



New Physics from Oscillations at the DUNE Near Detector

In collaboration with P. Coloma, J. López-Pavón & S. Urrea, based on JHEP 08 (2021) 065

Salvador Rosauro-Alcaraz 02/06/2023



Introduction

Standard neutrino oscillations

What we know (at 1σ)

I. Esteban et al. 2007.14792 www.nu-fit.org

Solar sector $\begin{cases} \sin^2 \theta_{12} = 0.304^{+0.012}_{-0.012} \\ \Delta m_{21}^2 = 7.42^{+0.21}_{-0.20} \cdot 10^{-5} eV^2 \end{cases}$

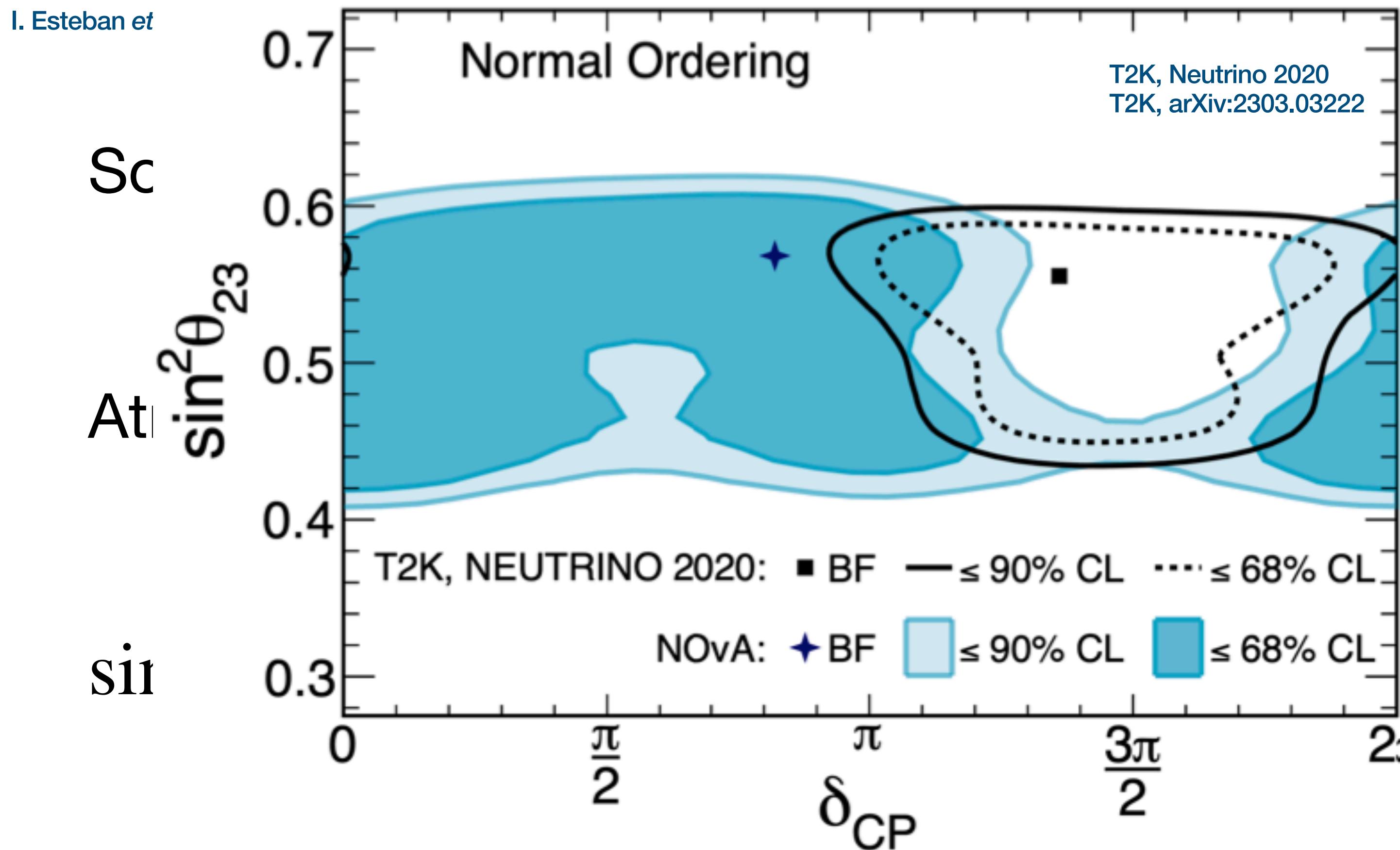
Atm. sector $\begin{cases} \sin^2 \theta_{23} = 0.573^{+0.016}_{-0.020} \\ |\Delta m_{31}^2| = 2.517^{+0.026}_{-0.028} \cdot 10^{-3} eV^2 \end{cases}$

$$\sin^2 \theta_{13} = 0.02219^{+0.00062}_{-0.00063}$$

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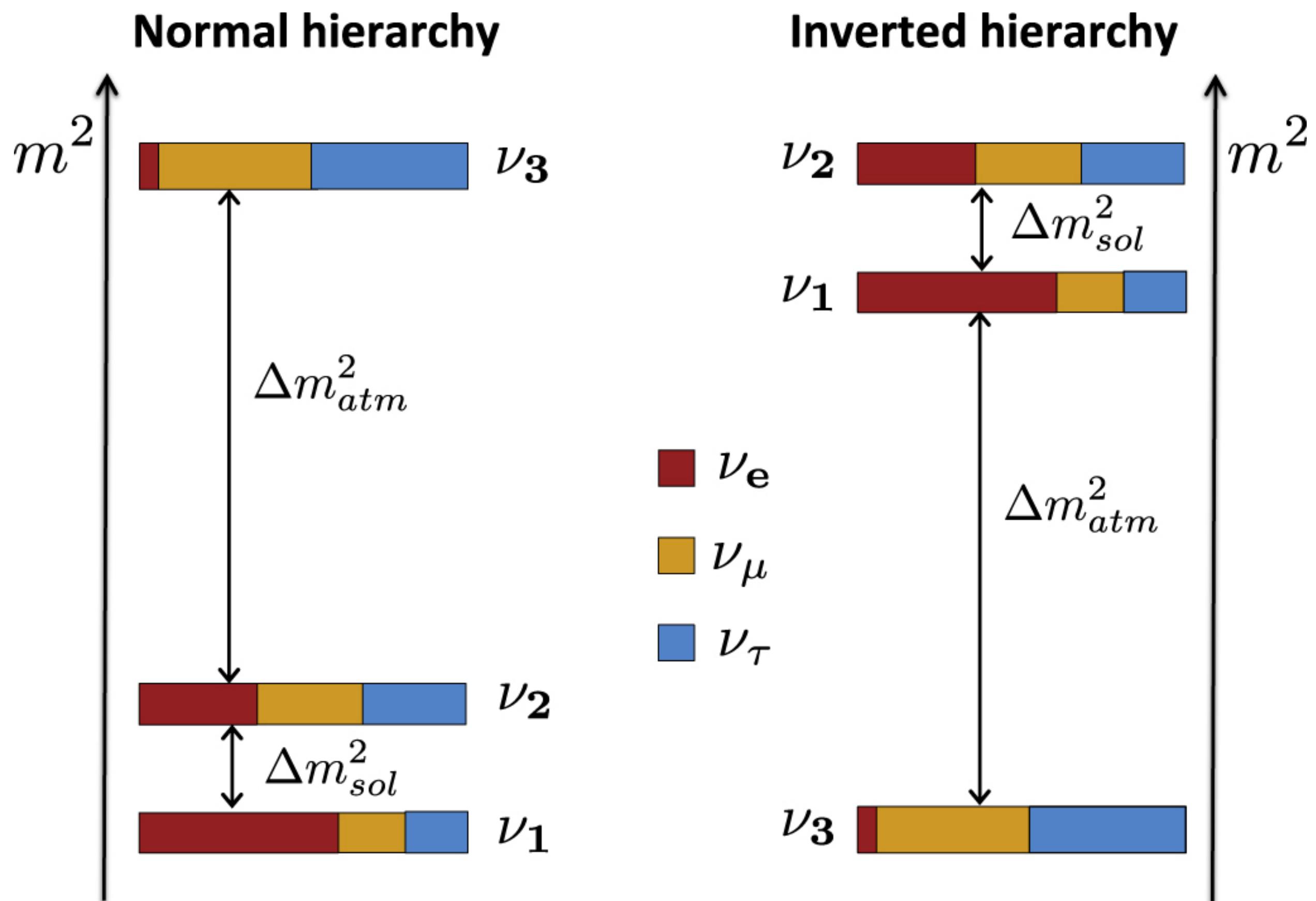


What we do not know (yet)

Is there leptonic
CP violation, i.e., $\delta \neq 0, \pi$?

Introduction

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Mass ordering: $sign(\Delta m_{31}^2)$

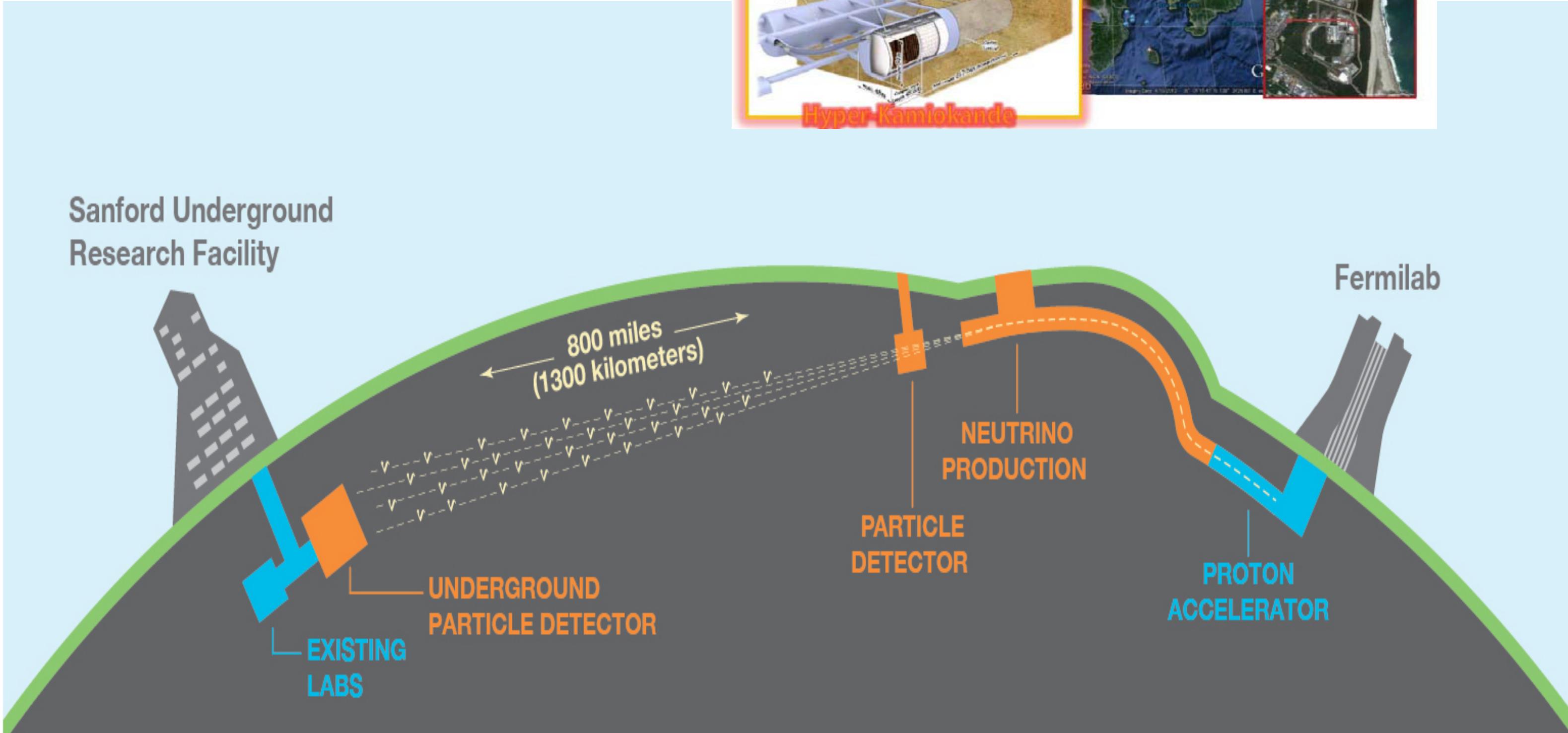
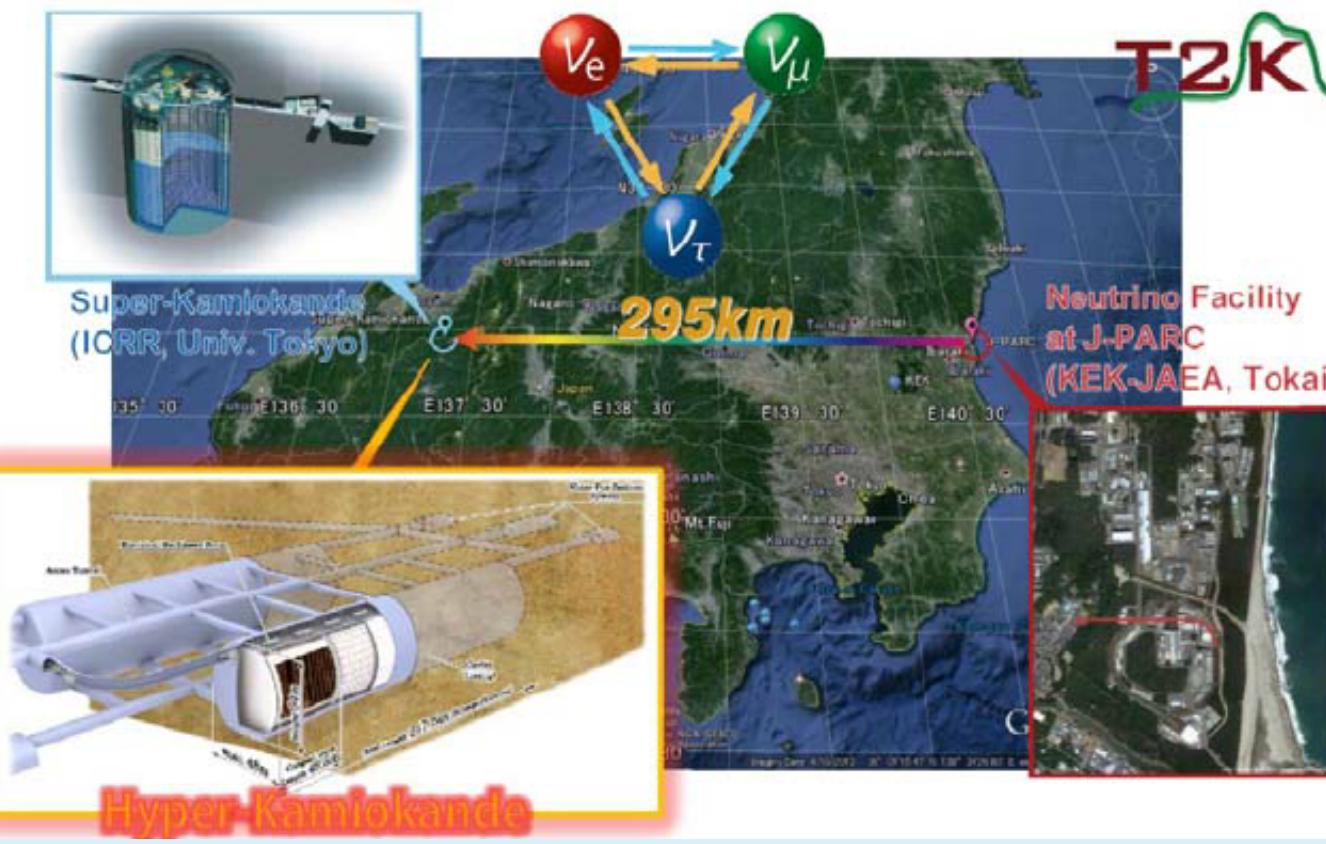
Octant of θ_{23}

Introduction

Future facilities

DUNE & T2HK

DUNE Collaboration, arXiv:2006.16043
T2HK Collaboration, arXiv:1412.4673



What we **do not** know (yet)

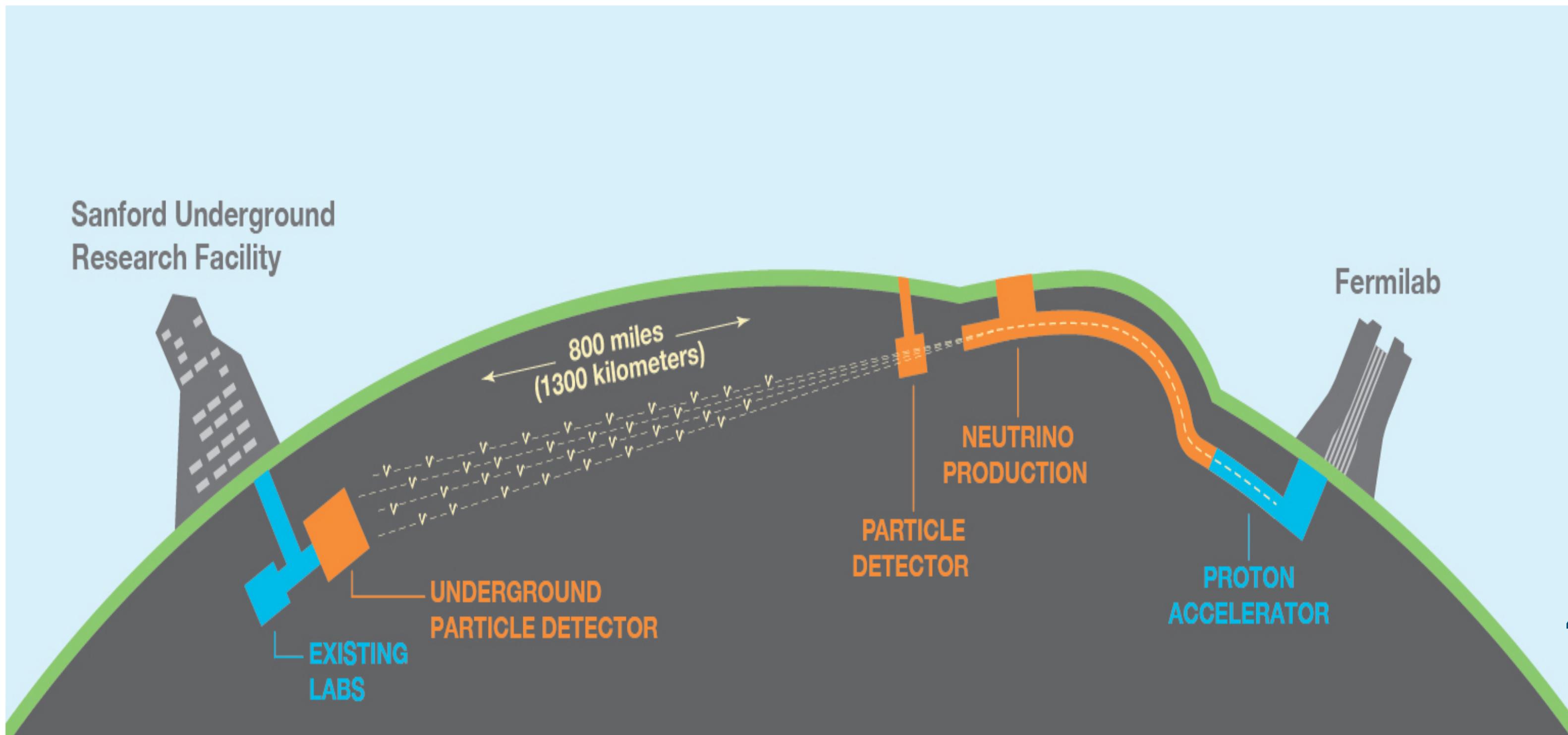
Is there leptonic CP violation, i.e., $\delta \neq 0, \pi$?

Mass ordering: $sign(\Delta m_{31}^2)$

Octant of θ_{23}

Recipe for good measurements

$\vec{\sigma}_{FD}$ correlated with $\vec{\sigma}_{ND}$



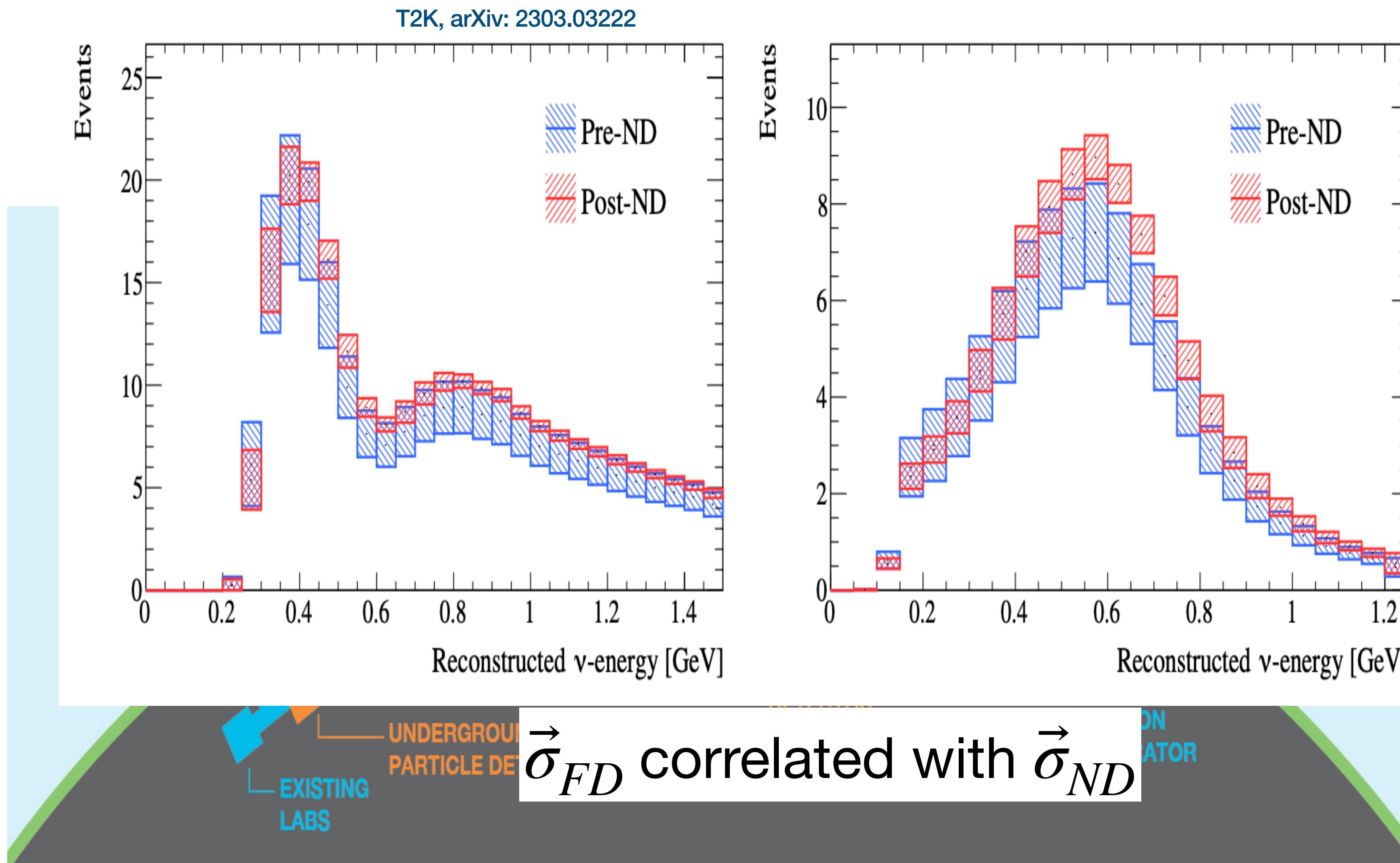
1 Measure events in ND,
 $N_\mu^{ND}(\vec{\sigma}_{ND})$

2 Measure events in FD,
 $N_e^{FD}(\vec{\theta}, \vec{\sigma}_{FD})$

3 Find ratio $R_{e\mu} = \frac{N_e^{FD}}{N_\mu^{ND}}$

4 Extract $\vec{\theta}$ from comparing it
with the extrapolation of N_μ^{ND}

Recipe for good measurements



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Physics at a ND

Beyond 3ν oscillations

Need right-handed neutrinos to explain light neutrino masses

$$\nu_\alpha = \sum_i^{3+N_s} \mathcal{U}_{\alpha i} n_i$$

$$\mathcal{U} = \begin{pmatrix} N & \Theta \\ R & S \end{pmatrix}$$

PMNS mixing matrix



Physics at a ND

Beyond 3ν oscillations

Need right-handed neutrinos to explain light neutrino masses

$$\nu_\alpha = \sum_{i=1}^3 N_{\alpha i} n_i + \sum_{i=4}^{N_s} \Theta_{\alpha i} n_i$$

$$U = \begin{pmatrix} N \\ R \\ S \end{pmatrix} \quad \Theta = \begin{pmatrix} \Theta \\ S \end{pmatrix}$$

$$N = (I - T)U$$

$$T = \begin{pmatrix} \alpha_{ee} & 0 & 0 \\ \alpha_{\mu e} & \alpha_{\mu\mu} & 0 \\ \alpha_{\tau e} & \alpha_{\tau\mu} & \alpha_{\tau\tau} \end{pmatrix}$$

Physics at a ND

Beyond 3ν oscillations

Need right-handed neutrinos to explain light neutrino masses

$$\nu_\alpha = \sum_{i=1}^3 N_{\alpha i} n_i + \sum_{i=4}^{N_s} \Theta_{\alpha i} n_i$$

Two regimes for non-unitarity:

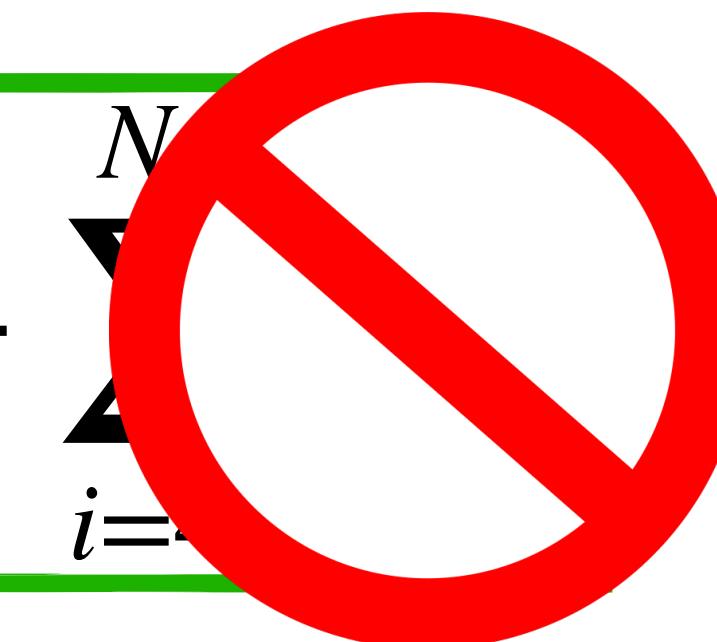
- n_i not produced in neutrino oscillation experiments
- n_i participating in neutrino oscillations

Physics at a ND

Beyond 3ν oscillations

Need right-handed neutrinos to explain light neutrino masses

$$\nu_\alpha = \sum_{i=1}^3 N_{\alpha i} n_i + \sum_{i=4}^N$$



Two regimes for non-unitarity:

- n_i not produced in neutrino oscillation experiments
- n_i participating in neutrino oscillations

Stronger bounds
elsewhere

Physics at a ND

Beyond 3ν oscillations

Need right-handed neutrinos to explain light neutrino masses

$$\nu_\alpha = \sum_{i=1}^3 N_{\alpha i} n_i + \sum_{i=4}^{N_s} \Theta_{\alpha i} n_i$$

Two regimes for non-unitarity:

- n_i not produced in neutrino oscillation experiments
- n_i participating in neutrino oscillations

$$\Delta m_{41}^2 \gg |\Delta m_{31}^2|$$

Study effects at ND

Physics at a ND

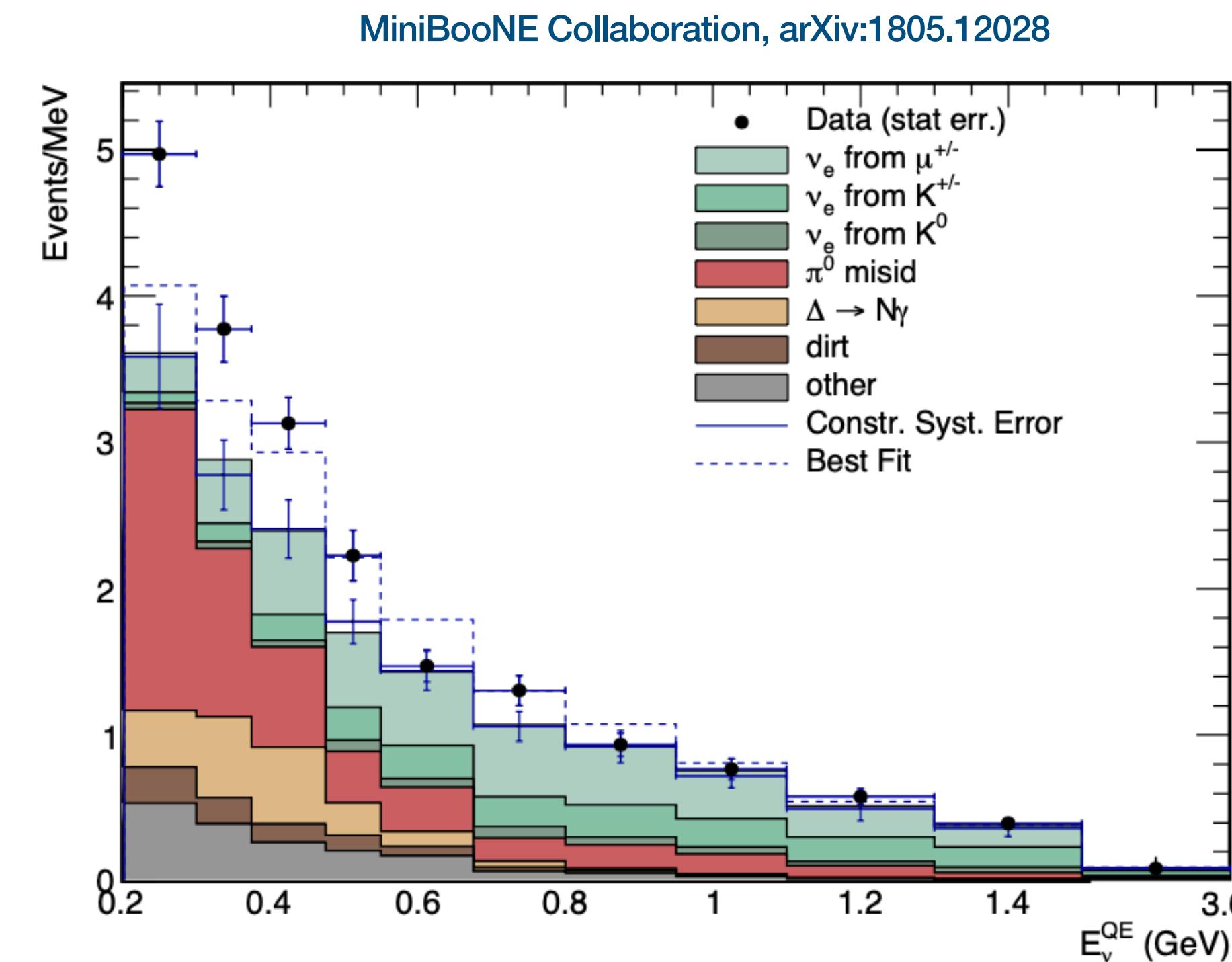
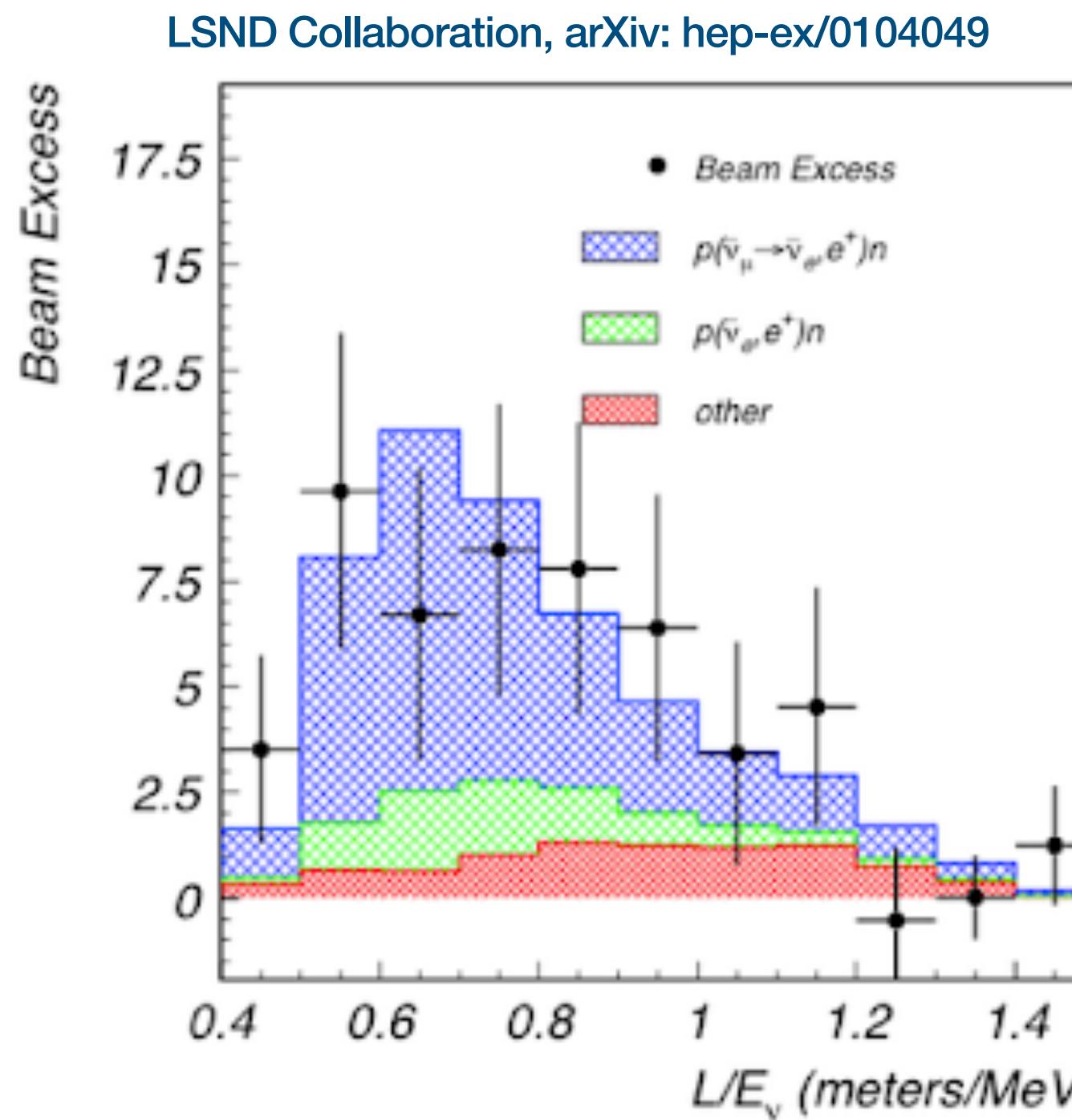
Sterile neutrino oscillations

$$P_{\gamma\beta} = \sin^2 2\theta_{\gamma\beta} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right), \quad \sin^2 \theta_{\gamma\beta} \equiv 4 \left| \mathcal{U}_{\gamma 4} \right|^2 \left| \mathcal{U}_{\beta 4} \right|^2$$

Physics at a ND

Sterile neutrino oscillations

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Anomalous $\nu_\mu \rightarrow \nu_e$

Physics at a ND

Non-unitarity from low-scale physics

$$P_{\gamma\beta} = \sin^2 2\theta_{\gamma\beta} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right), \quad \sin^2 \theta_{\gamma\beta} \equiv 4 \left| \mathcal{U}_{\gamma 4} \right|^2 \left| \mathcal{U}_{\beta 4} \right|^2$$

Averaged-out regime $\Delta m_{41}^2 L/E \gg 1$

$$P_{\gamma\beta} = 2 \left| \alpha_{\gamma\beta} \right|^2 = 2 \left| \mathcal{U}_{\gamma 4} \right|^2 \left| \mathcal{U}_{\beta 4} \right|^2$$

Similar to NU at high scales

Physics at a ND

Non-standard neutrino interactions

General 4-fermion effective operator $\mathcal{L}_{NSI} = -2\sqrt{2}G_F\epsilon_{\alpha\beta}\left(\bar{\nu}_\alpha\gamma_\mu P_L\nu_\beta\right)\left(\bar{f}\gamma^\mu P_{L(R)}f\right)$

Used to “reconcile” NOvA and T2K results

P. Denton, J. Gherlein & R. Pestes, arXiv: 2008.01110
S. Chatterjee & A. Palazzo, arXiv: 2008.04161

Physics at a ND

Non-standard neutrino interactions

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$$P_{\gamma\beta} = \left| \epsilon_{\gamma\beta}^d \right|^2 + \left| \epsilon_{\gamma\beta}^s \right|^2 + 2 \left| \epsilon_{\gamma\beta}^d \right| \left| \epsilon_{\gamma\beta}^s \right| \cos(\Delta\Phi_{\gamma\beta})$$

Physics at a ND

Non-standard neutrino interactions

General 4-fermion effective operator $\mathcal{L}_{NSI} = -2\sqrt{2}G_F \epsilon_{\alpha\beta} (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P_{L(R)} f)$

$$P_{\gamma\beta} = |\epsilon_{\gamma\beta}^d|^2 + |\epsilon_{\gamma\beta}^s|^2 + 2 |\epsilon_{\gamma\beta}^d| |\epsilon_{\gamma\beta}^s| \cos(\Delta\Phi_{\gamma\beta})$$

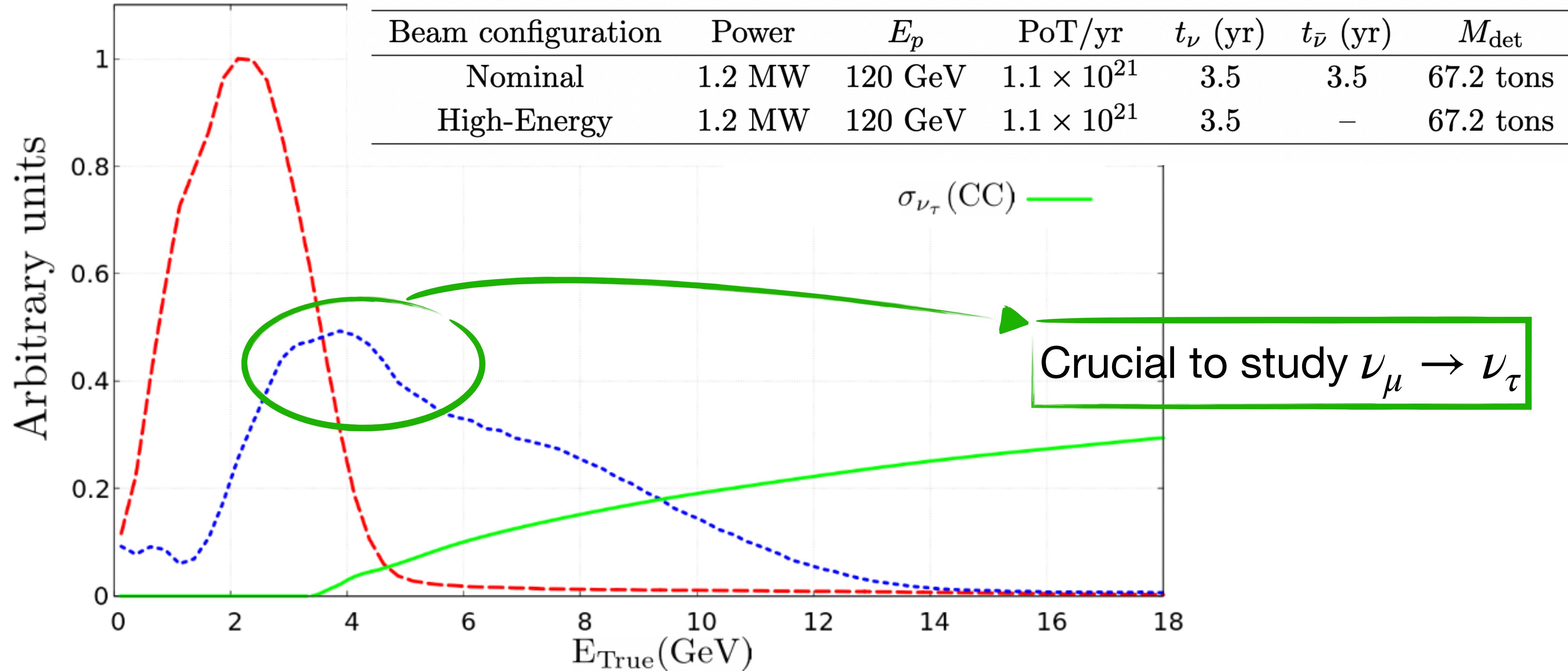
Averaged-out regime $\Delta m_{41}^2 L/E \gg 1$

$$P_{\gamma\beta} = 2 |\alpha_{\gamma\beta}|^2$$

Translate bounds from $\alpha_{\gamma\beta}$ to $\epsilon_{\gamma\beta}$

Simulation details

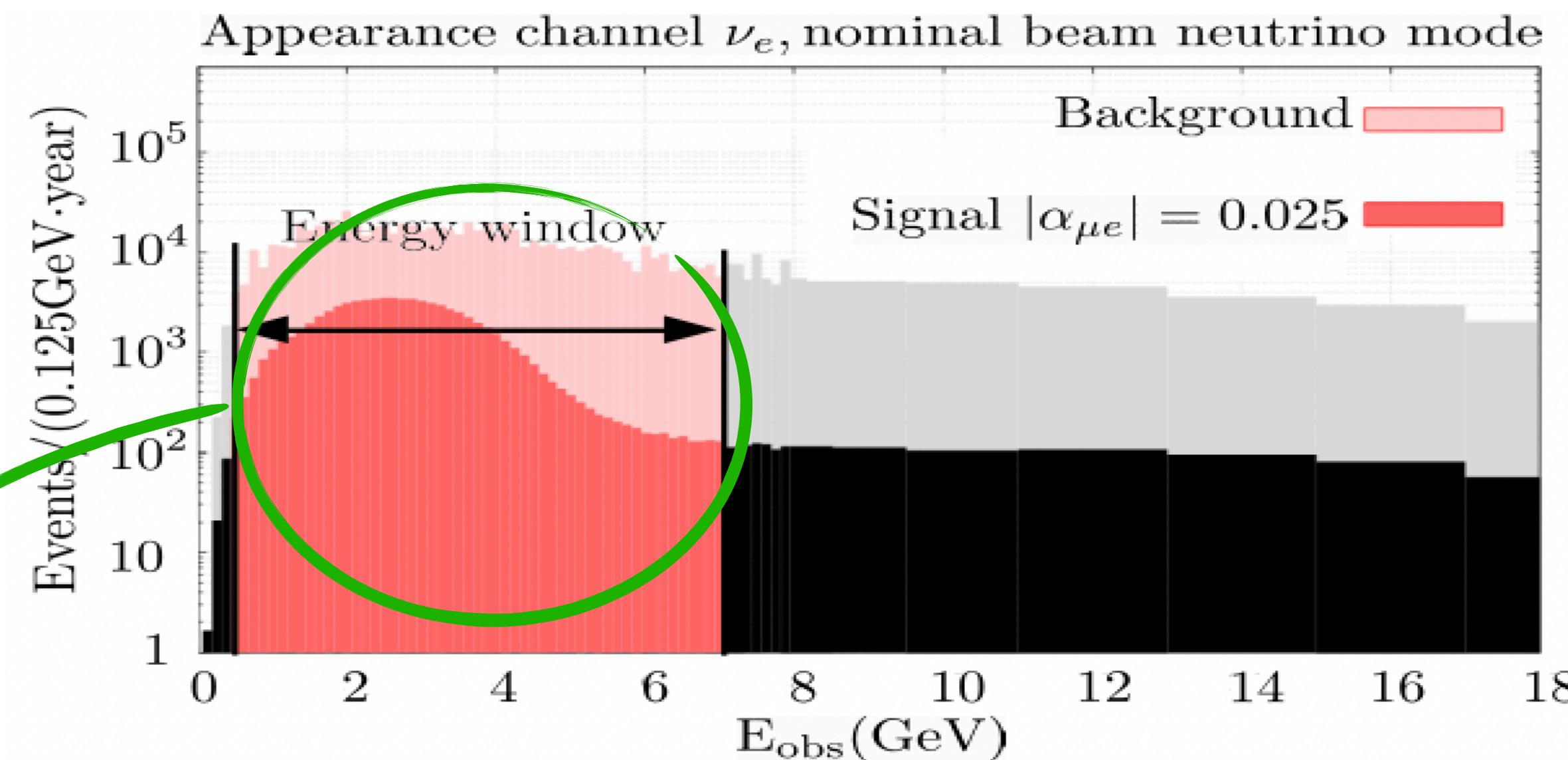
DUNE flux & detector simulation



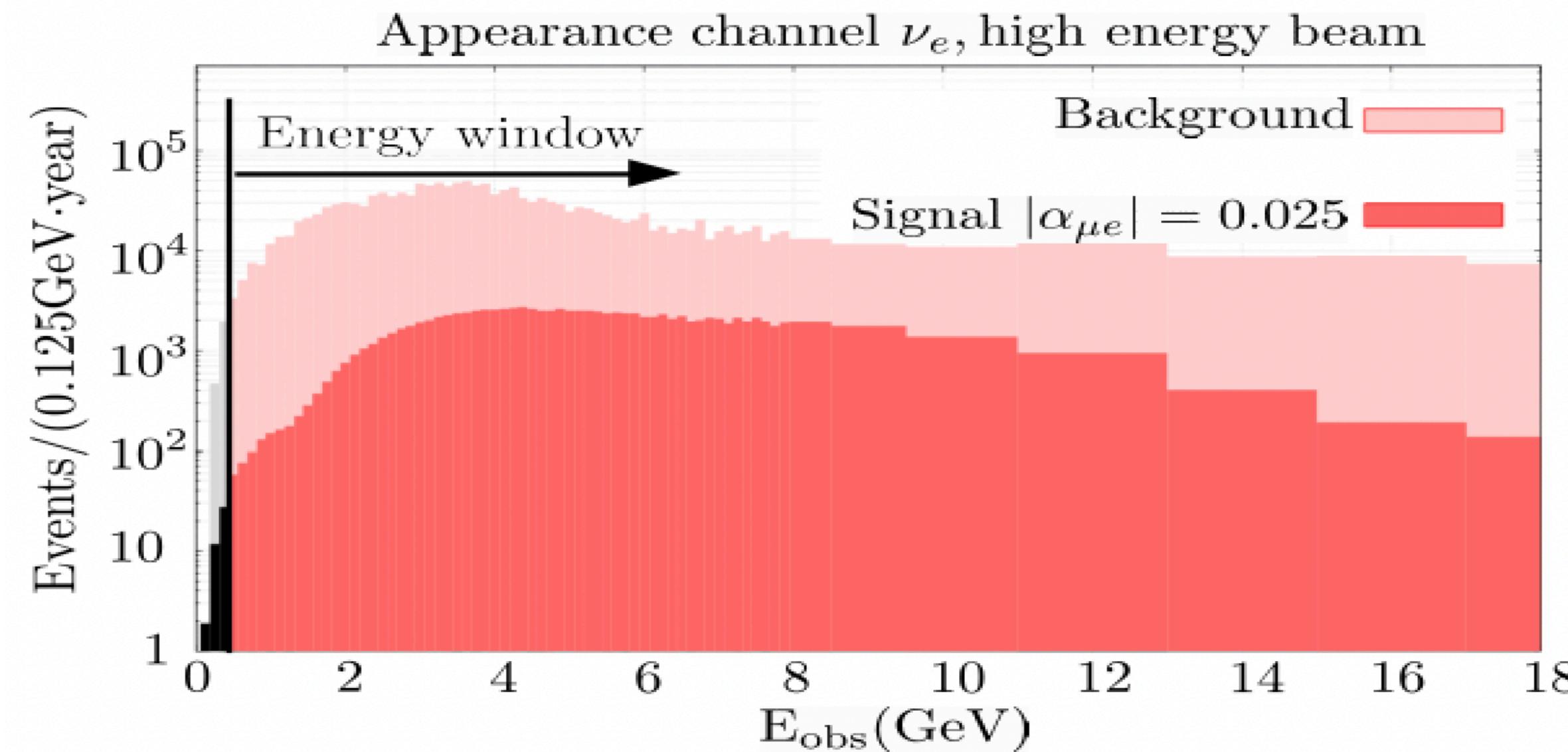
Simulation details

Event rates

Events included
in the analysis



Nominal beam

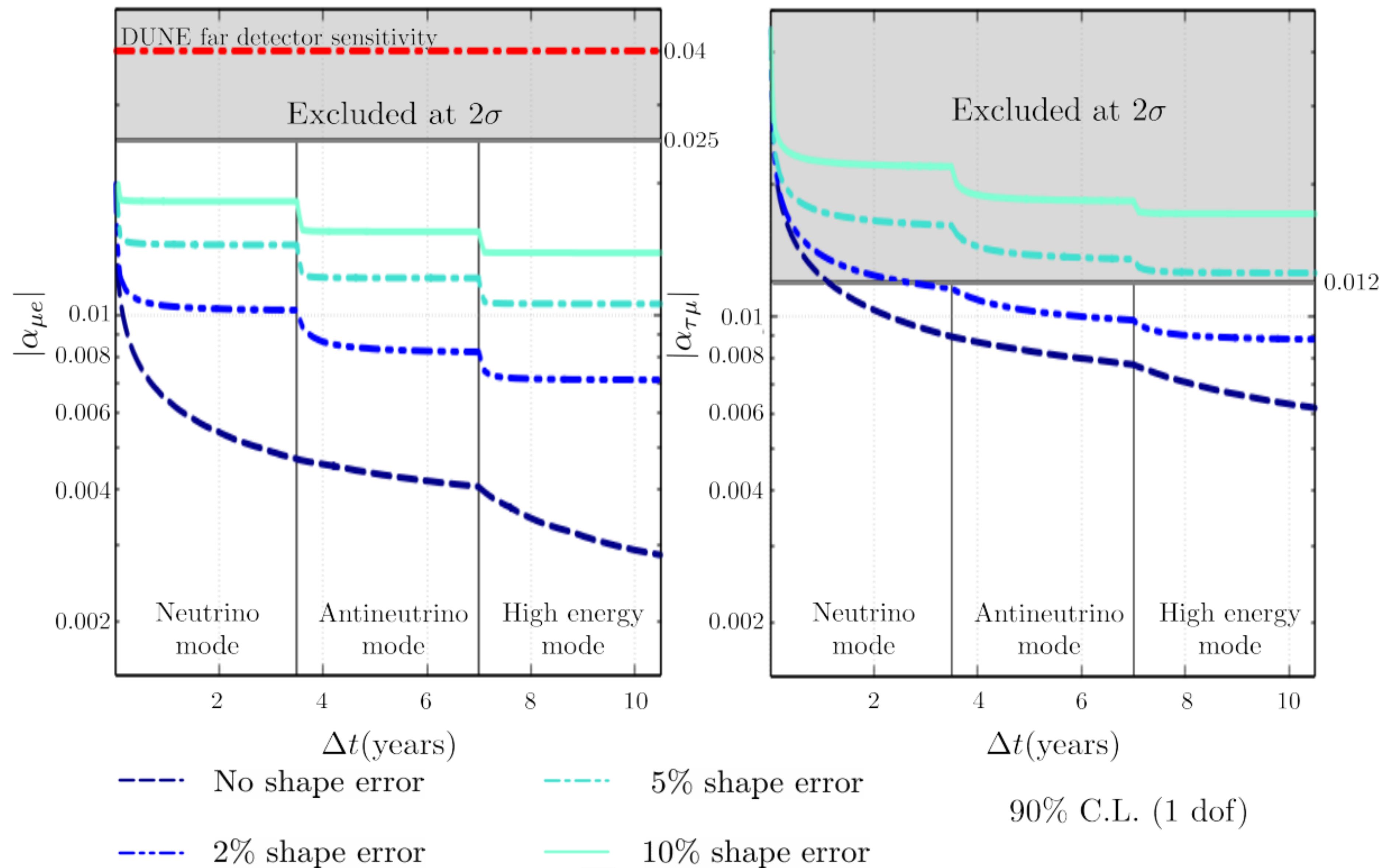


High energy beam

Results

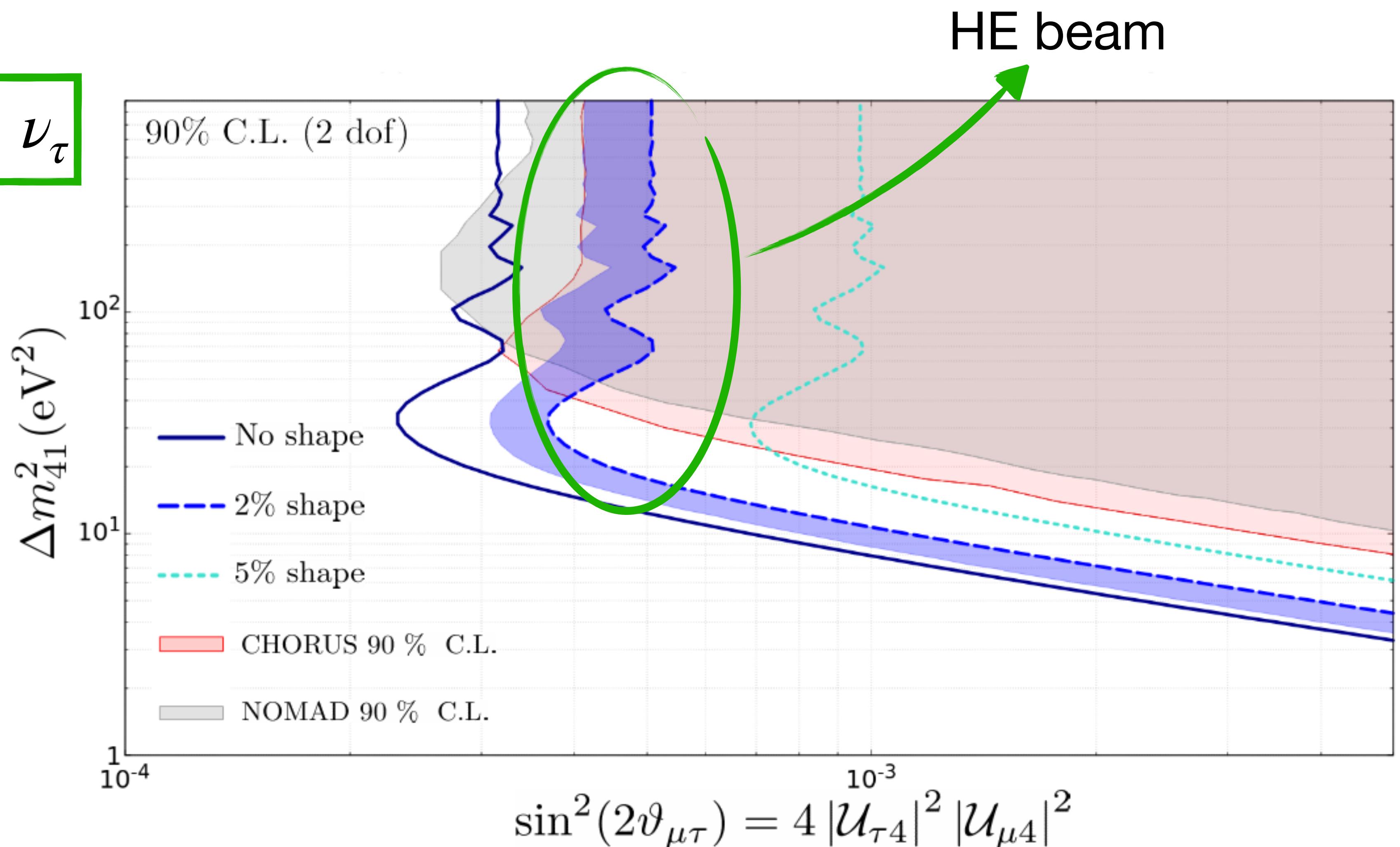
Non-unitarity

$$\Delta m_{41}^2 L/E \gg 1$$



Results

Sterile oscillations



Results

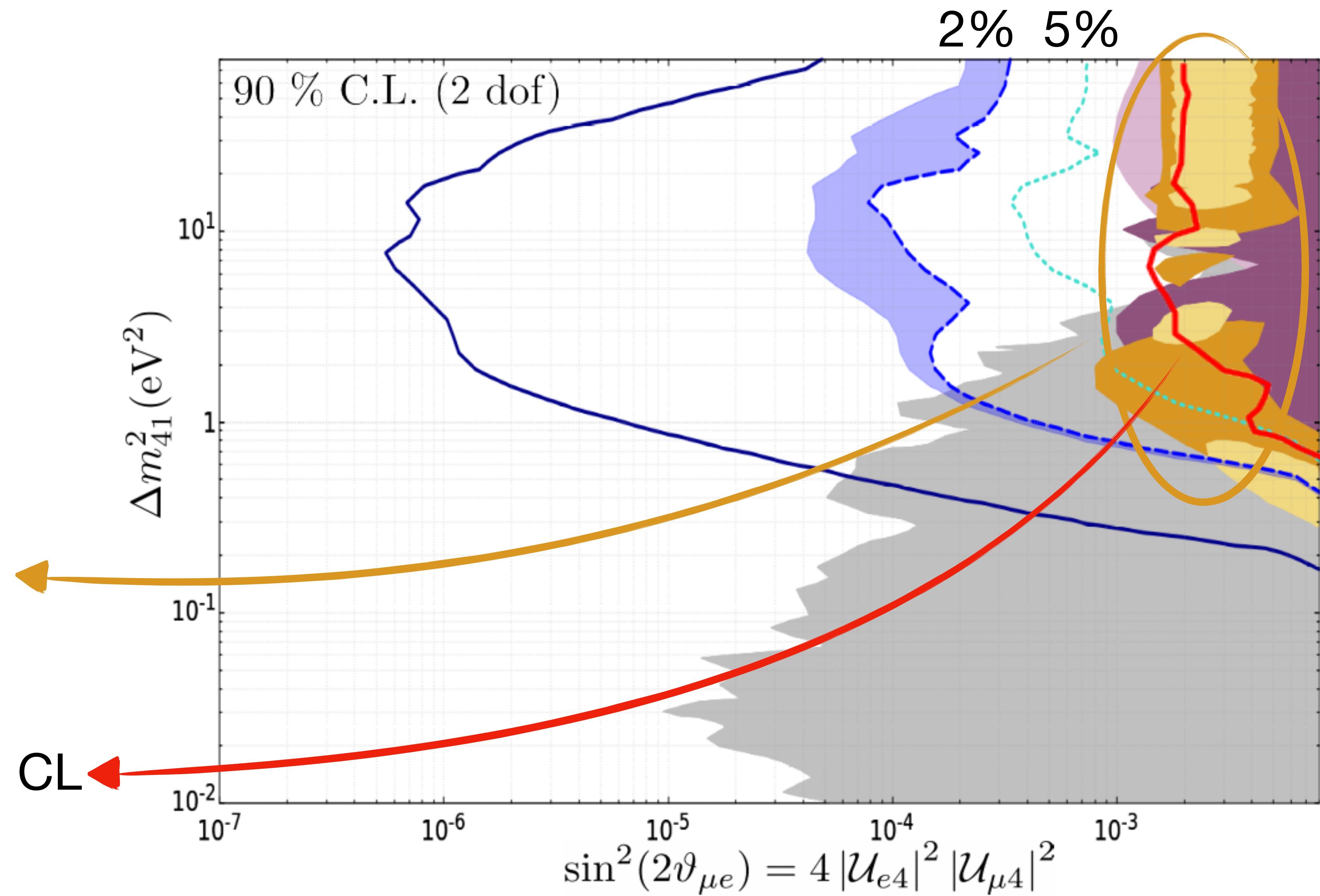
Sterile oscillations

LSND & MiniBooNE
preferred region @ 99% CL

MicroBooNE @ 95% CL

MicroBooNE, arXiv: 2210.10216

$$\nu_\mu \rightarrow \nu_e$$



Conclusions

Near detectors can be useful to study physics beyond 3ν oscillations

Systematic uncertainties play a crucial role

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Systematic uncertainties play a crucial role

Study ν_τ appearance in DUNE

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Near detectors can be useful to study physics beyond 3ν oscillations

Systematic uncertainties play a crucial role

Study ν_τ appearance in DUNE

2% shape uncertainties:

- Reduce considerably the sensitivity to $\nu_\mu \rightarrow \nu_e$ around 2 orders of magnitude
- Reduce bound on NU parameters about a factor 2

Thank you!

Back-up slides

Simulation details

ν_τ detection

PDG, P. A. Zayla *et al.*, PTEP 2020
GENIE Collaboration, J. Tena-Vidal *et al.*, arXiv:2104.09179
A. de Gouvêa *et al.*, arXiv:1904.07265

Produce τ through CC interactions

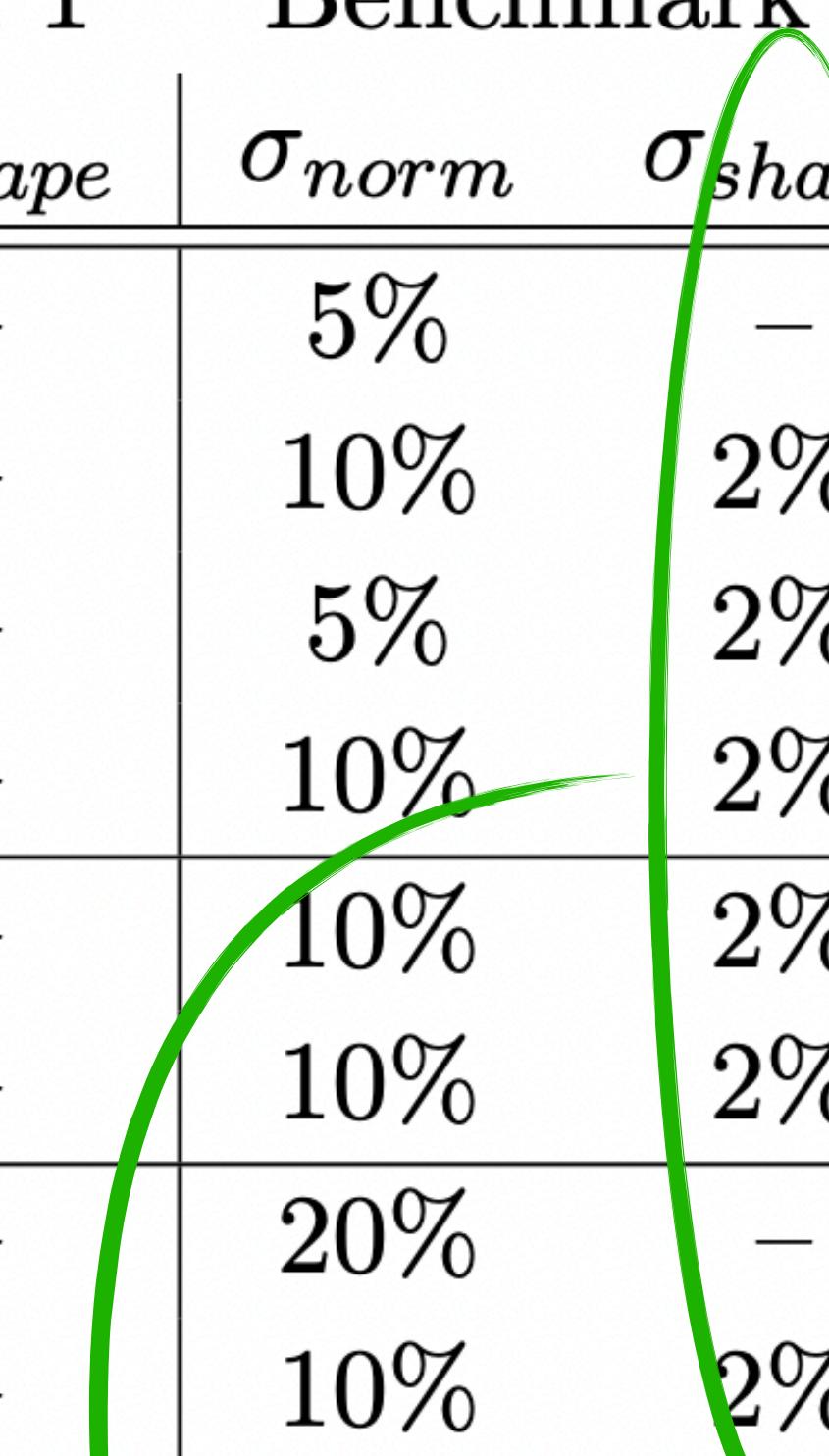
Consider only τ hadronic decay ($BR(\tau \rightarrow had) \sim 65\%$)

Main background due to NC

Simulation details

Systematic uncertainties

Event sample	Contribution	Benchmark 1		Benchmark 2		Benchmark 3	
		σ_{norm}	σ_{shape}	σ_{norm}	σ_{shape}	σ_{norm}	σ_{shape}
ν_e -like	Signal	5%	–	5%	–	5%	–
	Intrinsic cont.	10%	–	10%	2%	10%	5%
	Flavor mis-ID	5%	–	5%	2%	5%	5%
	NC	10%	–	10%	2%	10%	5%
ν_μ -like	$\nu_\mu, \bar{\nu}_\mu$ CC (signal)	10%	–	10%	2%	10%	5%
	NC	10%	–	10%	2%	10%	5%
ν_τ -like	Signal	20%	–	20%	–	20%	–
	NC	10%	–	10%	2%	10%	5%



Allows every bin to vary independently

Results

Non-standard neutrino interactions

$$P_{\gamma\beta} = \left| \epsilon_{\gamma\beta}^d \right|^2 + \left| \epsilon_{\gamma\beta}^s \right|^2 + 2 \left| \epsilon_{\gamma\beta}^d \right| \left| \epsilon_{\gamma\beta}^s \right| \cos(\Delta\Phi_{\gamma\beta})$$

Results

Non-standard neutrino interactions

Appearance channel ν_τ 90% C.L.

