

# Beam Spin Asymmetry for Deeply Virtual $\pi^0$ Electroproduction with CLAS12

International workshop on CLAS12 physics  
and future perspectives at JLab

March 22, 2023

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(University of Connecticut)

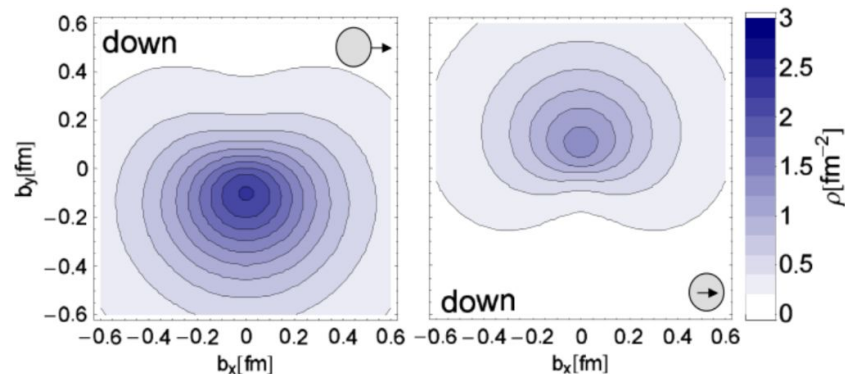
		Quark polarization		
		U	L	T
Nucleon polarization	U	$H$		$\bar{E}_T$
	L		$\tilde{H}$	
	T	$E$		$H_T, \tilde{H}_T$

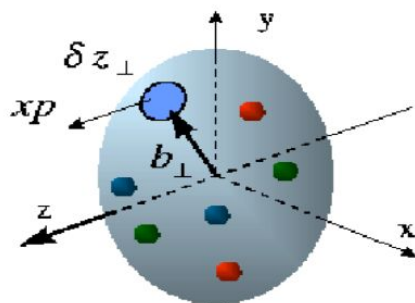
Chiral-odd GPD results:

- Deeply virtual meson production
- Lattice QCD by Gockeler *et al*

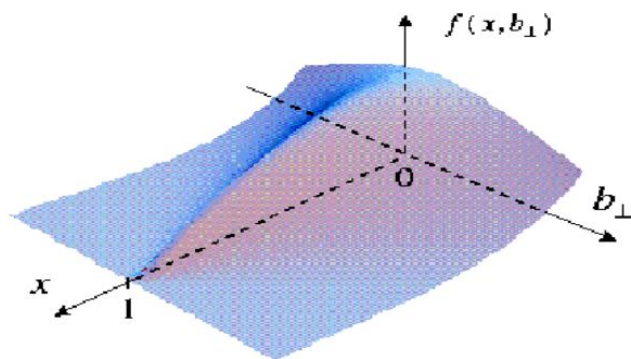
Chiral even GPDs:

- DVCS on unpolarized and polarized targets with polarized beam by HERMES, JLAB and COMPASS





- 4 chiral-even GPDs:  $H$ ,  $E$ ,  $\tilde{H}$ ,  $\tilde{E}$
- 4 chiral-odd GPDs:  $H_T$ ,  $E_T$ ,  $\tilde{H}_T$ ,  $\tilde{E}_T$



Meson	GPD flavor composition
$\tilde{H}, \tilde{E}$	$\pi^+$ $\Delta u - \Delta d$
$H_T, \bar{E}_T$	$\pi^0$ $2\Delta u + \Delta d$
	$\eta$ $2\Delta u - \Delta d$
$H, E$	$\rho^0$ $2u + d$
	$\rho^+$ $u - d$
	$\omega$ $2u - d$

PDFs:

in the forward limit

$$\xi = t = 0:$$

$$H^q(x, 0, 0) = q(x)$$

$$\tilde{H}^q(x, 0, 0) = \Delta q(x)$$

Form Factors:

$$\int dx H^q(x, \xi, t) = F_1(t)$$

$$\int dx E^q(x, \xi, t) = F_2(t)$$

$$\int dx \tilde{H}^q(x, \xi, t) = G_A(t)$$

$$\int dx \tilde{E}^q(x, \xi, t) = G_P(t)$$

X. Ji, Phys. Rev. Lett. 78, 610 (1997):

$$J^q = \int x dx [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$

- Proton anomalous tensor magnetic moment

$$\kappa_T^u = \int dx \bar{E}_T^u(x, \xi, t=0)$$

$$\kappa_T^d = \int dx \bar{E}_T^d(x, \xi, t=0)$$

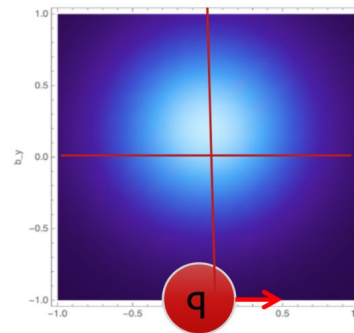
- Proton tensor charge

$$\delta_T^u = \int dx H_T^u(x, \xi, t=0)$$

$$\delta_T^d = \int dx H_T^d(x, \xi, t=0)$$

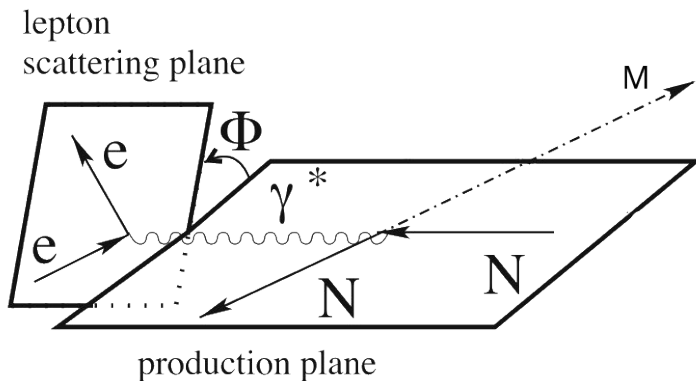
- Density of transversity polarized quarks in an unpolarized proton in the transverse plane

$$\delta(x, \vec{b}) = \frac{1}{2} \left[ H(x, \vec{b}) - \frac{b_y}{m} \frac{\partial}{\partial b^2} \bar{E}_T(x, \vec{b}) \right]$$

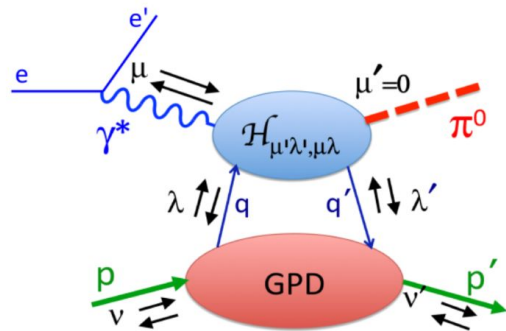


$$\frac{d^4\sigma}{dQ^2 dx_B dt d\Phi} = \Gamma(Q^2, x_B, E) \frac{1}{2\pi} \left\{ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\Phi) + \sqrt{\epsilon(2\epsilon + 1)} \frac{d\sigma_{LT}}{dt} \cos(\Phi) + \lambda \sqrt{2\epsilon(1 - \epsilon)} \frac{d\sigma_{LT'}}{dt} \sin(\Phi) \right\}$$

where  $\lambda$  is the helicity state of the incident electron beam



$$\sigma = \sigma_0 + \sqrt{2\epsilon(1+\epsilon)}\sigma_{LT}^{\cos\phi} \cos\phi + \epsilon\sigma_{TT}^{\cos 2\phi} \cos 2\phi + \lambda_e\sqrt{2\epsilon(1-\epsilon)}\sigma_{LT'}^{\sin\phi} \sin\phi$$



PHYSICAL REVIEW D 84, 034007 (2011)

## Flexible parametrization of generalized parton distributions from deeply virtual Compton scattering observables

Gary R. Goldstein,<sup>1,\*</sup> J. Osvaldo Gonzalez Hernandez,<sup>2,†</sup> and Simonetta Liuti<sup>2,‡</sup>

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(Received 16 February 2011; published 5 August 2011)

Eur. Phys. J. A (2011) 47: 112  
DOI 10.1140/epja/2011-11112-6

Regular Article – Theoretical Physics

## Transversity in hard exclusive electroproduction of pseudoscalar mesons

S.V. Goloskokov<sup>1,\*</sup> and P. Kroll<sup>2,3,§</sup>

THE EUROPEAN  
PHYSICAL JOURNAL A

$$\langle F \rangle = \sum_{\lambda} \int_{-1}^1 dx \mathcal{H}_{0\lambda, \mu\lambda} (x, \xi, Q^2, t) F(x, \xi, t)$$

## Goloskokov-Kroll model:

$$\sigma_L \sim \left\{ (1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

$$\sigma_T \sim \left[ (1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\sigma_{LT} \sim \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

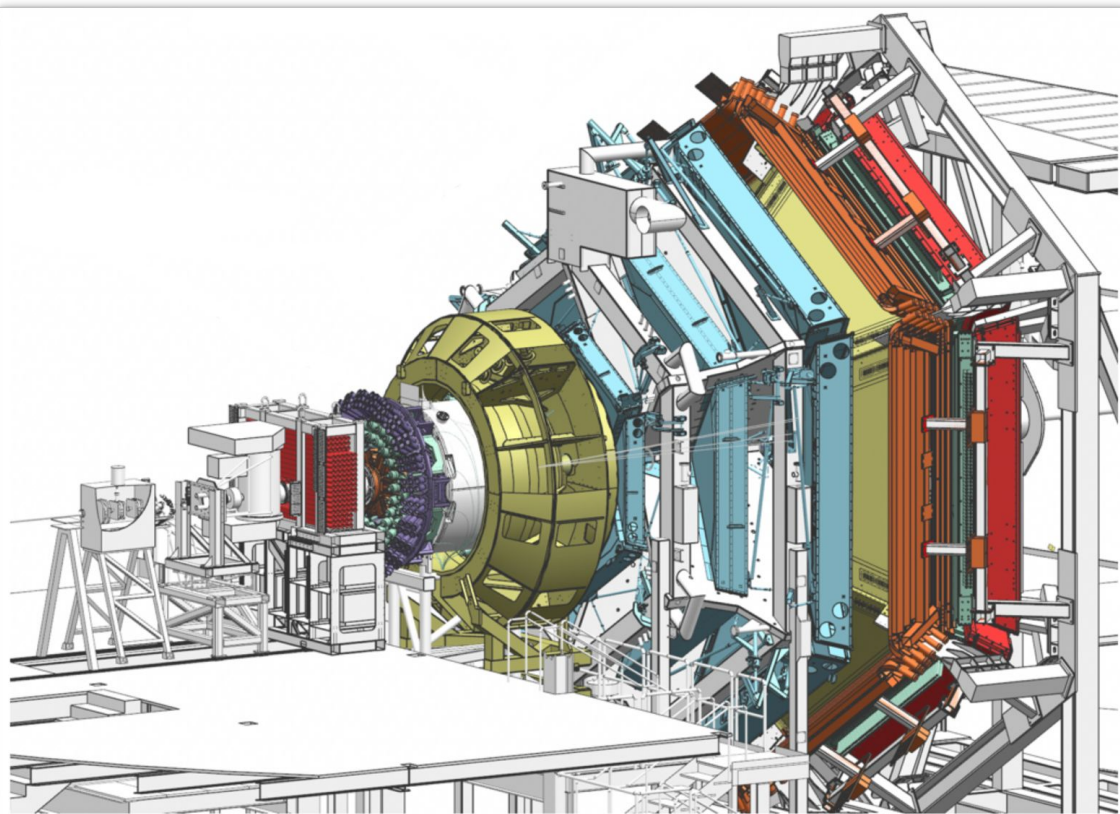
$$\sigma_{TT} \sim \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

$$\sigma_{LT'} \sim \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \text{Im} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$



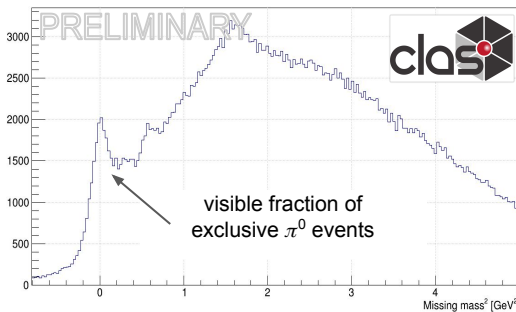
- CEBAF Large Acceptance Spectrometer
- First CLAS experiment since 12 GeV upgrade
- 10.6 GeV longitudinally polarized electron beam
- 86 % electron beam polarization
- Unpolarized liquid hydrogen fixed target
- Access to the  $Q^2$  up to 10 GeV<sup>2</sup>

$$ep \rightarrow e+p+\pi^0 \rightarrow e+p+\gamma+\gamma$$



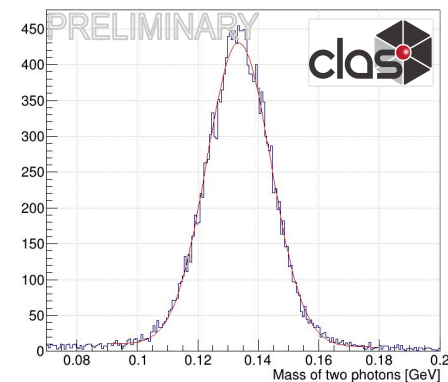


Raw  $MM^2(epX)$  distribution

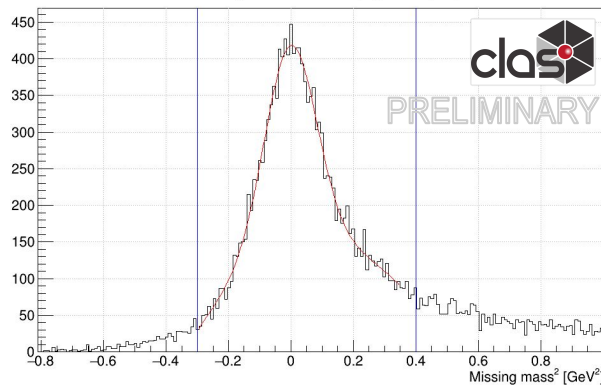


- All final state particle detected:  
electron, proton, two photons
- Reconstructed mass of two photons is used to identify  $\pi^0$  candidates
- The momentum and energy conservation laws provide powerful constraints to suppress background: the plots show clear exclusive sample after application of exclusivity cuts

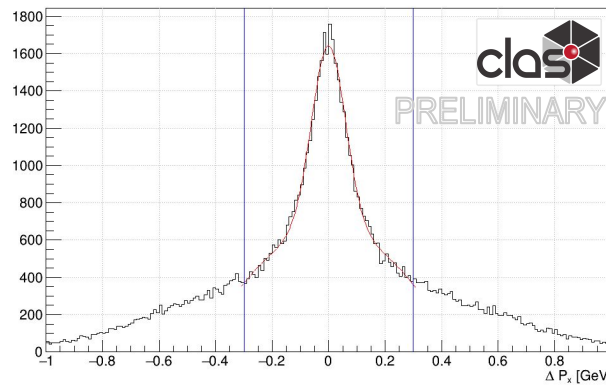
Mass of two photons



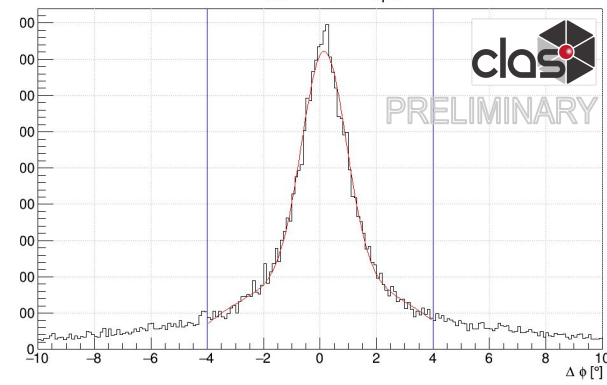
Missing mass<sup>2</sup> for ( $ep \rightarrow epX$ )

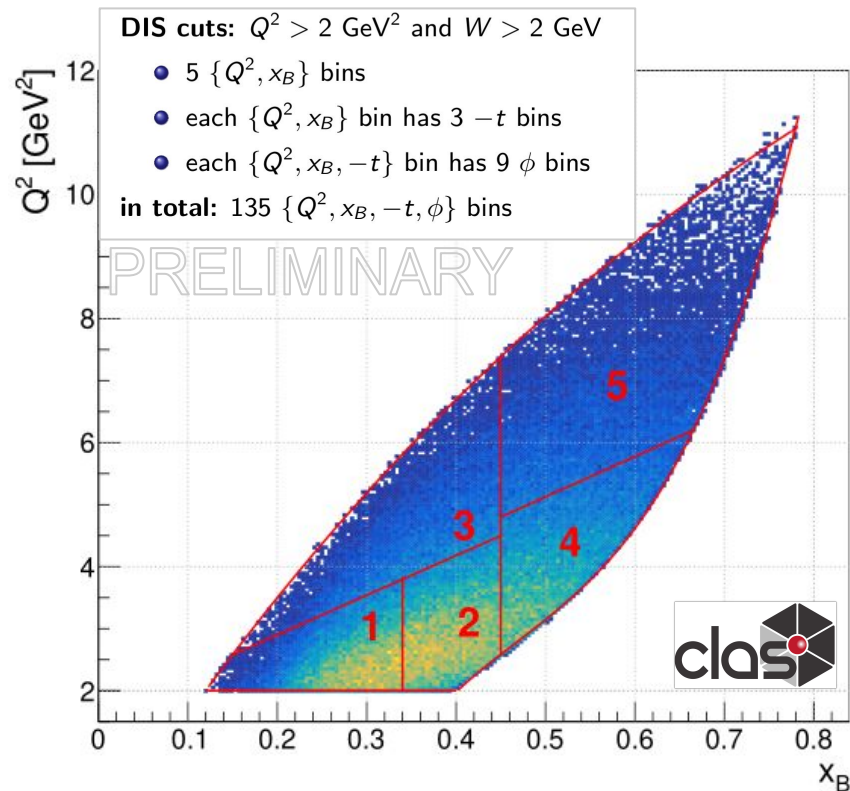
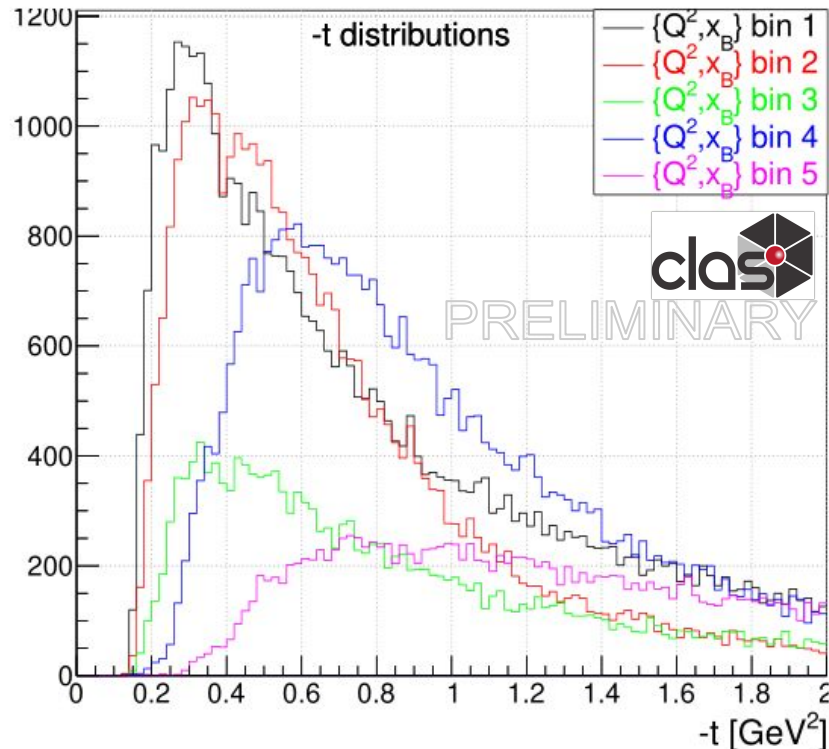


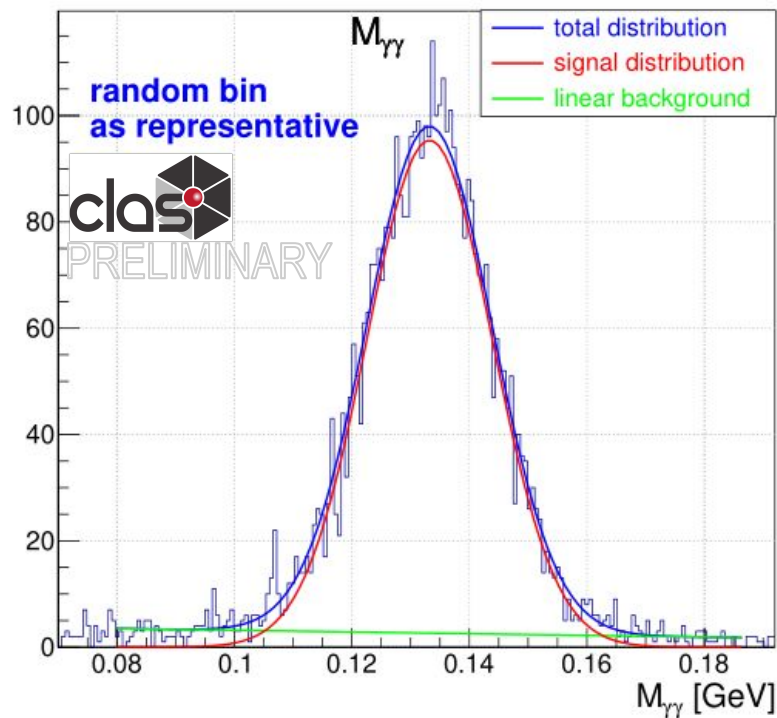
$\Delta P_x$  for ( $ep \rightarrow ep\pi^0X$ )



$\Delta \phi_{\pi X} = \phi_{\pi} - \phi_{epX}$

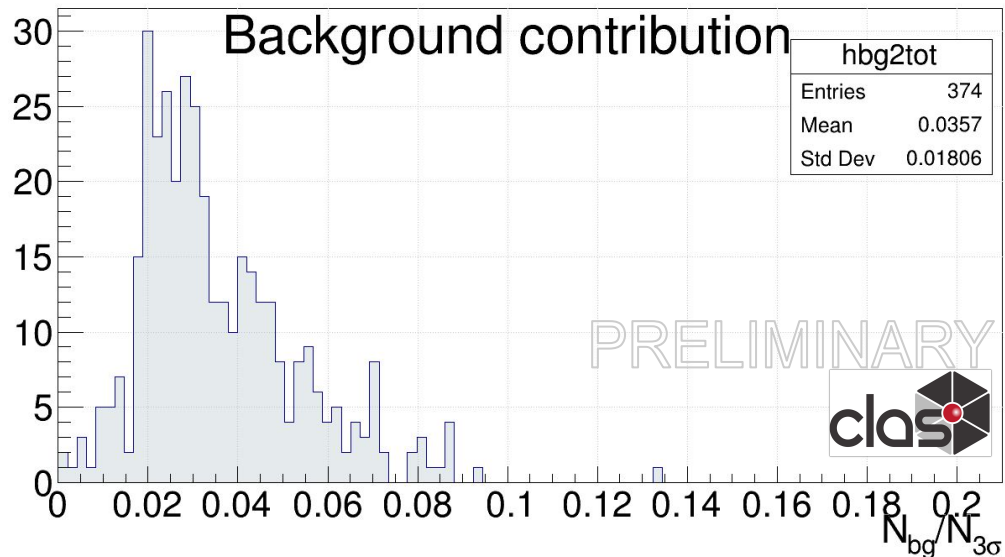


(a)  $Q^2$  vs  $x_B$  coverage(b)  $-t$  distributions for each  $\{Q^2, x_B\}$  bin



Sideband background subtraction:

$$N_{signal} = N_{3\sigma} - N_{bg}$$



$$BSA = \frac{1}{P_b} \frac{n^+ - n^-}{n^+ + n^-},$$

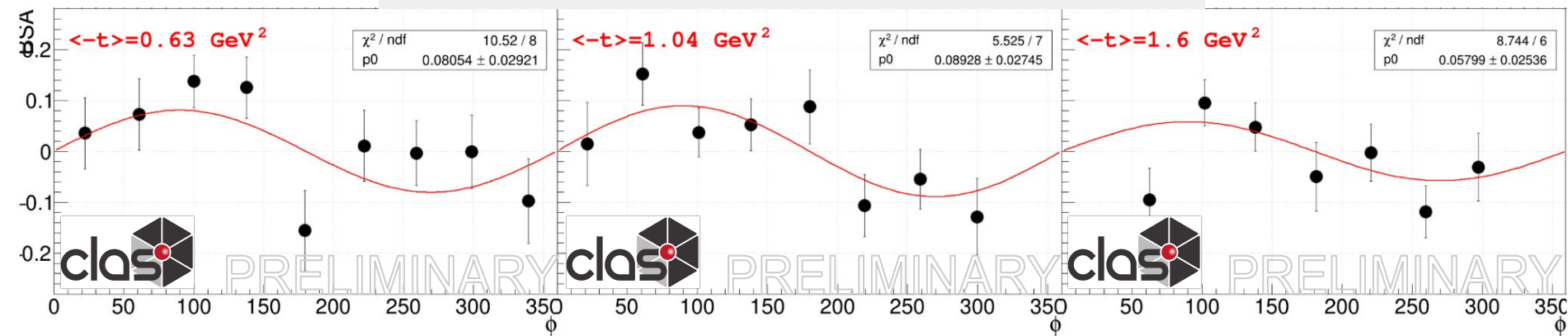
where  $P_b$  is an average electron beam polarization

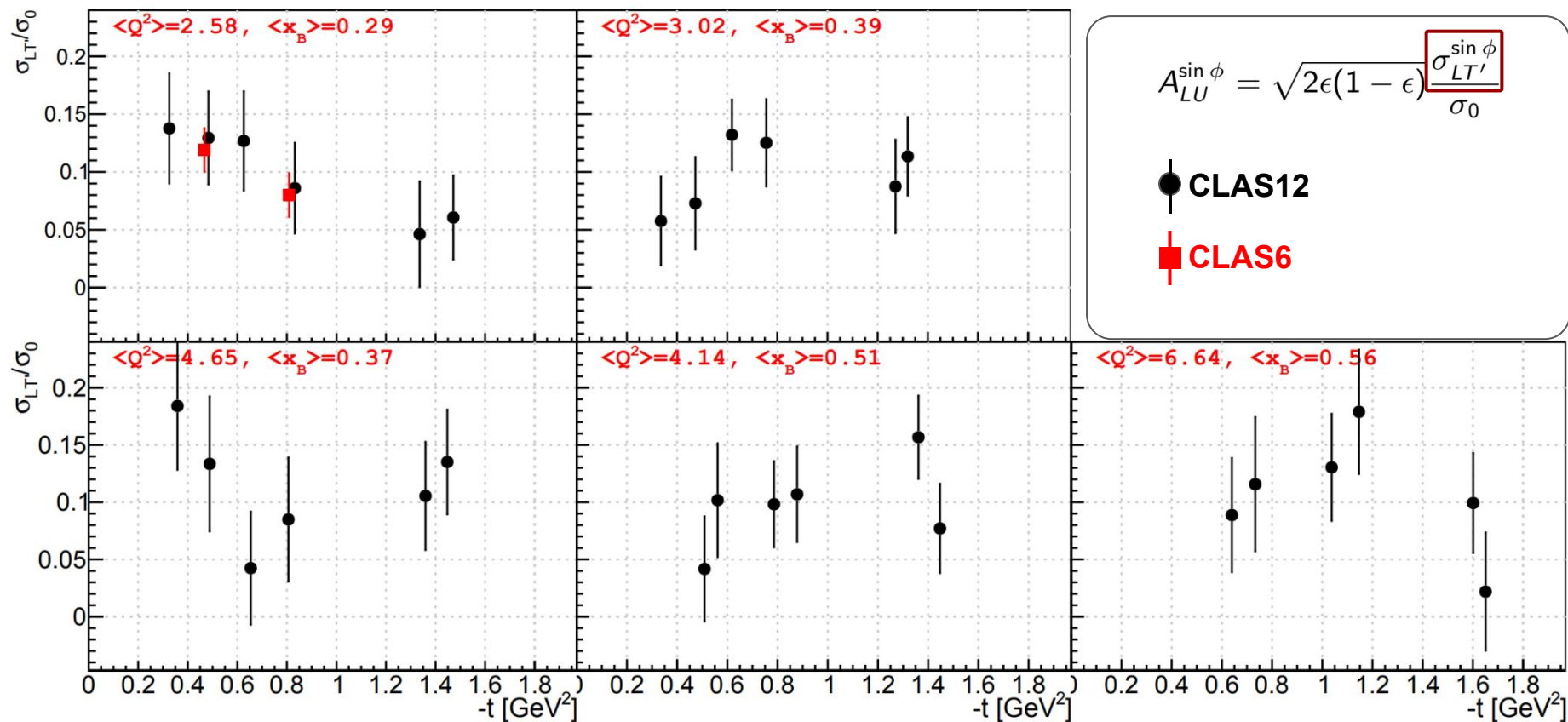
$$\sigma = \sigma_0 + \sqrt{2\epsilon(1+\epsilon)} \sigma_{LT}^{\cos \phi} \cos \phi + \epsilon \sigma_{TT}^{\cos 2\phi} \cos 2\phi + \lambda_e \sqrt{2\epsilon(1-\epsilon)} \sigma_{LT'}^{\sin \phi} \sin \phi$$

$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} \propto A_{LU}^{\sin \phi} \sin \phi$$

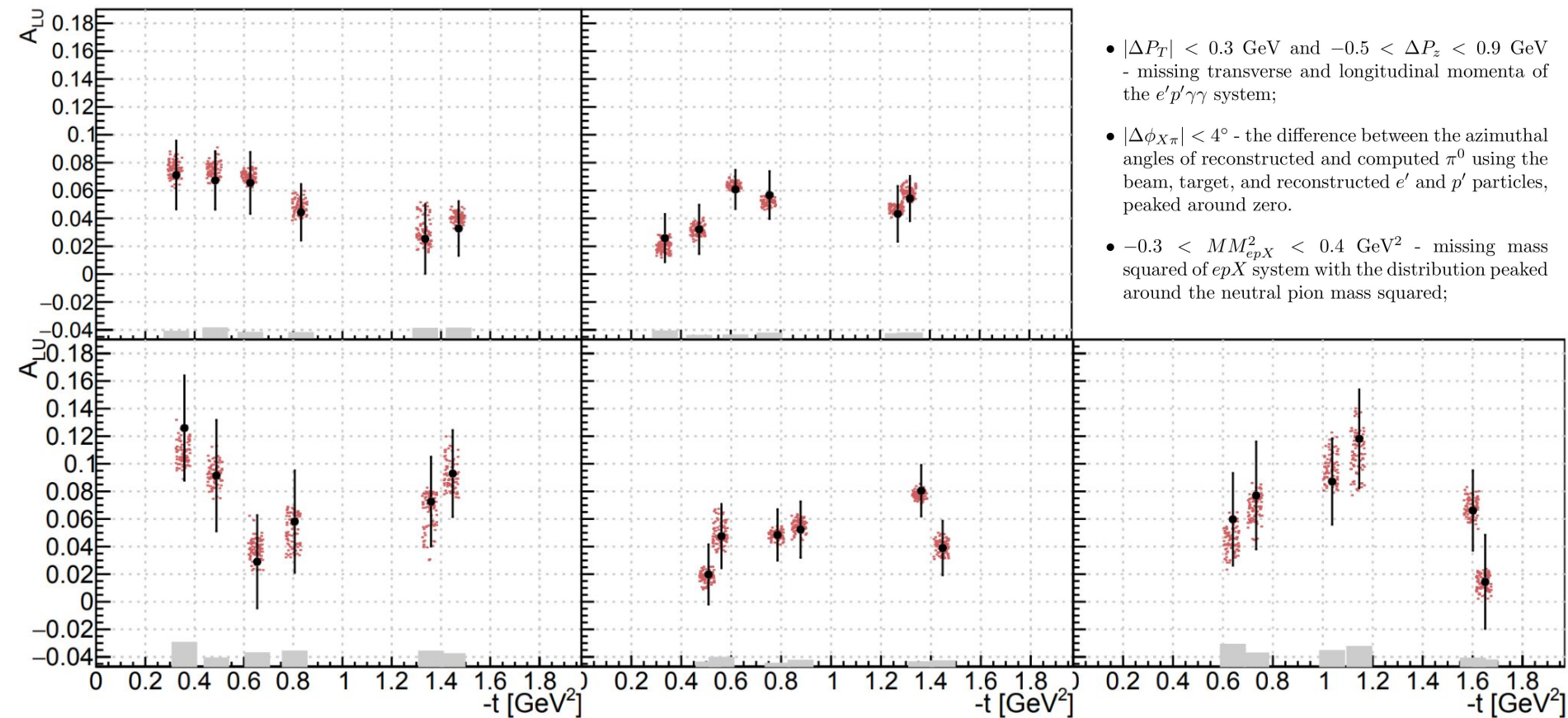
$$A_{LU}^{\sin \phi} = \sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}^{\sin \phi}}{\sigma_0}$$

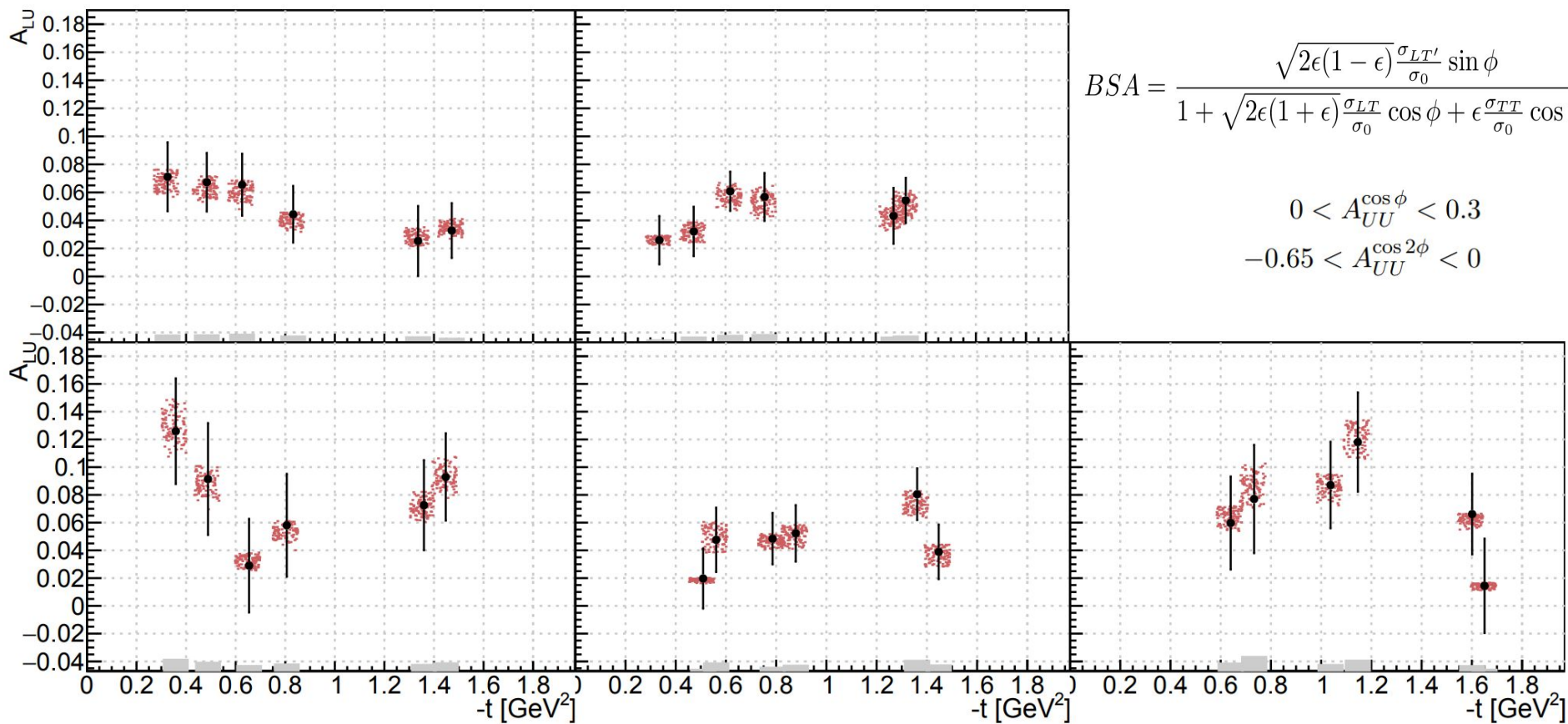
Beam spin asymmetries for 3  $\langle -t \rangle$  bins in the 5th  $\langle Q^2, x_B \rangle$  bin



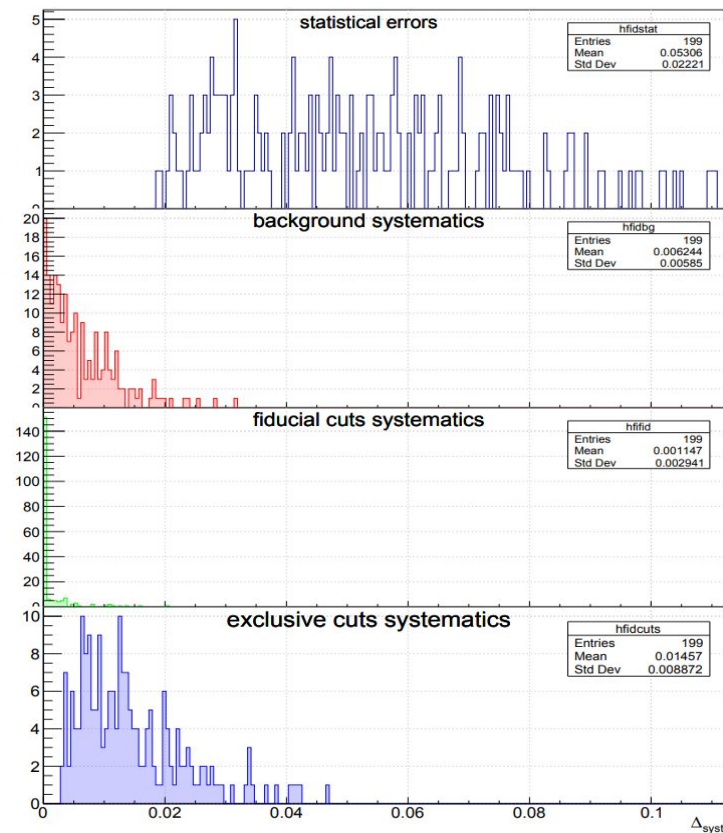
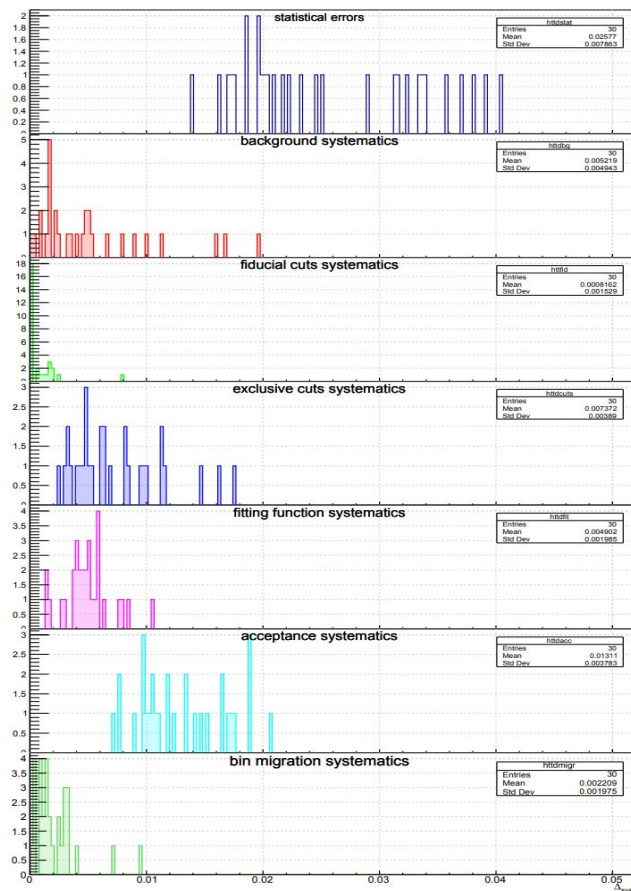


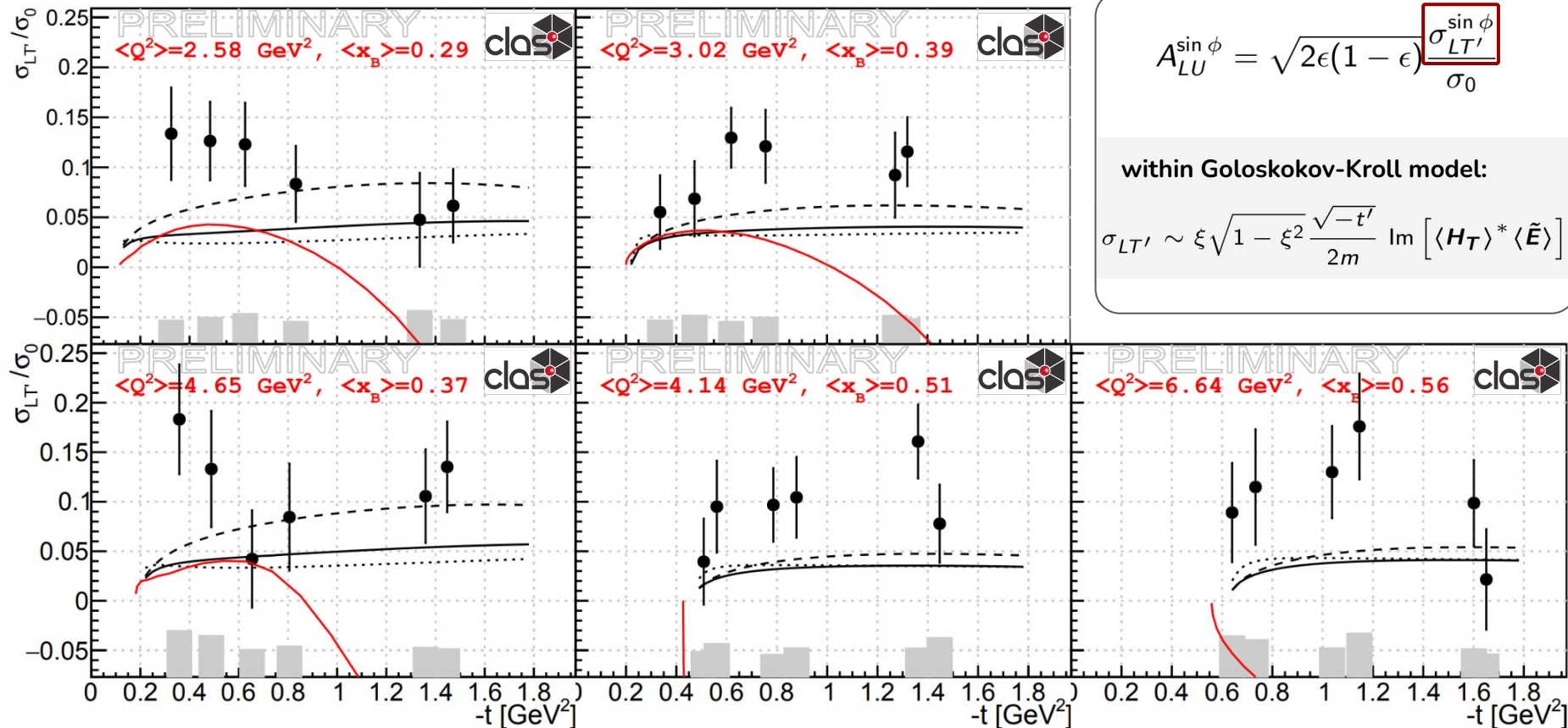








(b) for  $A_{LU}$  BSA measurements



$$A_{LU}^{\sin \phi} = \sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}^{\sin \phi}}{\sigma_0}$$

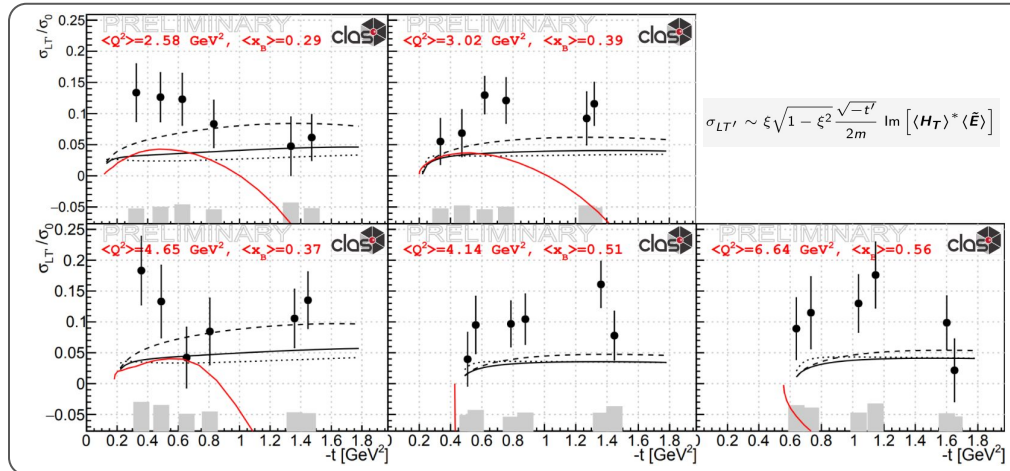
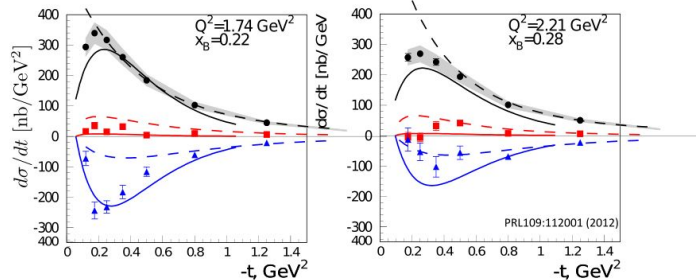
within Goloskokov-Kroll model:

$$\sigma_{LT'} \sim \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \text{Im} \left[ \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

## Unpolarized structure functions

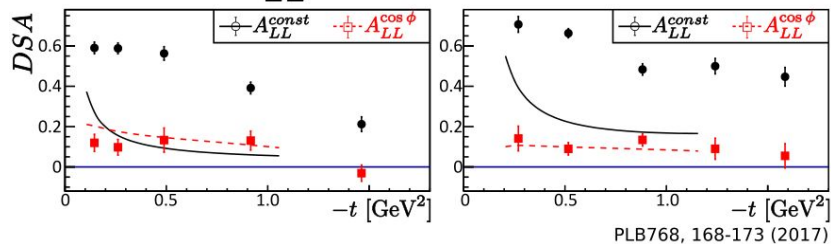
$$\sigma_T \sim (1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2$$

$$\sigma_{TT} \sim \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

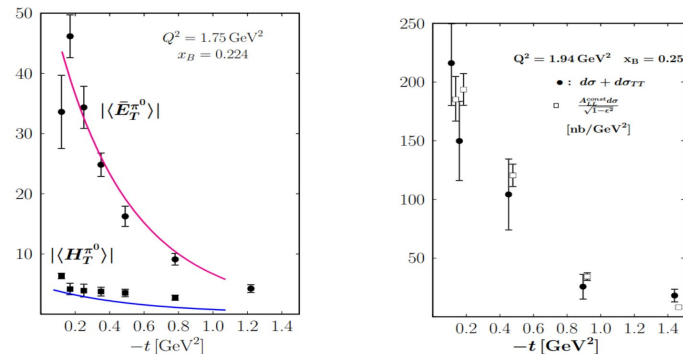


## Double spin asymmetry

$$A_{LL}^{const} \sim |\langle H_T \rangle|^2$$

 $H_T$  is underestimated in GK model

$$A_{LL}^{const} \sim |\langle H_T \rangle|^2$$



Few-Body Syst 57, 1041-1050 (2016)

## Measurement of Exclusive $\pi^0$ Electroproduction Structure Functions and their Relationship to Transverse Generalized Parton Distributions

I. Bedlinskiy *et al.* (CLAS Collaboration)  
Phys. Rev. Lett. **109**, 112001 – Published 10 September 2012

## Rosenbluth Separation of the $\pi^0$ Electroproduction Cross Section

M. Defurne *et al.* (Jefferson Lab Hall A Collaboration)  
Phys. Rev. Lett. **117**, 262001 – Published 23 December 2016

## Exclusive $\eta$ electroproduction at $W > 2$ GeV with CLAS and transversity generalized parton distributions

I. Bedlinskiy *et al.* (CLAS Collaboration)  
Phys. Rev. C **95**, 035202 – Published 10 March 2017

## Rosenbluth Separation of the $\pi^0$ Electroproduction Cross Section Off the Neutron

M. Mazouz *et al.* (Jefferson Lab Hall A Collaboration)  
Phys. Rev. Lett. **118**, 222002 – Published 1 June 2017



A multidimensional study of the structure function ratio  $\sigma_{LT'}/\sigma_0$  from hard exclusive  $\pi^+$  electro-production off protons in the GPD regime

S. Diehl<sup>a,h,\*</sup>, A. Kim<sup>f</sup>, K. Joo<sup>f</sup>



## Measurement of $ep \rightarrow ep\pi^0$ beam spin asymmetries above the resonance region

R. De Masi *et al.* (CLAS Collaboration)  
Phys. Rev. C **77**, 042201(R) – Published 14 April 2008



Physics Letters B  
Volume 789, 10 February 2019, Pages 426-431



## Measurement of the beam spin asymmetry of $\vec{e}p \rightarrow e'p'\eta$ in the deep-inelastic regime with CLAS

B. Zhao<sup>a,b</sup>, A. Kim<sup>a</sup>, K. Joo<sup>a</sup>, I. Bedlinskiy<sup>c</sup>, W. Kim<sup>d</sup>, V. Kubarovsky<sup>e,f</sup>, M. Ungaro<sup>g</sup>,



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Target and double spin asymmetries of deeply virtual  $\pi^0$  production with a longitudinally polarized proton target and CLAS

A. Kim<sup>b,a,\*</sup>, H. Avakian<sup>c</sup>, V. Burkert<sup>c</sup>, K. Joo<sup>a</sup>, W. Kim<sup>b</sup>



## 1 Beam Spin Asymmetry Measurements of Deeply Virtual $\pi^0$ Production with CLAS12

A. Kim,<sup>1</sup> K. Joo,<sup>1</sup> and S. Diehl<sup>2,1</sup>  
(The CLAS Collaboration)

<sup>1</sup>University of Connecticut, Storrs, Connecticut 06269

<sup>2</sup>II. Physikalisches Institut der Universität Gießen, 35392 Gießen, Germany

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5

- The preliminary Beam Spin Asymmetry for deeply virtual exclusive  $\pi^0$  electroproduction was measured with CLAS12 and longitudinally polarized electron beam
- New CLAS12 measurements extend experimental reach up to 10 GeV<sup>2</sup>, twice the range compared to CLAS6 data
- Sizable extracted BSA moments indicate significant contributions from chiral-odd GPDs
- Chiral-odd based theoretical models will benefit from additional constraints provided by new measurements in a wide kinematic range