

*LAL Cours d'automne 2009*

# Potential Discoveries at the Large Hadron Collider

Chris Quigg

*Fermilab*

# Lecture 3: Standard Model Opportunities

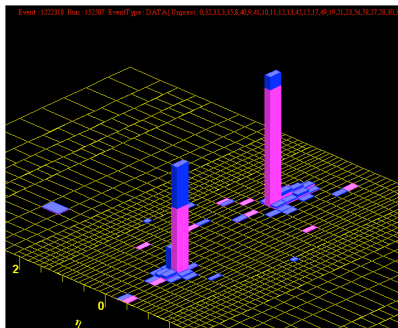
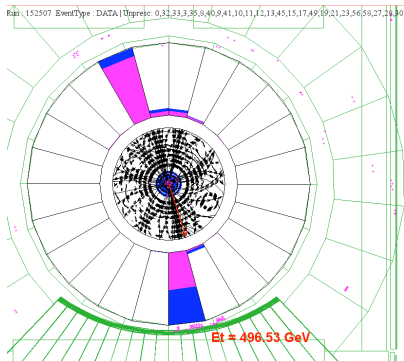
## Themes of Contemporary Particle Physics

- Elementarity: *structureless quarks and leptons?*
- Symmetry: *EWSB and the 1-TeV scale; new forces*  
(origin of gauge symmetries; discrete symmetries)
- Unity: *quarks + leptons; gauge symmetries*  
(constituents + forces; incorporation of gravity)
- Identity: *fermion masses, mixings; CP violation*  
(Majorana  $\nu$ ?)
- Topography: *fabric of space and time*

# Elementarity

## The World's Most Powerful Microscopes

*nanonanophysics*

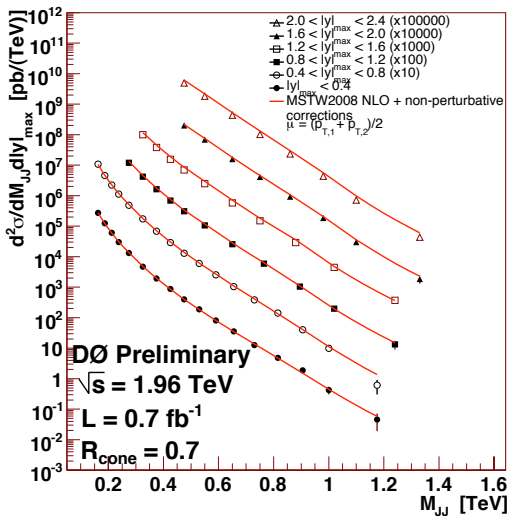


CDF dijet event ( $\sqrt{s} = 1.96$  TeV):  $E_T = 1.364$  TeV

$q\bar{q} \rightarrow \text{jet} + \text{jet}$

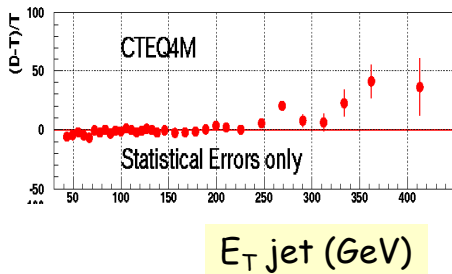
# Elementarity

## $\text{D}\bar{\text{D}}$ Jet Cross Sections



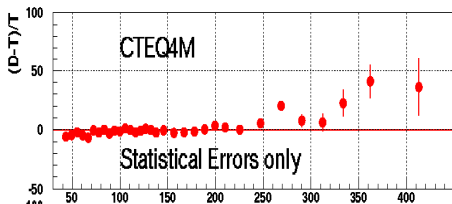
# Elementarity

## The CDF Experience

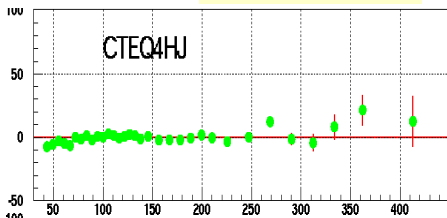


# Elementarity

## The CDF Experience



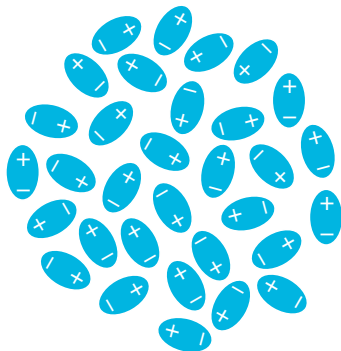
$E_T \text{ jet (GeV)}$



$E_T \text{ jet (GeV)}$

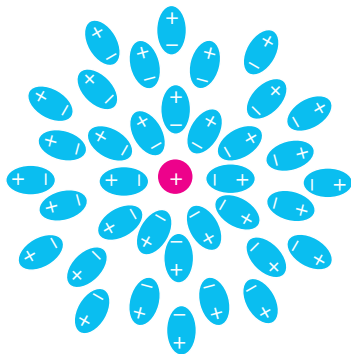
# Charge screening in electrodynamics

Dielectric (polarizable) medium ...



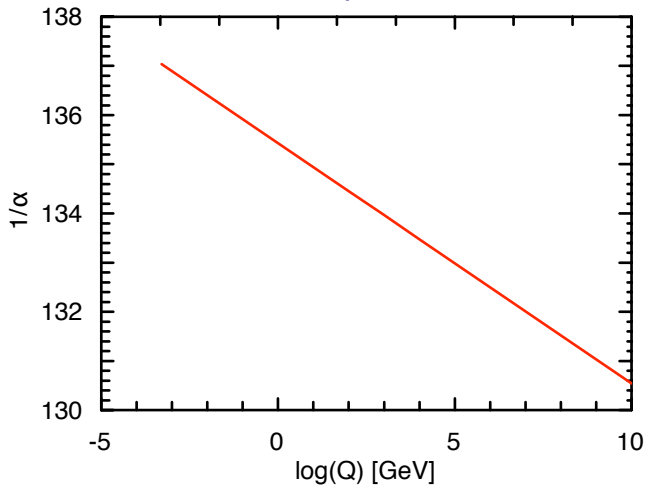
# Charge screening in electrodynamics

Dielectric (polarizable) medium ...



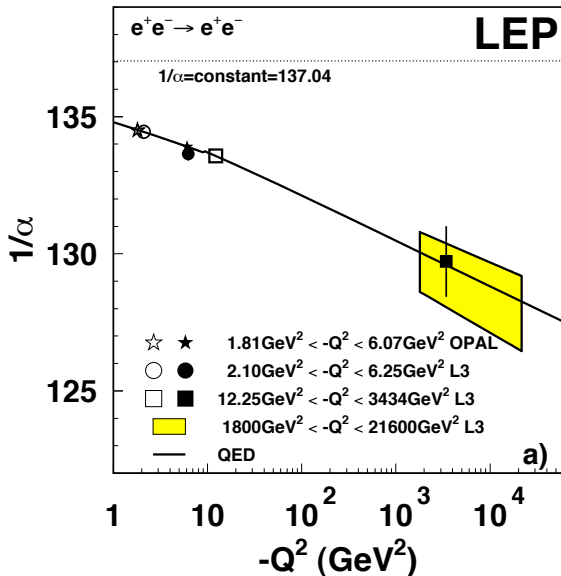


# Charge screening in QED (electrons + photons)

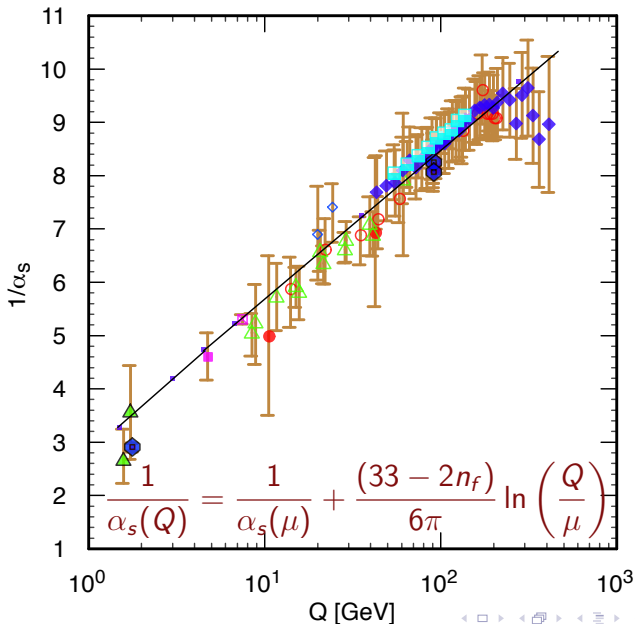


$$1/\alpha(Q) = 1/\alpha_0 - \frac{2}{3\pi} \ln\left(\frac{Q}{m}\right)$$

# Charge screening in QED (real world)

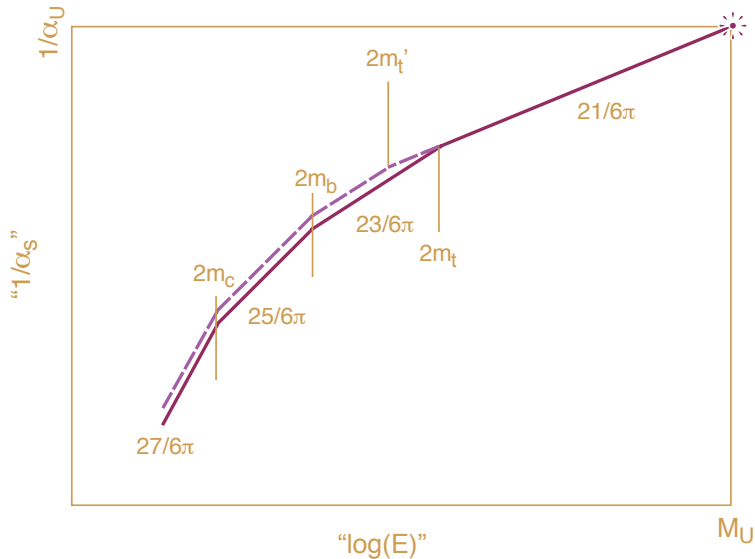


# Evolution of $\alpha_s(Q^2)$



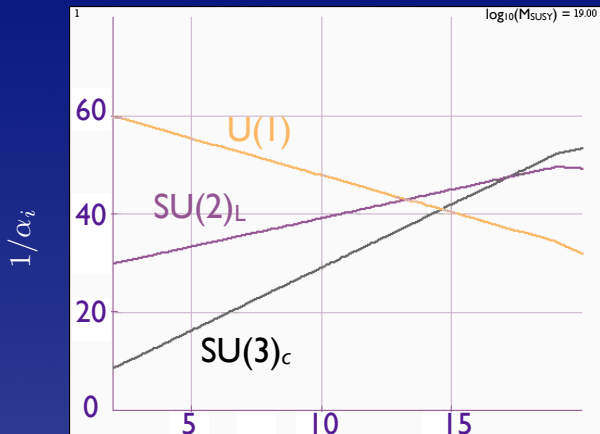
# Evolution of $\alpha_s(Q^2)$

Influence of  $m_t$



# Coupling Constant Unification

Different running of  $U(1)_Y$ ,  $SU(2)_L$ ,  $SU(3)_c$   
gives possibility of coupling constant unification

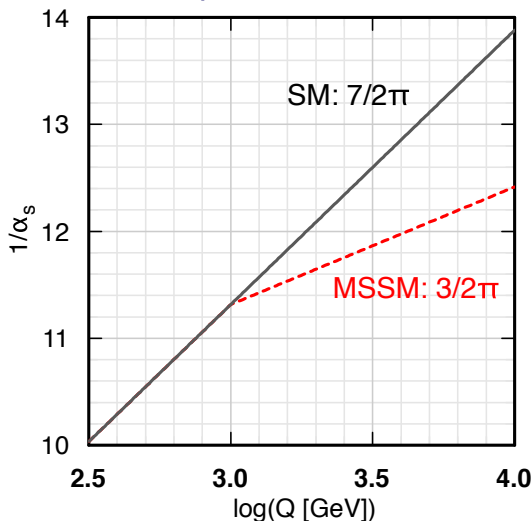


$$\alpha^{-1} = \frac{5}{3}\alpha_1^{-1} + \alpha_2^{-1}$$

$$\log_{10}(E[\text{GeV}])$$

# Can LHC See Change in Evolution?

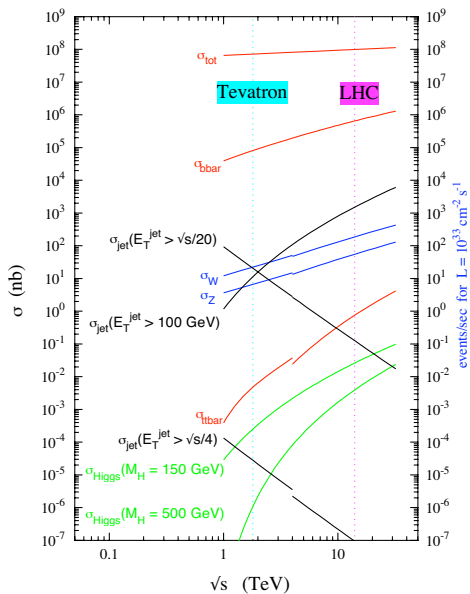
Sensitive to new colored particles



(sharp threshold illustrated)

... also for  $\sin^2 \theta_W$

# Sample event rates in $p^\pm p$ collisions



# Supermodels

## New physics possibilities in very early running

### Supermodels for early LHC

Christian W. Bauer,<sup>1,2</sup> Zoltan Ligeti,<sup>1,2</sup> Martin Schmaltz,<sup>1,2,3</sup> Jesse Thaler,<sup>1,2</sup> and Devin G. E. Walker<sup>1,2,4</sup>

<sup>1</sup>Theoretical Physics Group, Lawrence Berkeley National Laboratory, Berkeley, CA 94720

<sup>2</sup>Berkeley Center for Theoretical Physics, University of California, Berkeley, CA 94720

<sup>3</sup>Physics Department, Boston University, Boston, MA 02215

<sup>4</sup>Center for the Fundamental Laws of Nature, Jefferson Physical Laboratory, Harvard University, Cambridge, MA 02138

We investigate what new physics signatures the LHC can discover in the 2009–2010 run, beyond the expected sensitivity of the Tevatron data by 2010. We construct “supermodels”, for which the LHC sensitivity even with only  $10 \text{ pb}^{-1}$  is greater than that of the Tevatron with  $10 \text{ fb}^{-1}$ . The simplest supermodels involve  $s$ -channel resonances in the quark-antiquark and especially in the quark-quark channels. We concentrate on easily visible final states with small standard model backgrounds, and find that there are simple searches, besides those for  $Z'$  states, which could discover new physics in early LHC data. Many of these are well-suited to test searches for “more conventional” models, often discussed for multi- $\text{fb}^{-1}$  data sets.

## Rules of the game:

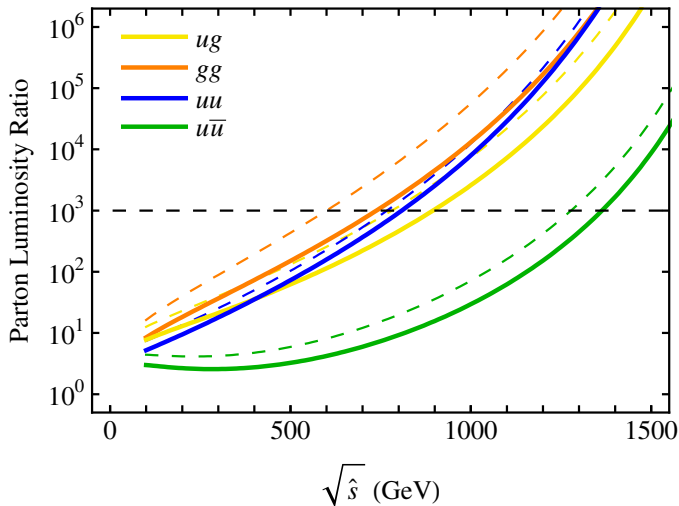
- $\gtrsim 10$  signal events in  $10 \text{ pb}^{-1}$  at LHC (25 better)
- No signal in  $10 \text{ fb}^{-1}$  at Tevatron
- Easily detected, low-background decay channel
- Consistent with existing constraints



# Supermodels

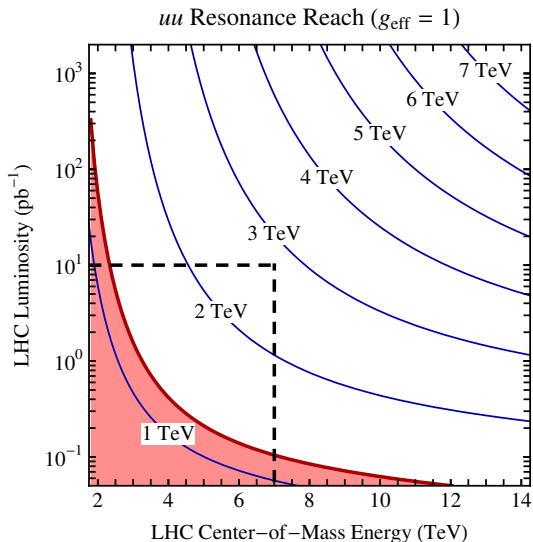
Need  $10^3$  advantage in parton luminosities

LHC (7 & 10 TeV) vs. Tevatron



# Supermodels

## Example: strongly coupled $qq$ resonance



# Supermodels

To observe diquark, require decays beyond  $qq$

An example: color-**6** diquark  $D$  + leptodiquark  $L$  (!)

$$\begin{array}{l} uu \rightarrow D \\ \quad \quad \quad \downarrow \\ \quad \quad \quad \ell^- L \\ \quad \quad \quad \quad \quad \downarrow \\ \quad \quad \quad \quad \quad \ell^+ jj \end{array}$$

*Doesn't respond to any needs, but ...*

final state familiar from  $W_R$  searches

Don't assume there is nothing to find at low  $\int \mathcal{L} dt$

# Electroweak Questions for the LHC. I

- What hides electroweak symmetry: a Higgs boson, or new strong dynamics?
- If a Higgs boson: one or several?
- Elementary or composite?
- Is the Higgs boson indeed light, as anticipated by the global fits to EW precision measurements?
- Does  $H$  only give masses to  $W^\pm$  and  $Z^0$ , or also to fermions? (Infer  $t\bar{t}H$  from production)
- Are the branching fractions for  $f\bar{f}$  decays in accord with the standard model?

If all this: what sets the fermion masses and mixings?

# Search for the Standard-Model Higgs Boson

$$\Gamma(H \rightarrow f\bar{f}) = \frac{G_F m_f^2 M_H}{4\pi\sqrt{2}} \cdot N_c \cdot \left(1 - \frac{4m_f^2}{M_H^2}\right)^{3/2}$$

$\propto M_H$  in the limit of large Higgs mass;  $\propto \beta^3$  for scalar

$$\Gamma(H \rightarrow W^+W^-) = \frac{G_F M_H^3}{32\pi\sqrt{2}} (1-x)^{1/2} (4-4x+3x^2) \quad x \equiv 4M_W^2/M_H^2$$

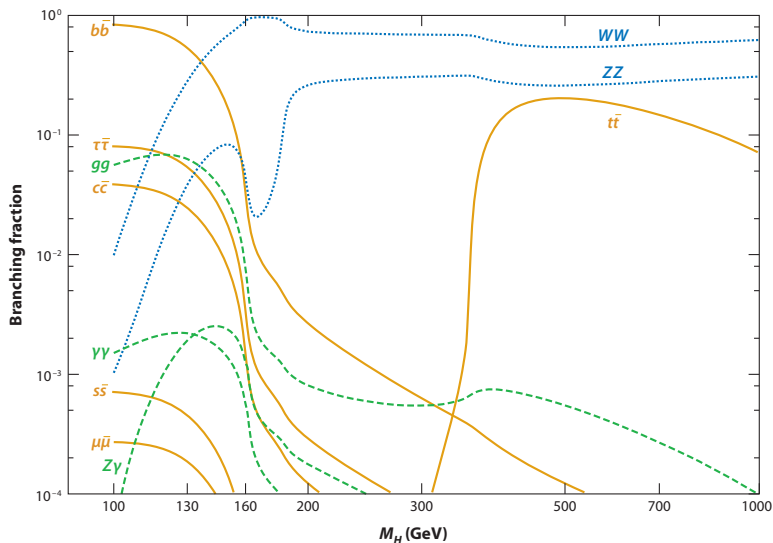
$$\Gamma(H \rightarrow Z^0Z^0) = \frac{G_F M_H^3}{64\pi\sqrt{2}} (1-x')^{1/2} (4-4x'+3x'^2) \quad x' \equiv 4M_Z^2/M_H^2$$

asymptotically  $\propto M_H^3$  and  $\frac{1}{2}M_H^3$ , respectively

$2x^2$  and  $2x'^2$  terms  $\Leftrightarrow$  decays into transverse gauge bosons

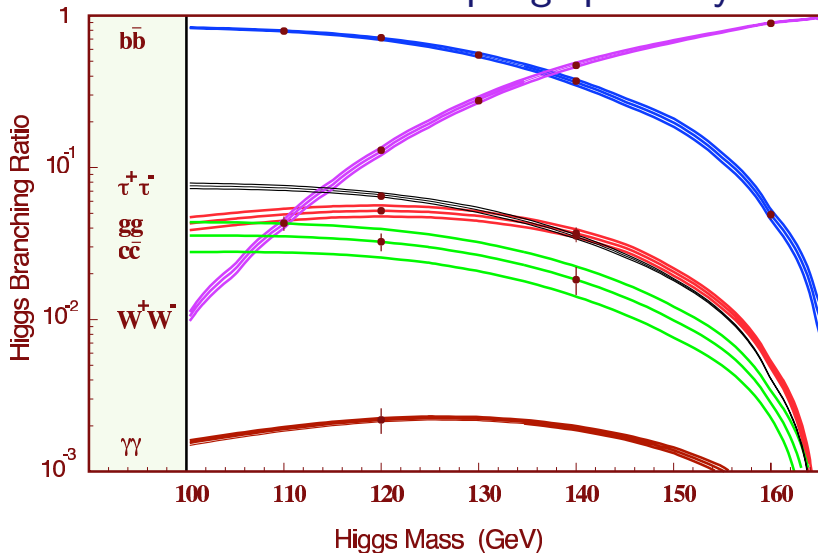
Dominant decays for large  $M_H$ : pairs of longitudinal weak bosons

# SM Higgs Boson Branching Fractions



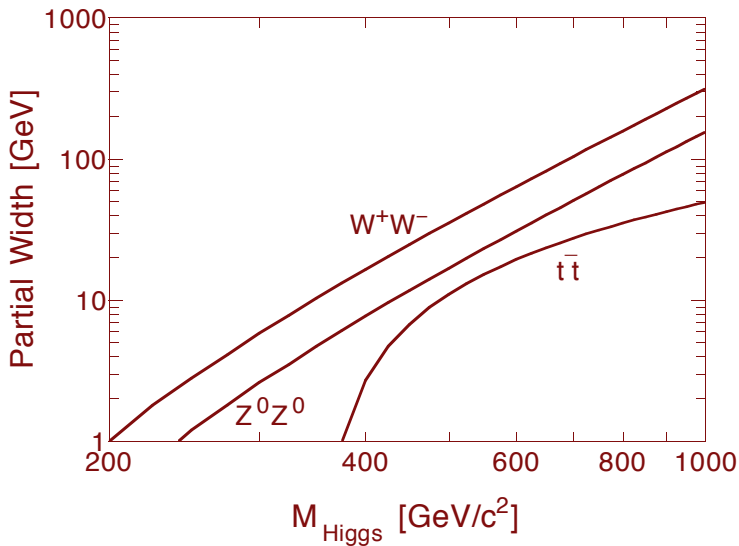
Djouadi, hep-ph/0503172

# ILC would measure $Hf\bar{f}$ couplings precisely



Points:  $500 \text{ fb}^{-1}$  @ 350 GeV    Bands: theory uncertainty ( $m_b$ )

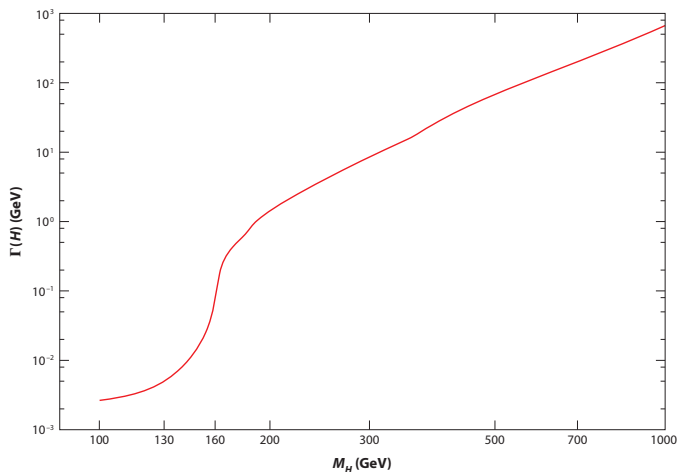
# Dominant decays at high mass



For  $M_H \rightarrow 1 \text{ TeV}$ , Higgs boson is *ephemeral*:  $\Gamma_H \rightarrow M_H$ .



# Total width of the standard-model Higgs boson



Below  $W^+W^-$  threshold,  $\Gamma_H \lesssim 1$  GeV

Far above  $W^+W^-$  threshold,  $\Gamma_H \propto M_H^3$

# A few words on Higgs production ...

$e^+e^- \rightarrow H$ : hopelessly small

$\mu^+\mu^- \rightarrow H$ : scaled by  $(m_\mu/m_e)^2 \approx 40\,000$

$e^+e^- \rightarrow HZ$ : prime channel

Hadron colliders:

$gg \rightarrow H \rightarrow b\bar{b}$ : background ?!

$gg \rightarrow H \rightarrow \tau\tau, \gamma\gamma$ : rate ?!

$gg \rightarrow H \rightarrow W^+W^-$ : best Tevatron sensitivity now

$\bar{p}p \rightarrow H(W, Z)$ : prime Tevatron channel for light Higgs

At the LHC:

Many channels accessible, search sensitive up to 1 TeV

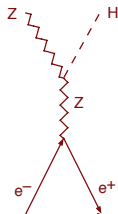
# Higgs search in $e^+e^-$ collisions

$\sigma(e^+e^- \rightarrow H \rightarrow \text{all})$  is *minute*,  $\propto m_e^2$

Even narrowness of low-mass  $H$  is not enough to make it visible ... Sets aside a traditional strength of  $e^+e^-$  machines—*pole physics*

Most promising:

associated production  $e^+e^- \rightarrow HZ$   
(has no small couplings)

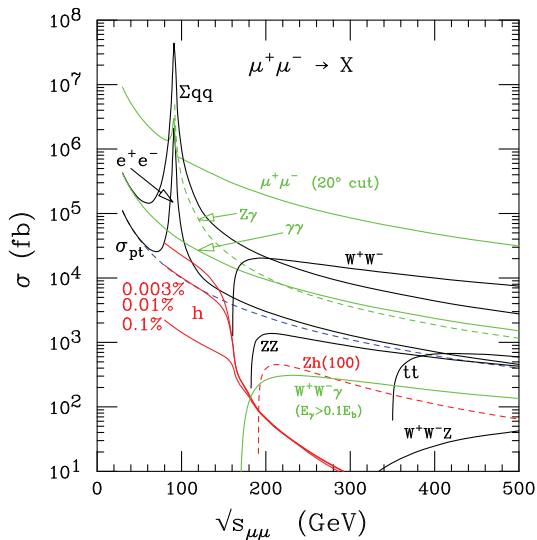


$$\sigma = \frac{\pi\alpha^2}{24\sqrt{s}} \frac{K(K^2 + 3M_Z^2)[1 + (1 - 4x_W)^2]}{(s - M_Z^2)^2 x_W^2(1 - x_W)^2}$$

$K$ : c.m. momentum of  $H$

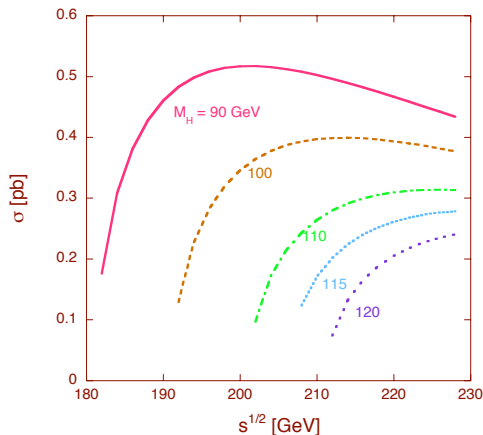
$x_W \equiv \sin^2 \theta_W$

$$l^+l^- \rightarrow X \dots$$



$$\sigma(e^+e^- \rightarrow H) = (m_e/m_\mu)^2 \sigma(\mu^+\mu^- \rightarrow H) \approx \sigma(\mu^+\mu^- \rightarrow H)/40\,000$$

# LEP Range Excitation Curves (ISR not included)

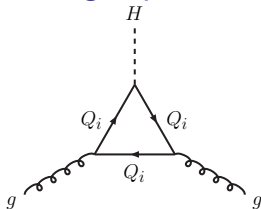


LEP 2: sensitive nearly to kinematical limit

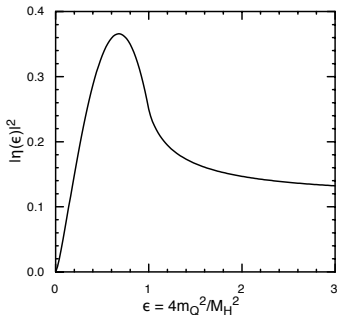
$$M_H^{\max} = \sqrt{s} - M_Z$$

LC: sensitive for  $M_H \lesssim 0.7\sqrt{s}$

# $H$ couples to gluons through quark loops

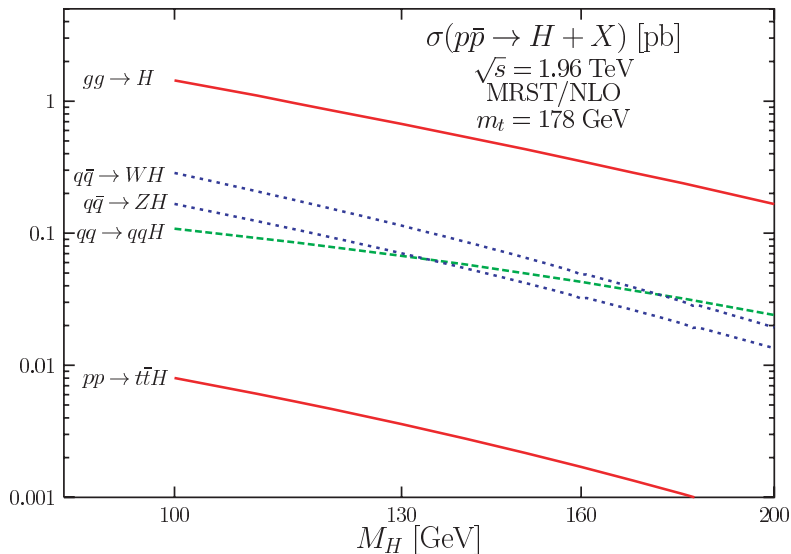


Only heavy quarks matter:



heavy 4th generation would have a big effect

# Higgs-boson production at the Tevatron



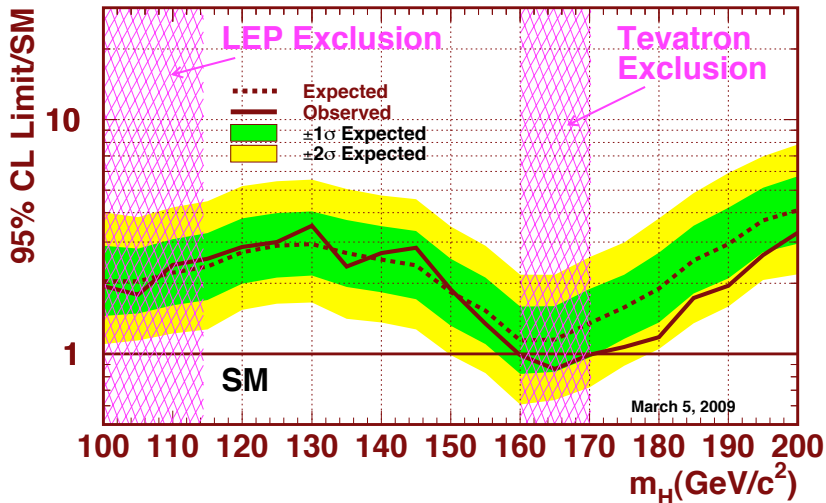
Djouadi

Update 1

Update 2

# Current Tevatron Sensitivity

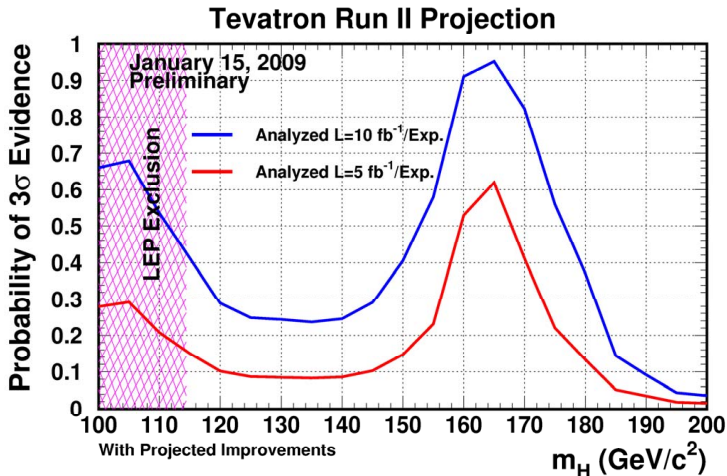
Tevatron Run II Preliminary,  $L=0.9-4.2 \text{ fb}^{-1}$



combining experiments, channels: Winter 2009

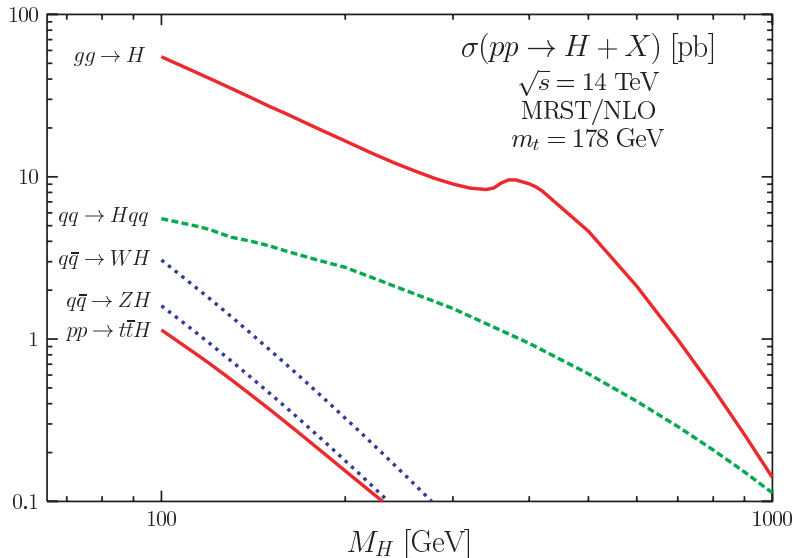


# Tevatron prospects ...



**> 40% probability to have a 3 sigma evidence for  $m_H = 115 \text{ GeV}$  Higgs with  $L=10 \text{ fb}^{-1}$**

# LHC cross sections ...



Djouadi

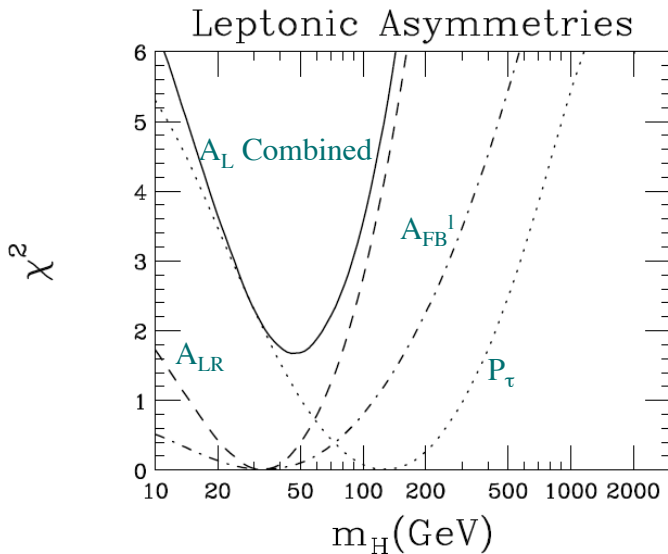
## A Cautionary Note

- $A_{FB}^b$ , which exerts the greatest “pull” on the global fit, is most responsible for raising  $M_H$  above the range excluded by direct searches.
- Leptonic and hadronic observables point to different best-fit values of  $M_H$
- Many subtleties in experimental and theoretical analyses

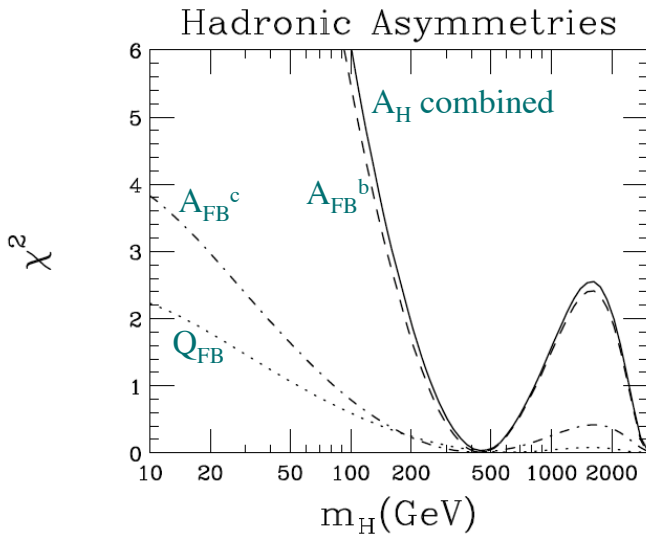
M. Chanowitz, [arXiv:0806.890](https://arxiv.org/abs/0806.890)

Introduction to global analyses: J. L. Rosner, [hep-ph/0108195](https://arxiv.org/abs/hep-ph/0108195);  
[hep-ph/0206176](https://arxiv.org/abs/hep-ph/0206176)

# $\chi^2$ Distributions: Leptonic Asymmetries



# $\chi^2$ Distributions: Hadronic Asymmetries



# The EW scale and beyond

EWSB scale,  $v = (G_F \sqrt{2})^{-\frac{1}{2}} \approx 246$  GeV, sets

$$M_W^2 = g^2 v^2 / 2 \quad M_Z^2 = M_W^2 / \cos^2 \theta_W$$

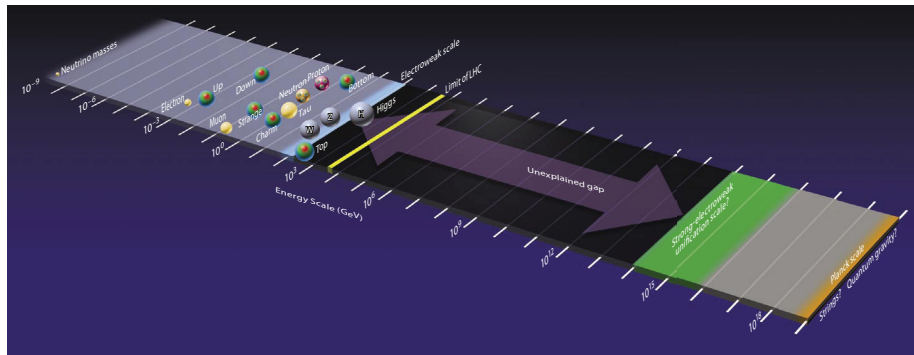
But it is not the only scale of physical interest

quasi-certain:  $M_{\text{Planck}} = 1.22 \times 10^{19}$  GeV

probable:  $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  unification scale  
 $\sim 10^{15-16}$  GeV

somewhere: flavor scale?

# The Hierarchy Problem



How to keep the distant scales from mixing in the face of quantum corrections? *OR*

How to stabilize the mass of the Higgs boson on the electroweak scale? *OR*

Why is the electroweak scale small?