# Impact of Neutron DVCS Measurements on Extraction of Compton Form Factors

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Modelling

# Outline

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# Deeply Virtual Compton Scattering

exclusive process, measured via leptoproduction of a photon



• interference with the Bethe-Heitler process gives unique access to both real and imaginary parts of the DVCS amplitude

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# Accessing GPDs

- at leading order four complex twist-two Compton form factors  $\mathcal{H}(\xi, t, Q^2)$ ,  $\mathcal{E}(\xi, t, Q^2)$ ,  $\widetilde{\mathcal{H}}(\xi, t, Q^2)$ ,  $\widetilde{\mathcal{E}}(\xi, t, Q^2)$
- factorization theorem [Collins et al. '98]



• CFFs are a convolution [Müller '92, et al. '94, Ji, Radyushkin '96]

$${}^{a}\mathcal{H}\left(\xi,t,Q^{2}\right) = \int \mathrm{d}x \ C^{a}\left(x,\xi,\frac{Q^{2}}{Q_{0}^{2}}\right) \underbrace{H^{a}\left(x,\eta=\xi,t,Q_{0}^{2}\right)}_{\mathsf{GPD}}, \ a=q,G$$

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# Types of models

- ① "Physical" GPD (and CFF) model
- ② Neural network parametrization of CFFs

# **Modelling GPDs**

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# GPD evolution

 evolution in x space complicated, we introduce conformal moments

$$F_n(\eta, t) = \int_{-1}^1 dx c_n(x, \eta) F(x, \eta, t)$$
$$c_n(x, \eta) = \eta^n \frac{\Gamma\left(\frac{3}{2}\right) \Gamma(1+n)}{2^n \Gamma\left(\frac{3}{2}+n\right)} C_n^{\frac{3}{2}}\left(\frac{x}{\eta}\right)$$

•  $C_n^{3/2}$  Gegenbauer polynomials

- analytic continuation  $n 
  ightarrow j \in \mathbb{C}$
- evolution diagonal in j space at LO

$$\mu \frac{d}{d\mu} F_j^q \left(\eta, t, \mu^2\right) = -\frac{\alpha_s(\mu)}{2\pi} \gamma_j^{(0)} F_j^q \left(\eta, t^2, \mu^2\right)$$

# Hybrid model

• valence quarks modelled in x space (q = u, d) at crossover line  $x = \eta$  (no  $Q^2$  evolution)

$$\Im\mathfrak{m}\mathcal{H}(\xi,t) \stackrel{LO}{=} \pi \left[ \frac{4}{9} H^{u_{\mathrm{val}}}(\xi,\xi,t) + \frac{1}{9} H^{d_{\mathrm{val}}}(\xi,\xi,t) + \frac{2}{9} H^{\mathrm{sea}}\left(\xi,\xi,t\right) \right]$$

- sea quarks modelled in j space
- SO(3) partial waves expansion
- leading contribution

$$H_{j}^{a}(\eta = 0, t) = N^{a} \frac{\mathbf{B} \left(1 - \alpha^{a} + j, \beta^{a} + 1\right)}{\mathbf{B} \left(2 - \alpha^{a}, \beta^{a} + 1\right)} \frac{\beta(t)}{1 - \frac{t}{\left(m_{j}^{a}\right)^{2}}},$$

$$(m_j^a)^2 = \frac{1+j-\alpha^a}{\alpha'^a}, \quad \beta(t) = \left(1 - \frac{t}{M^2}\right)^{-p}, \quad a = \{s, g\}$$

• full NLO QCD  $Q^2$  evolution

# **Dispersion relations**

CFFs constrained by dispersion relations

$$\mathfrak{Re}\,\mathcal{H}(\xi,t) \stackrel{LO}{=} \Delta(t) + \frac{1}{\pi} \mathrm{P.V.} \int_0^1 \mathrm{d}x \left(\frac{1}{\xi-x} - \frac{1}{\xi+x}\right) \Im\mathfrak{m}\,\mathcal{H}(x,t)$$

- only imaginary part of CFFs and one subtraction constant  $\Delta(t)$  are modelled

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# Neural networks constrained by dispersion relations



• Only imaginary part of CFFs and one subtraction constant  $\Delta(t)$  are parametrized by neural nets

# **Results**

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#### Proton DVCS

## Extraction of 6 CFFs

## [M. Č., K. Kumerički, A. Schäfer, '20], from JLab Hall A data





#### [Benali et al. '20], DVCS off a deuterium target



Using isospin symmetry (e.g.  $H_{u,\text{proton}}^{\text{val}} = H_{d,\text{neutron}}^{\text{val}}$ ) we combine proton and neutron DVCS data to separate up and down quark contributions to CFFs.

- Flavor separation
  - separate model for each flavor CFF:  $\mathcal{H}_u$ ,  $\mathcal{H}_d$
  - fKM20 "physical" flavored model, fNNDR neural nets and dispersion relations



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Flavor CFFs		

 $\mathcal{H} = \frac{4}{2}\mathcal{H}_u + \frac{1}{2}\mathcal{H}_d$ 

 ${\ }^{\bullet}$  up and down contributions to CFF  ${\mathcal H}$  cleanly separated

$$\begin{array}{c} \mathbf{g} \quad \mathbf{g} \\ \mathbf{$$

$$x_B = 0.36$$
$$Q^2 = 4 \text{ GeV}^2$$

Flavor separation

### • $\mathcal{E}$ cannot be separated



# CLAS 12 GeV predictions

proton and neutron beam spin asymmetry

$$A_{LU} = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}} \propto \Im \mathfrak{m} \left\{ F_1 \mathcal{H} + \xi \left( F_1 + F_2 \right) \widetilde{\mathcal{H}} - \frac{\Delta^2}{4M^2} F_2 \mathcal{E} \right\} \sin(\phi)$$

- we analyse harmonics with beam energy 10.4 GeV
- physical model only assumes isospin rotation

$$\mathcal{H}_n^{\text{val}} = \frac{2e_d^2 + e_u^2}{2e_u^2 + e_d^2} \mathcal{H}^{\text{val}} = \frac{2}{3} \mathcal{H}^{\text{val}} , \quad \mathcal{H}_n^{\text{sea}} = \mathcal{H}^{\text{sea}}$$

# 2020 models predictions







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Reweighting		

- reweighting neural nets according to their  $\chi^2$
- for the 2020 paper we had 20 neural nets in each model. Now we generated additional 80 nets (same training procedure, same old data) to "improve statistics"
- most  $\chi^2$  were still too large, formal reweighting yields just one "best" net
- hand selecting the best nets still yields too few nets
- it seems that reweighting procedure is appropriate for the situation where new data is just a refinement of the old data, and not when completely new observables are being measured

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#### CLAS22 predictions

# New fits with CLAS 2022 data

• We performed new NN and NNDR fits on harmonics with new CLAS 2022 data (39 points with  $-t < 0.5 \text{ GeV}^2$ ) and previously available JLab data (257 points)

	2020	2023
fNN	1.5	1.25
fNNDR	1.5	>3

Table:  $\chi^2/N_{pts}$ 

- new data clearly exclude the fNNDR model
- we also trained NN and NNDR models on only CLAS data (excluding Hall A data): fNNC23 and fNNDRC23 models
- fNNC23 behaves better than fNNDRC23, just like fNN23 outperforms fNNDR23, but fNN23 still outperforms all of them

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#### CLAS22 predictions

# Proton data







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# New CFF extraction



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# Flavor separation of CFF ${\cal H}$



# Flavor separation of CFF $\mathcal{E}$



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- new data favours models without dispersion relations
- 2020 data allows for flavour separation of CFF  $\mathcal{H}$ , but not  $\mathcal{E}$
- 2022 data allows for flavour separation of  $\mathfrak{Im}\,\mathcal{H}$  and  $\mathfrak{Re}\,\mathcal{E}$