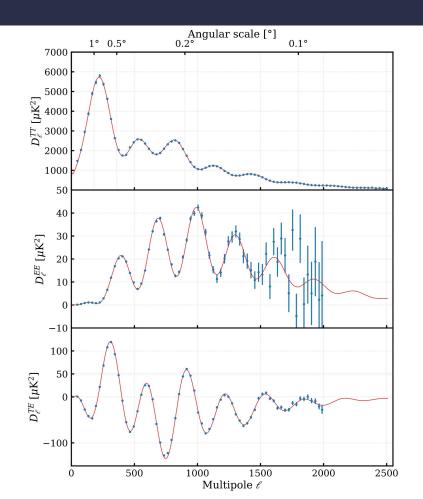
$$\langle a_{\ell m}^T a_{\ell' m'}^{T*} \rangle = \delta_{\ell \ell'} \delta_{m m'} C_{\ell}^{TT}$$

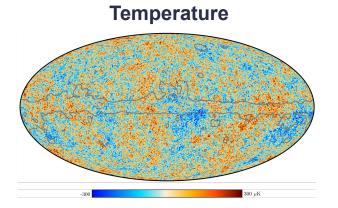




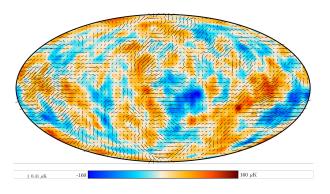
#### **Polarization E-modes**

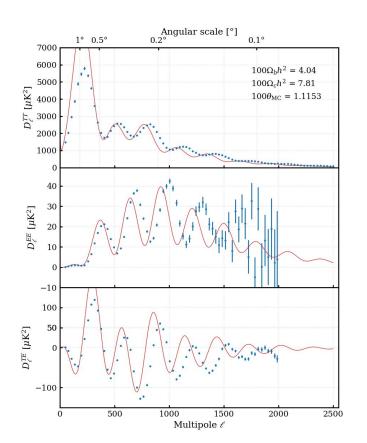








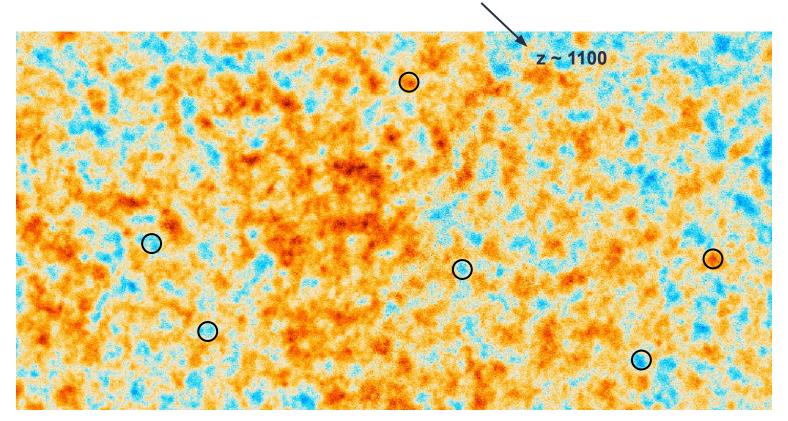


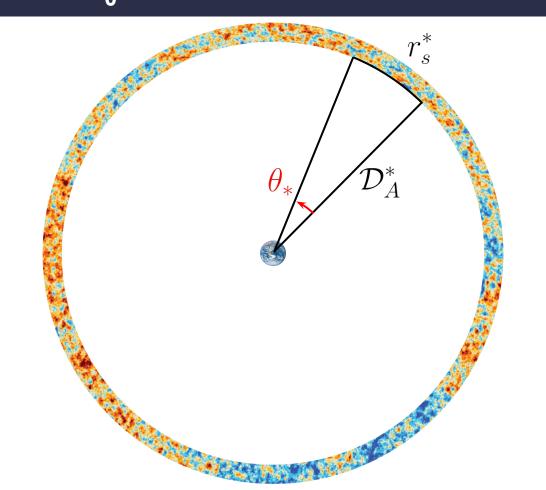


$$\rightarrow \theta_* \rho_b^0 \rho_c^0$$

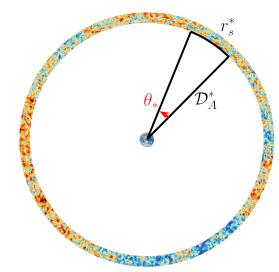
CMB standard ruler: size of the sound horizon at decoupling imprinted in the CMB radiation

**CMB standard ruler**: size of the sound horizon at decoupling imprinted in the CMB radiation

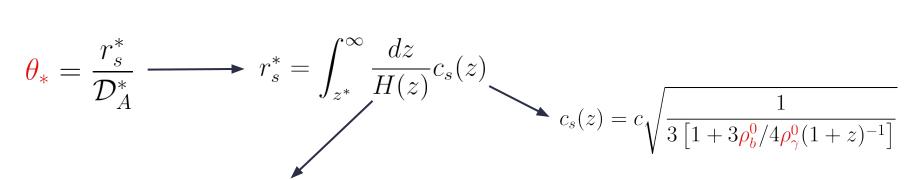




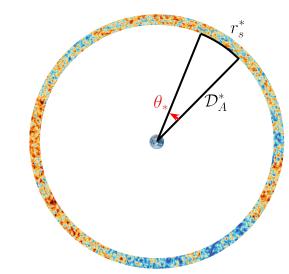
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$



## How to measure $H_0$ from the CMB?

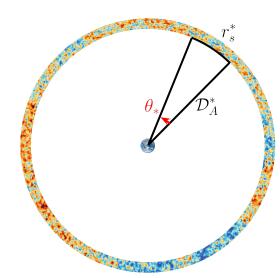


$$H_{\text{early}}^{2}(z) = \frac{8\pi G}{3} \left[ \rho_{r}^{0} (1+z)^{4} + (\rho_{b}^{0} + \rho_{c}^{0})(1+z)^{3} \right]$$



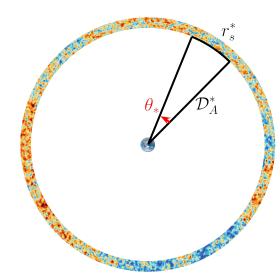
Now  $\mathcal{D}_A^*$  is known

$$heta_* = rac{r_s}{\mathcal{D}_s}$$

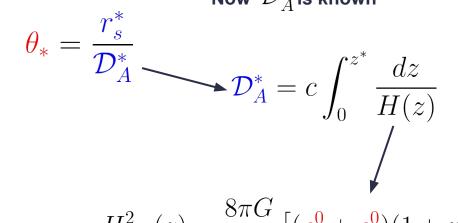




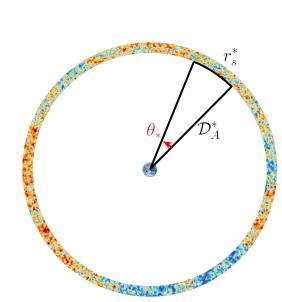
$$heta_* = rac{r_s^*}{\mathcal{D}_A^*}$$
  $heta_A$  is known  $\mathcal{D}_A$  is







$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} \left[ (\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_{\Lambda} \right]$$

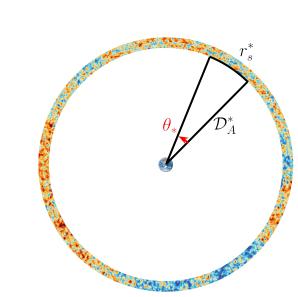


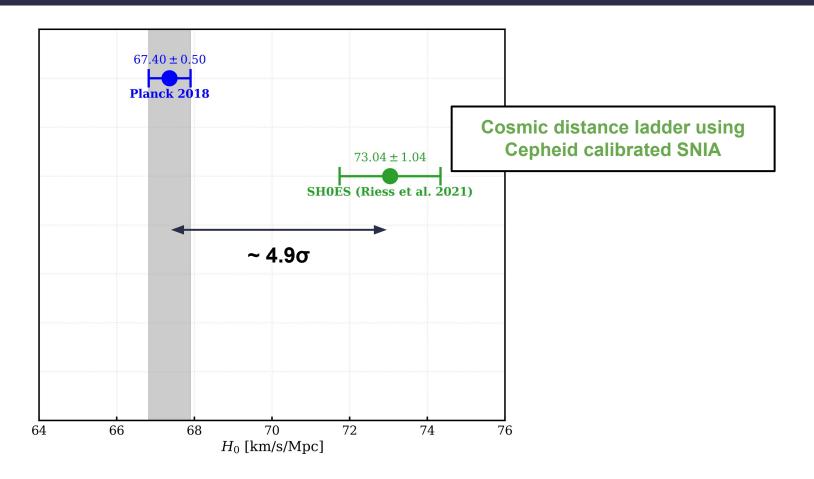
#### Now $\mathcal{D}_A^*$ is known

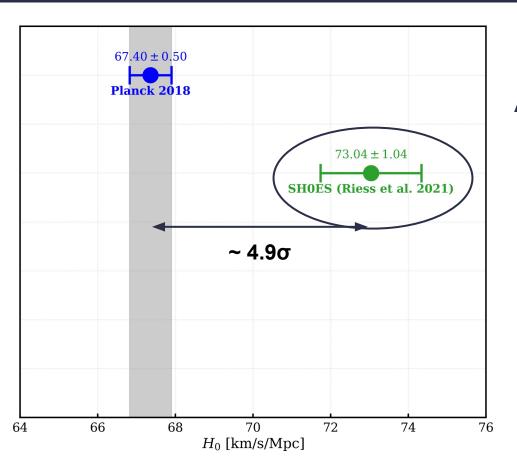
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow \mathcal{D}_A^* = c \int_0^{z^*} \frac{dz}{H(z)}$$

$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} \left[ (\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda \right]$$

$$H_0^2 = \frac{8\pi G}{3} \left[ \rho_b^0 + \rho_c^0 + \rho_\Lambda \right]$$

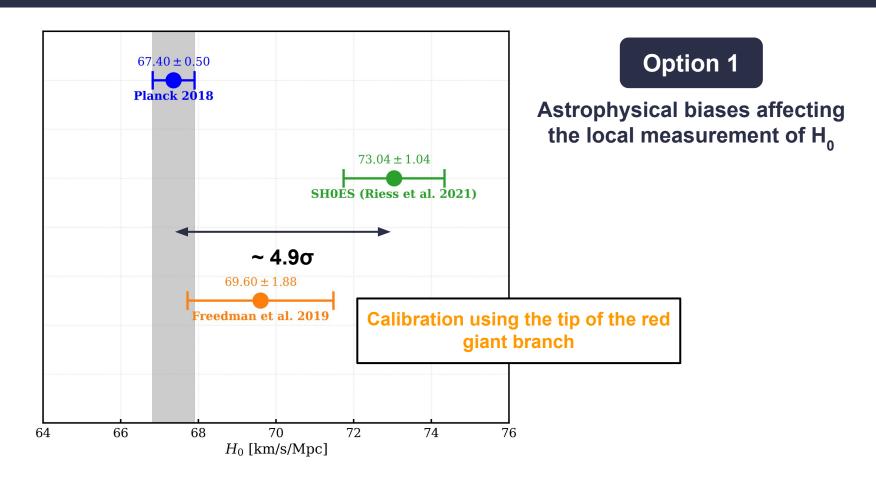


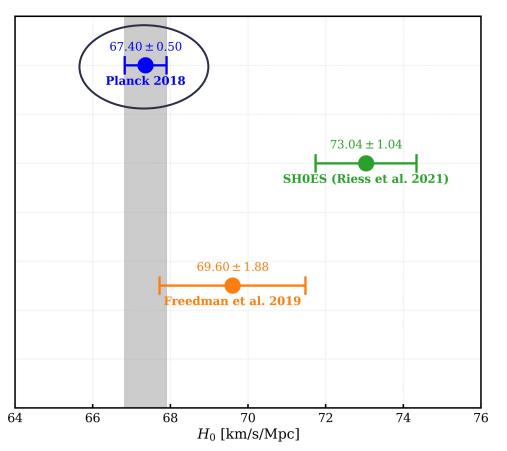




#### **Option 1**

Astrophysical biases affecting the local measurement of H<sub>0</sub>



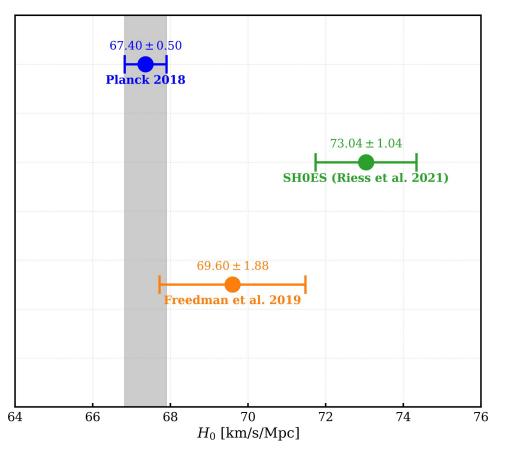


#### **Option 1**

Astrophysical biases affecting the local measurement of H<sub>0</sub>

#### Option 2

Instrumental systematic effect biasing the value of H<sub>0</sub> inferred from the CMB



#### **Option 1**

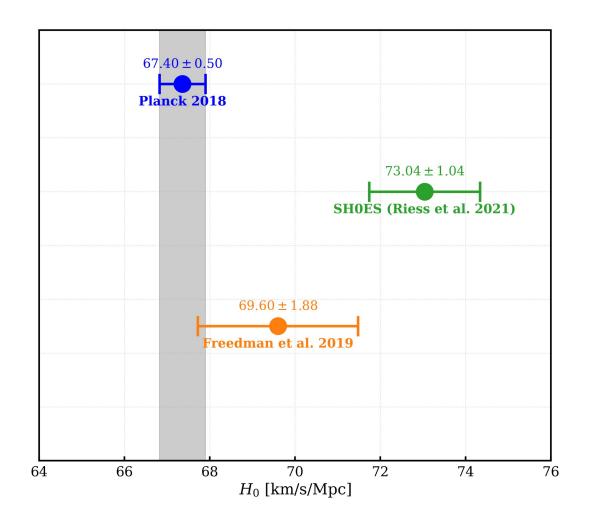
Astrophysical biases affecting the local measurement of H<sub>0</sub>

#### Option 2

Instrumental systematic effect biasing the value of H<sub>0</sub> inferred from the CMB

Option 3

Physics beyond ΛCDM



#### Option 1

Astrophysical biases affecting the local measurement of H<sub>0</sub>

Option 2

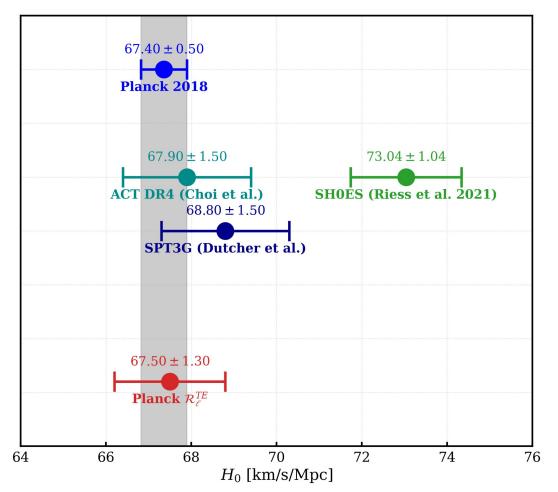
Instrumental systematic effect biasing the value of H<sub>0</sub> inferred from the CMB

Option 3

Physics beyond ΛCDM

# Independent measurements of H<sub>0</sub> from the ground





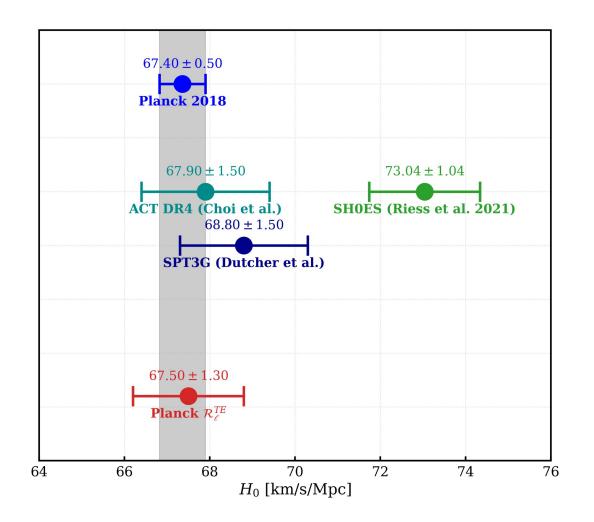
#### Option 2

Instrumental systematic effect biasing the value of H<sub>0</sub> inferred from the CMB



Hard to shift the CMB inferred H<sub>0</sub> with a systematic effect :

- Independent measurements from Planck, ACT and SPT
- Constraint from the correlation coefficient, robust against multiplicative systematics



#### Option 1

Astrophysical biases affecting the local measurement of H<sub>0</sub>

Option 2

Instrumental systematic effect biasing the value of H<sub>0</sub> inferred from the CMB

**Option 3** 

Physics beyond ΛCDM

### Early-time modification to ΛCDM

**Motivation :** obtain a higher value of  $H_0$  from the CMB

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**Motivation**: obtain a higher value of  $H_0$  from the CMB  $\longrightarrow$  lower  $\mathcal{D}_A^*$ 

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$$\theta_* = \frac{r_s^*}{D_A^*} \longrightarrow \text{ Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z) \\ \frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z)$$
 observations

#### One proposed solution : Early Dark Energy

**Motivation**: obtain a higher value of  $H_0$  from the CMB  $\longrightarrow$  lower  $\mathcal{D}_A^*$ 

$$\theta_* = \frac{r_s^*}{D_A^*} \longrightarrow \text{ Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$
 
$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z) + \rho_{\text{EDE}}(z)$$
 observations

The EDE component is described as a scalar field  $\,\phi$  (Poulin+ 2019, Smith+ 2019)

Background evolution : 
$$\ddot{\phi}+3H\dot{\phi}+V'(\phi)=0$$
 axion-like potential 
$$V(\phi)=m^2f^2\left[1-\cos\left(\frac{\phi}{f}\right)\right]^3$$

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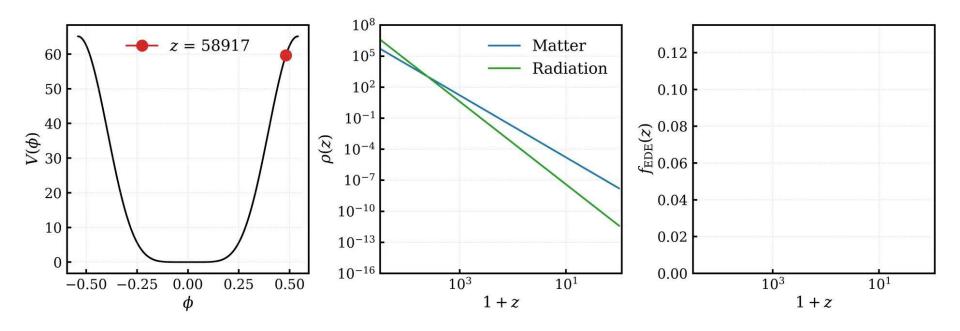
 $\phi_i$ : initial field value

#### Early Dark Energy: frozen at early times

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

The field is initially frozen due to Hubble friction (H >> m)

acts as dark energy (w= - 1)

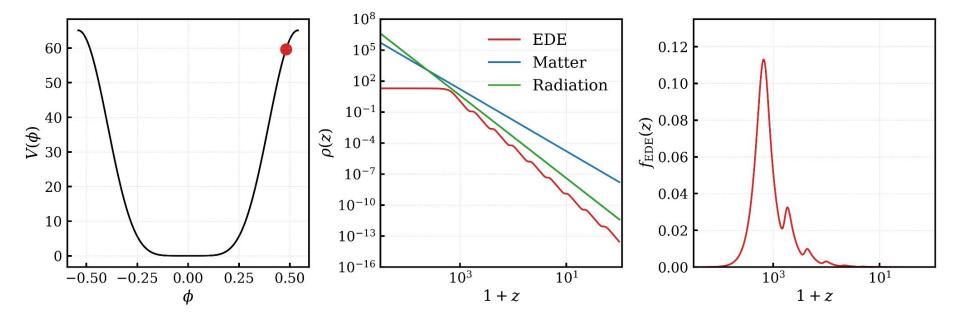


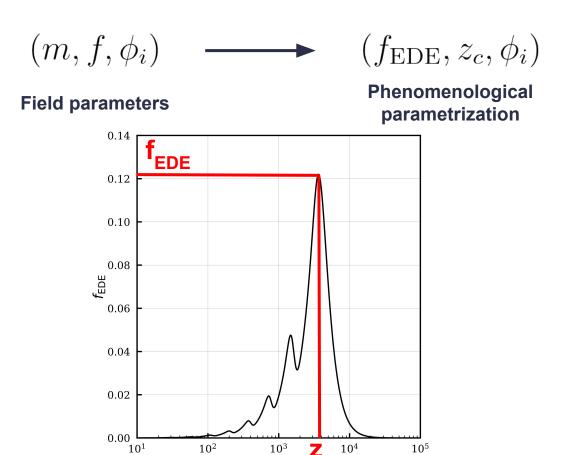
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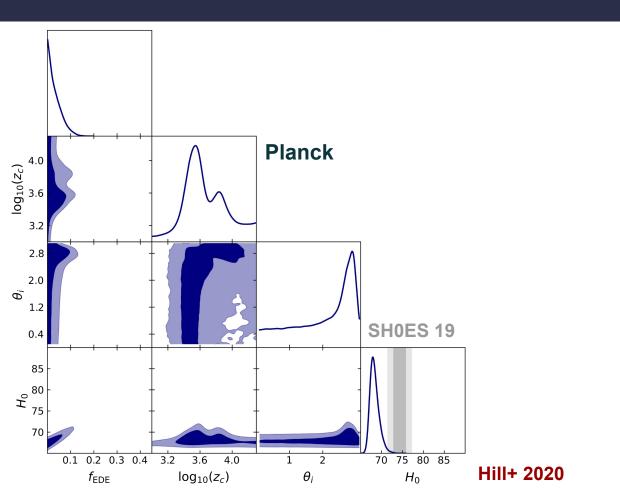
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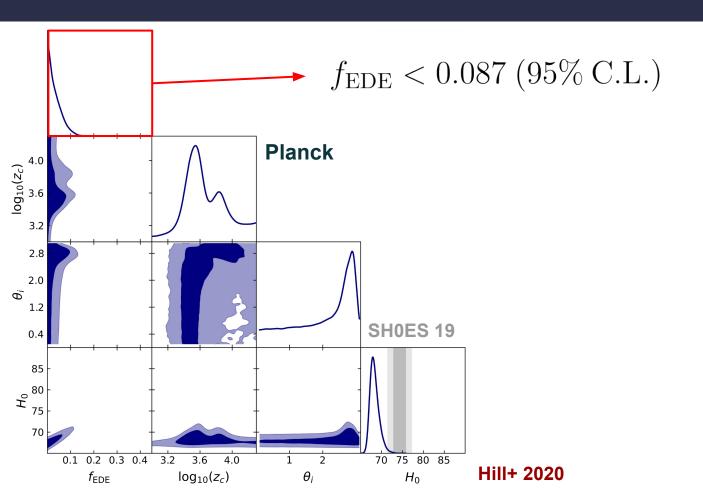




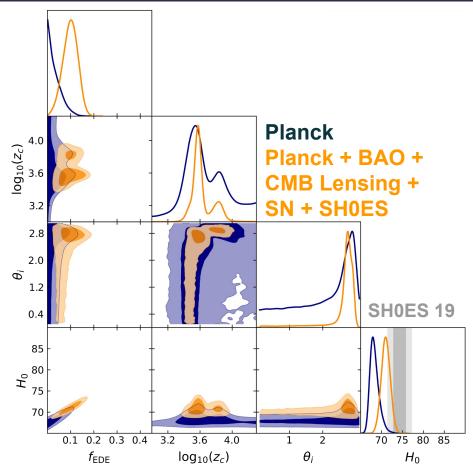
#### **Constraints on EDE from Planck data**



#### **Constraints on EDE from Planck data**

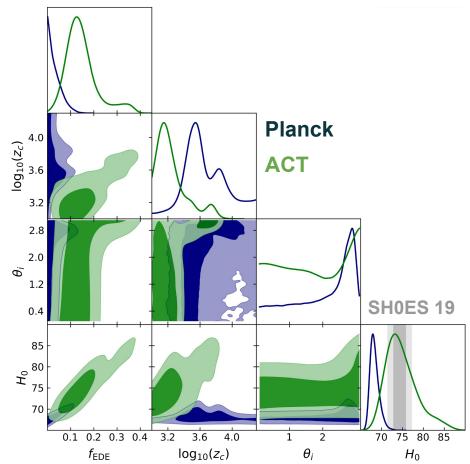


#### Results for a combination of Planck and SH0ES data

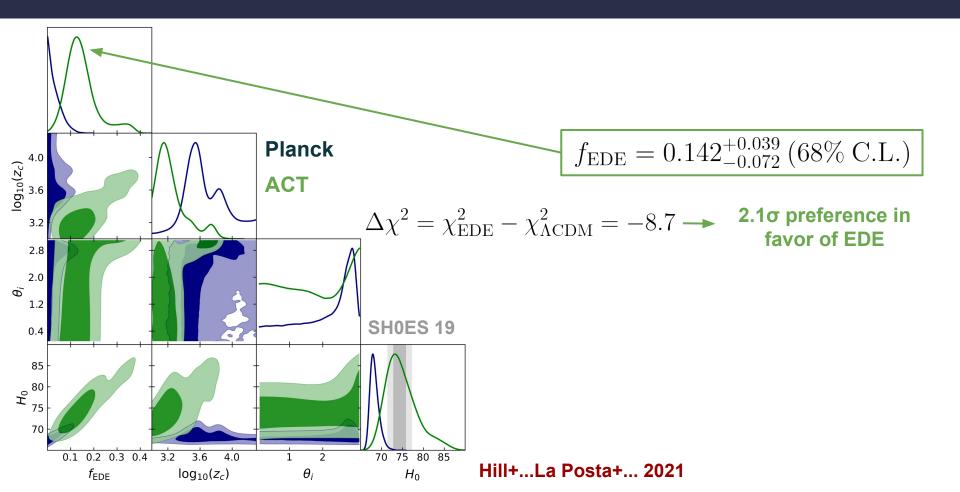


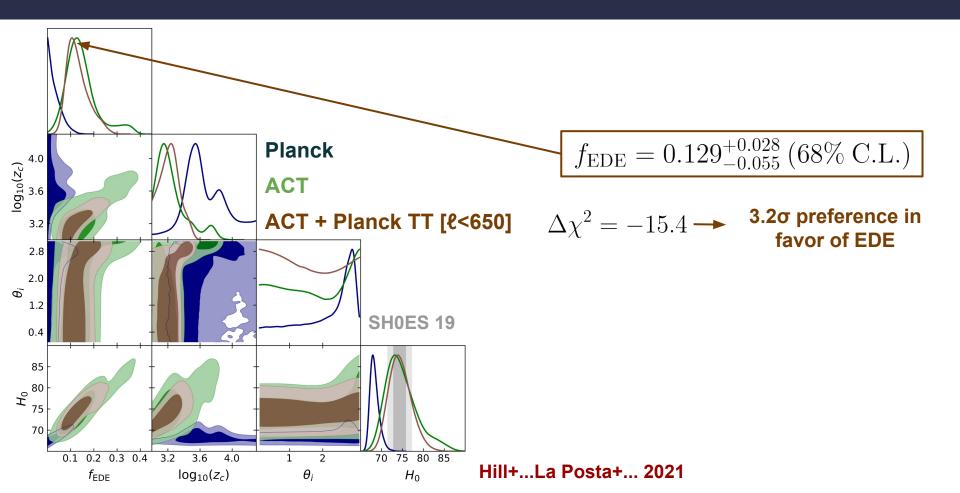
Poulin+ 2019, Smith+ 2019, Hill+ 2020

#### Additional constraints from ACTPol



Hill+...La Posta+... 2021





### Summary of EDE results

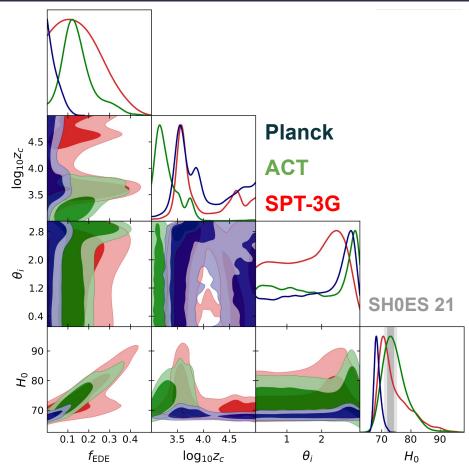
- Planck data alone don't favor high f<sub>EDE</sub> values (Hill+ 2020)
- Planck data in combination with SH0ES show a preference for non-zero f<sub>EDF</sub> (Poulin+ 2019, Smith+ 2019)
- ACT data alone favors EDE over ΛCDM (Hill+...La Posta+... 2021)

### **Summary of EDE results**

- Planck data alone don't favor high f<sub>EDE</sub> values (Hill+ 2020)
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- ACT data alone favors EDE over ΛCDM (Hill+...La Posta+... 2021)

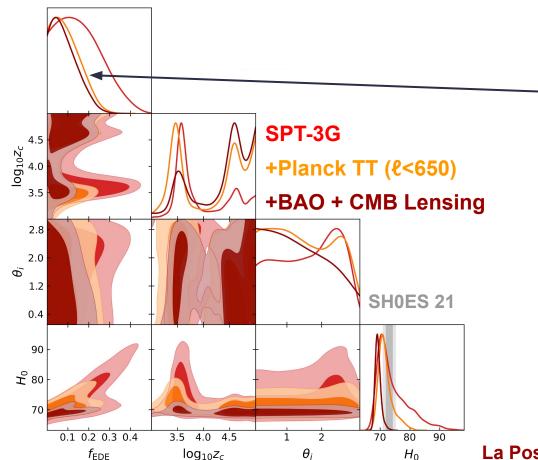
→ Motivates an analysis of EDE with public SPT-3G data

# **Constraints from SPT-3G public data**



La Posta+ 2021 [arXiv:2112.10754]

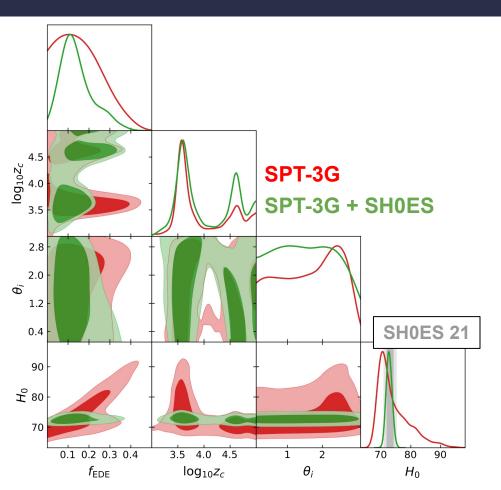
### **Constraints from SPT-3G public data**



We tighten the constraint on  $f_{EDE}$  when we combine SPT3G and Planck TT ( $\ell$ <650) or when we add LSS probes

La Posta+ 2021 [arXiv:2112.10754]

# **Combining with SH0ES constraint**

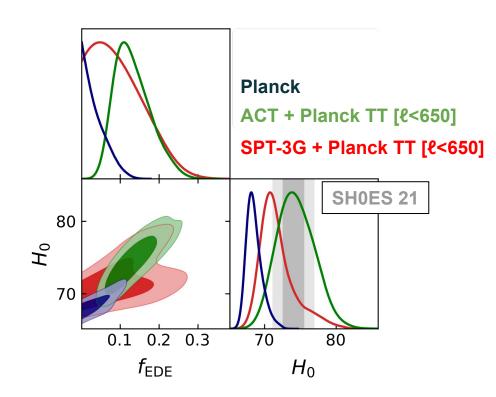


$$\Delta\chi^2_{\rm SPT-3G} = -6.3$$

improvement of the fit to SPT-3G data (with respect to ΛCDM)

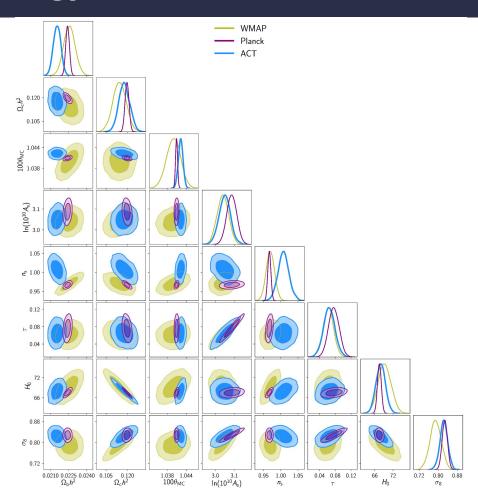
#### Conclusions

- Planck data alone do not favor high f<sub>EDE</sub> values
- Planck + SH0ES show a preference for f<sub>EDE</sub> ~ 10%
- ACT DR4 data favors EDE over ΛCDM (with f<sub>EDE</sub>~ 10%)
- SPT-3G is not as constraining as ACT and Planck : but sees some degree of EDE when combined with SH0ES

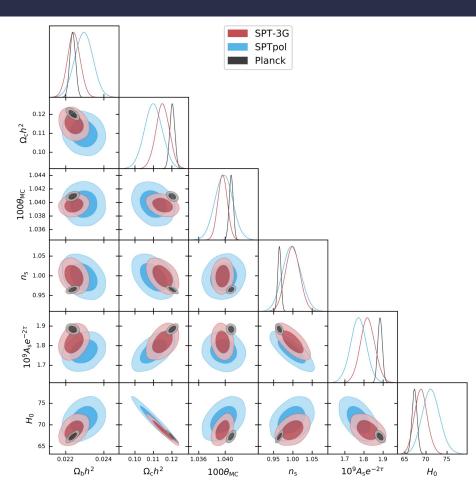


#### **Extra-slides**

# **ACTPol** cosmology

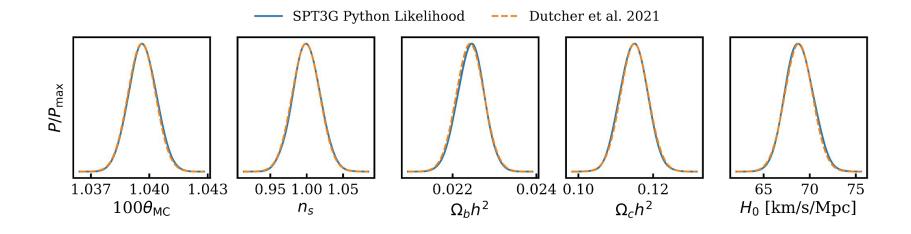


# **SPT-3G cosmology**

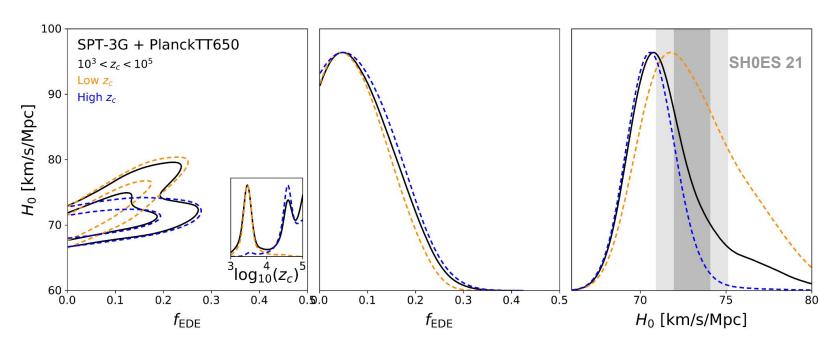


**Dutcher+ 2021** 

# SPT-3G python implementation

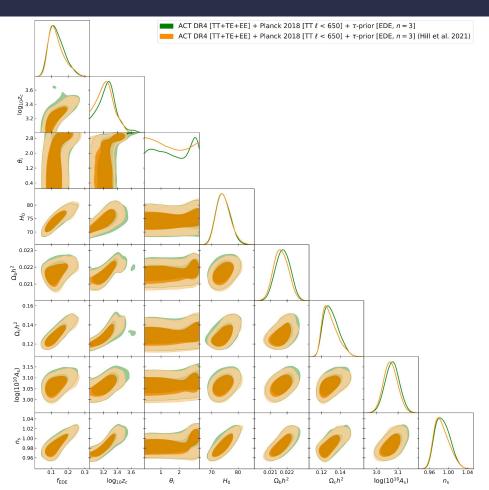


# Impact of the z<sub>c</sub> prior



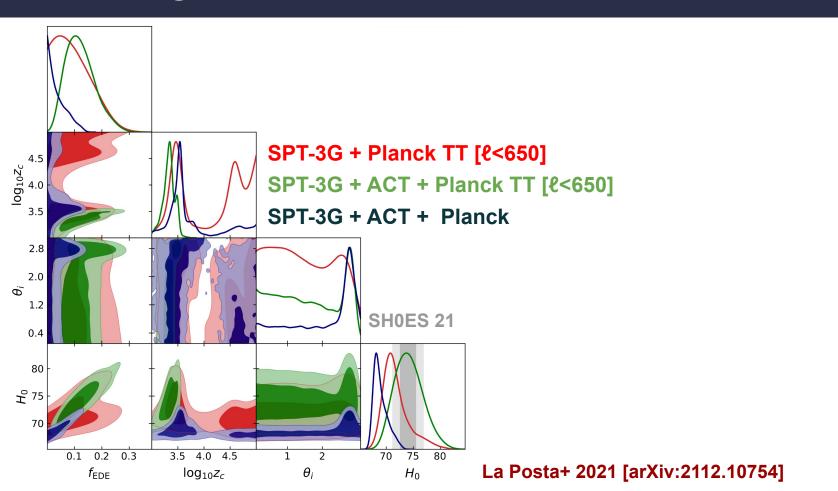
La Posta+ 2021 [arXiv:2112.10754]

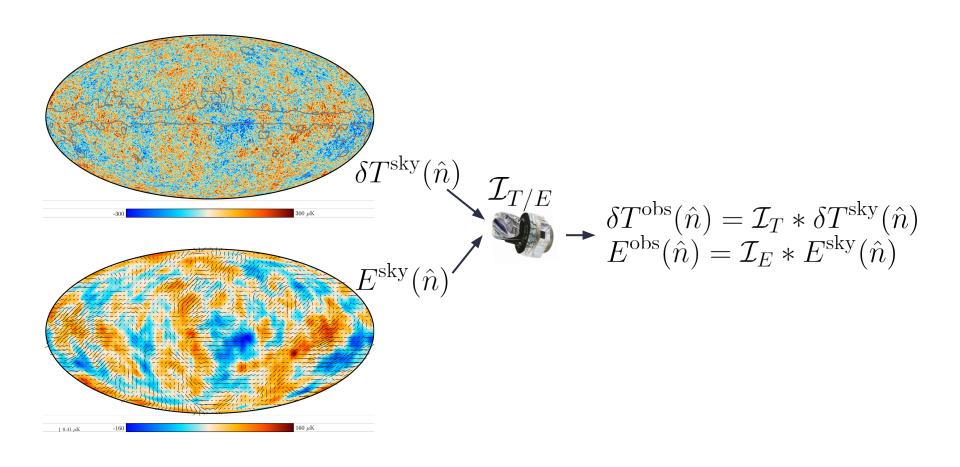
#### **CAMB/CLASS EDE models**



Hill+...La Posta+... 2021

### **Combining with other CMB datasets**





$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$
$$\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$$

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

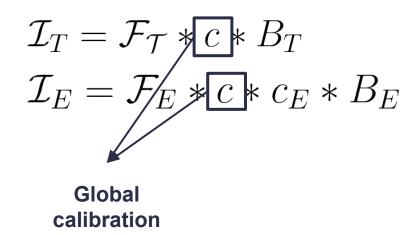
• Finite angular resolution (beams)

$$\mathcal{I}_T = \mathcal{F}_{\mathcal{T}} * c * B_T$$
 $\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$ 

Temperature (Polarization) beam

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration



$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
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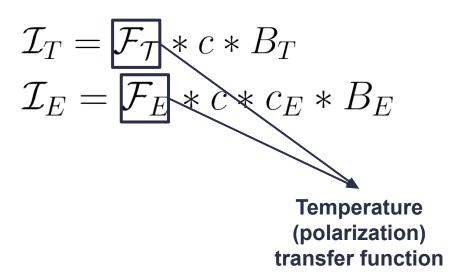
- Finite angular resolution (beams)
- Calibration
- Polarization efficiency

$$\mathcal{I}_T = \mathcal{F}_{\mathcal{T}} * c * B_T$$
 $\mathcal{I}_E = \mathcal{F}_E * c * \overline{c_E} * B_E$ 

Polarization efficiency

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)



$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
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- Polarization efficiency
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# These instrumental effects are multiplicative in harmonic space

$$C_{\ell}^{TT,\text{obs}} = (\mathcal{F}_{\ell}^{T})^{2} c^{2} (B_{\ell}^{T})^{2} C_{\ell}^{TT}$$

$$C_{\ell}^{EE,\text{obs}} = (\mathcal{F}_{\ell}^{E})^{2} c^{2} c_{E}^{2} (B_{\ell}^{E})^{2} C_{\ell}^{EE}$$

$$C_{\ell}^{TE,\text{obs}} = \mathcal{F}_{\ell}^{T} \mathcal{F}_{\ell}^{E} c^{2} c_{E} B_{\ell}^{T} B_{\ell}^{E} C_{\ell}^{EE}$$

#### **Correlation coefficient of T and E modes**

$$\mathcal{R}_{\ell}^{TE} = \frac{\left\langle a_{\ell m}^{T} a_{\ell m}^{E*} \right\rangle}{\sqrt{\left\langle a_{\ell m}^{T} a_{\ell m}^{T*} \right\rangle \left\langle a_{\ell m}^{E} a_{\ell m}^{E*} \right\rangle}} = \frac{C_{\ell}^{TE}}{\sqrt{C_{\ell}^{TT} C_{\ell}^{EE}}}$$

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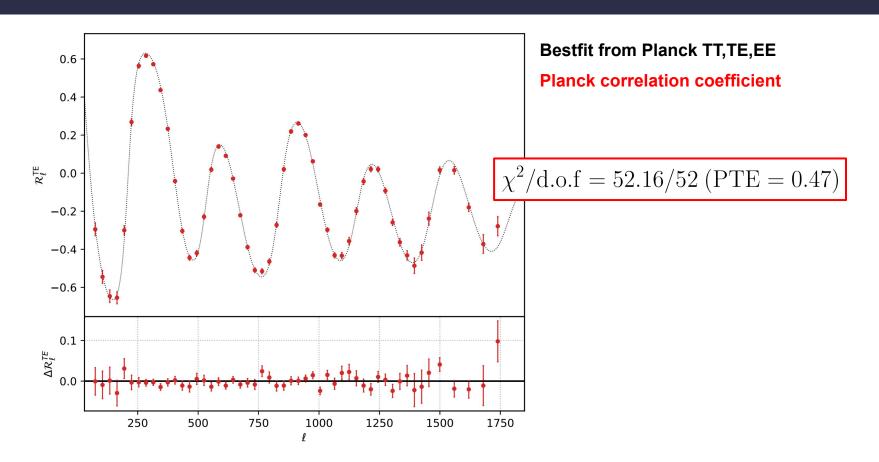
$$\mathcal{R}_{\ell}^{TE, \text{obs}} = \frac{\mathcal{F}_{\ell}^{T} \mathcal{F}_{\ell}^{E} c^{2} c_{E} B_{\ell}^{T} B_{\ell}^{E} C_{\ell}^{TE}}{\sqrt{(\mathcal{F}_{\ell}^{T})^{2} c^{2} (B_{\ell}^{T})^{2} C_{\ell}^{TT} \times (\mathcal{F}_{\ell}^{E})^{2} c^{2} c_{E}^{2} (B_{\ell}^{E})^{2} C_{\ell}^{EE}}}$$

#### Correlation coefficient of T and E modes

$$\mathcal{R}_{\ell}^{TE} = \frac{\left\langle a_{\ell m}^{T} a_{\ell m}^{E*} \right\rangle}{\sqrt{\left\langle a_{\ell m}^{T} a_{\ell m}^{T*} \right\rangle \left\langle a_{\ell m}^{E} a_{\ell m}^{E*} \right\rangle}} = \frac{C_{\ell}^{TE}}{\sqrt{C_{\ell}^{TT} C_{\ell}^{EE}}}$$

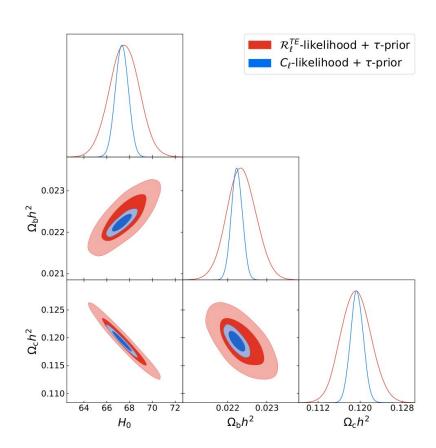
$$\mathcal{R}_{\ell}^{TE, \text{obs}} = \frac{\mathcal{F}_{\ell}^{T} \mathcal{F}_{\ell}^{E} c^{2} c_{E} B_{\ell}^{T} B_{\ell}^{E} C_{\ell}^{TE}}{\sqrt{(\mathcal{F}_{\ell}^{T})^{2} c^{2} (B_{\ell}^{T})^{2} C_{\ell}^{TT} \times (\mathcal{F}_{\ell}^{E})^{2} c^{2} c_{E}^{2} (B_{\ell}^{E})^{2} C_{\ell}^{EE}}} = \mathcal{R}_{\ell}^{TE}$$

#### Planck correlation coefficient



La Posta+ 2021 [Phys. Rev. D 104, 023527]

### Cosmological results from R<sup>TE</sup>

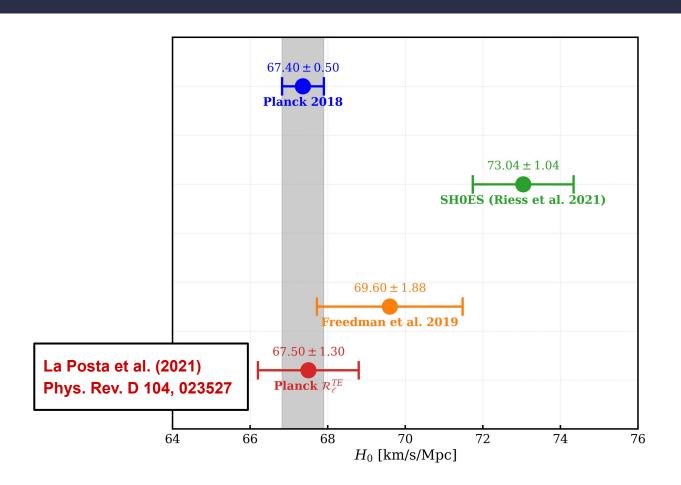


**3.3σ** away from the latest SH0ES measurement

 $H_0 = 67.5 + /- 1.3 [km/s/Mpc]$ 

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#### **Hubble tension**



# Independent measurements of H<sub>0</sub> from the ground

**Atacama Cosmology Telescope** 

6m telescope in the Atacama desert (Chile ~5000m high)

**ACT DR4 (Choi+ 2020, Aiola+ 2020)** 

data collected from 2013 to 2016

Cosmological analysis on ~5400 deg<sup>2</sup>
observed at 98 and 150 GHz

**South Pole Telescope** 

10m primary mirror (South Pole ~2800m high)

SPT-3G results (Dutcher+ 2021)

4 month period in 2018

observed at 95, 150 and 220 GHz

Cosmological analysis on ~1500 deg<sup>2</sup>