



# Freeze-in baryogenesis and early matter domination

Based on [arXiv:2111.05740](https://arxiv.org/abs/2111.05740), [arXiv:2204.13554](https://arxiv.org/abs/2204.13554), [arXiv:2304.07345](https://arxiv.org/abs/2304.07345)  
In collaboration with I. Dalianis, D. Karamitros, P. Papachristou, V. Spanos

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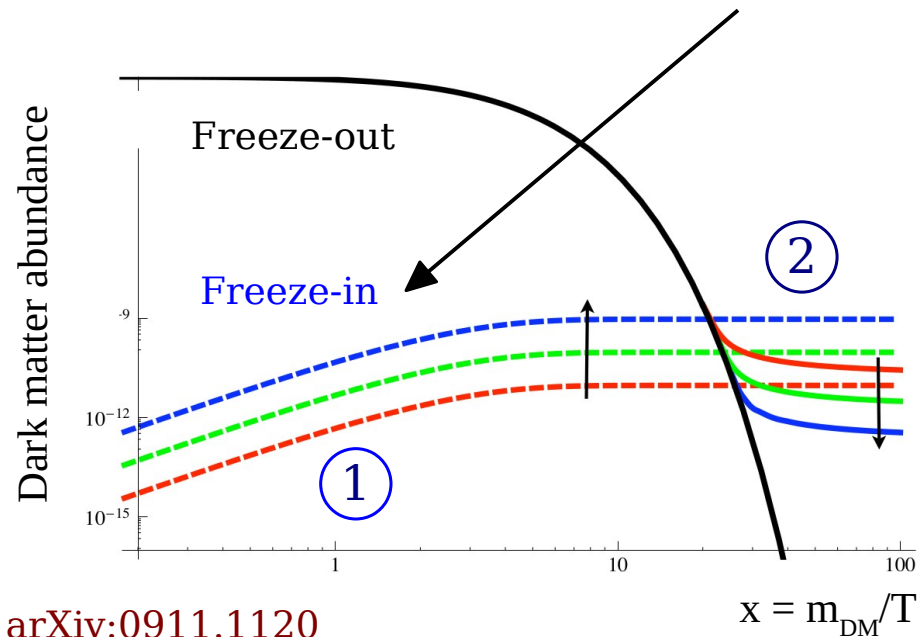
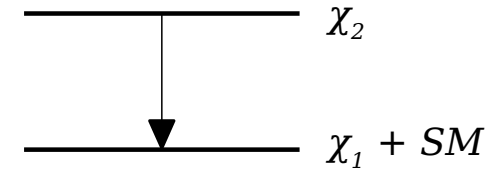
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# FIMPy baryogenesis: general idea

arXiv:2004.00636, arXiv:2201.11502,  
arXiv:2111.05740, arXiv:2204.13554

Freeze-in involves *very* weakly (“feebly”) interacting particles (FIMPs) that don’t reach thermal equilibrium with the SM thermal bath in the early Universe.

• Such particles can be produced *e.g.* from the decay of some heavier state or from annihilations of bath particles.

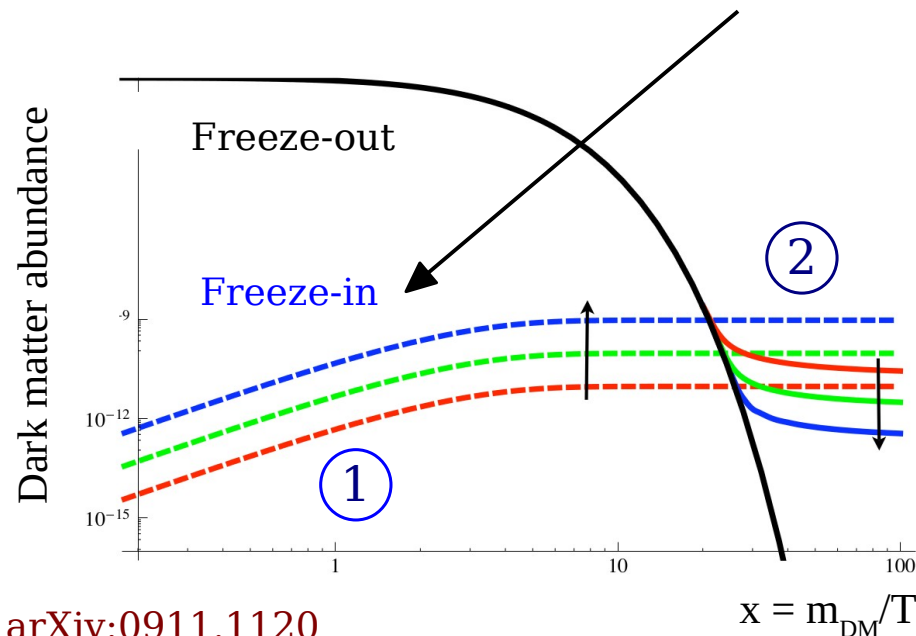
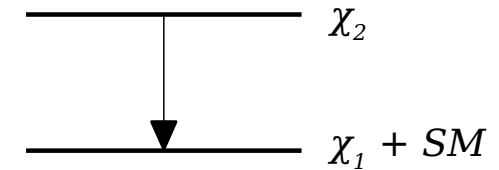


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• Such particles can be produced *e.g.* from the decay of some heavier state or from annihilations of bath particles.



• The decays and/or annihilations that are responsible for dark matter production can also violate both the baryon number  $B$  (or  $L$ ) and  $C/CP$ .

• But, by construction, in freeze-in these processes are also out-of-equilibrium.

→ All three Sakharov conditions can be satisfied.

NB: in arXiv:2004.00636 and arXiv:2201.11502  
 $CP$  violation is rather due to DM oscillations

# A concrete realization: toy model

arXiv:2204.13554

Consider an extension of the SM by a complex scalar field and two vector-like fermions, described by the Lagrangian :

$$\mathcal{L}_{\text{int}} = \frac{\lambda_1}{2\Lambda} (\bar{e}F_1) \varphi^* \varphi^* + \frac{\lambda_2}{2\Lambda} (\bar{e}F_2) \varphi^* \varphi^* + \frac{\kappa}{\Lambda^2} (\bar{e}F_1) (\bar{F}_2 e) + \text{h.c.}$$

- The bulk of dark matter production takes place at the highest considered temperature (in this framework the “reheating temperature”  $T_{\text{RH}}$ ).
- Scattering processes dominate both DM production and baryogenesis.

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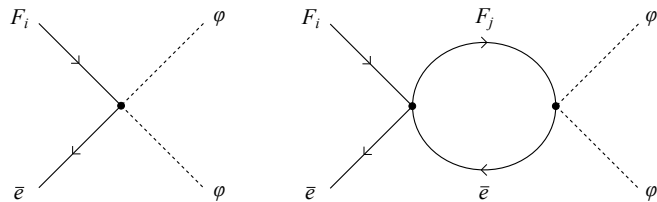
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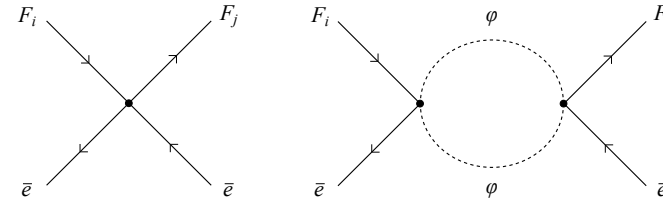
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**NB:** Qualitatively similar (but quantitatively different) results are obtained if we, instead, assume a fermion DM candidate and operators of higher dimension.



Dark Matter



Baryogenesis

$$\Omega_{\text{DM}} h^2 \approx \frac{2.4 \times 10^{23}}{g_{*s} \sqrt{g_{*p}}} \frac{A 4^n n! (n+1)!}{2n-1} \frac{T_{\text{RH}}^{2n-1} m_{\text{DM}}}{\Lambda^{2n}}$$

$$\begin{aligned} -sHT \frac{dY_{\Delta F_i}}{dT} &= -[F_i \bar{e} \leftrightarrow \varphi \varphi] - [F_i \bar{e} \leftrightarrow F_j \bar{e}] + (-1)^i [F_i \bar{F}_j \leftrightarrow e \bar{e}] \\ -sHT \frac{dY_{B-L_{\text{SM}}}}{dT} &= -\sum_i [F_i \bar{e} \leftrightarrow \varphi \varphi] \end{aligned}$$

# FIMPy baryogenesis + EMD

arXiv:2304.07345

So far, we have placed ourselves within the simplest of cosmological scenarios : inflation was followed by an uninterrupted period of radiation domination until matter-radiation equality.

What if the history of the Universe involved additional epochs ?

Consider the case in which, at some point after inflation (*and* freeze-in), the Universe became temporarily dominated by some fluid  $X$  behaving as matter, which subsequently decayed (symmetrically) into SM particles.

Upon the decay of  $X$ , entropy is injected in the plasma and all quantities become diluted

Note also that :

- In the case of scalar DM,  $Y_{\text{DM}} \sim T_{\text{RH}}/\Lambda$ , whereas  $Y_{\text{B-LSM}} \sim T_{\text{RH}}^4/\Lambda^6$
- In the case of fermion DM,  $Y_{\text{DM}} \sim T_{\text{RH}}^3/\Lambda^4$ , whereas  $Y_{\text{B-LSM}} \sim T_{\text{RH}}^8/\Lambda^{10}$

→ The dilution process may impact dark matter and baryogenesis in different ways

# The radiation - condensate system

Once the scalar condensate decays, the relativistic degrees of freedom that are present in the Universe are diluted by

$$\Delta_{\text{EMD}} \equiv \frac{S_{\text{final}}}{S_{\text{initial}}} \approx \frac{T_{\text{dom},X}}{T_{\text{dec},X}}$$

where  $S$  are the comoving entropies of the Universe at times well before and well after the decay of  $X$ .

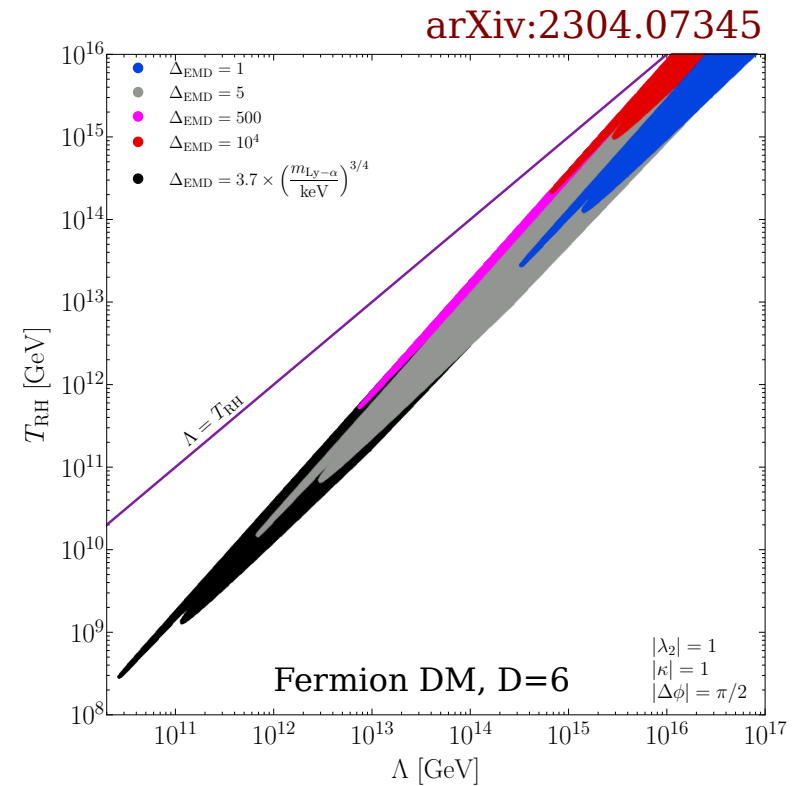
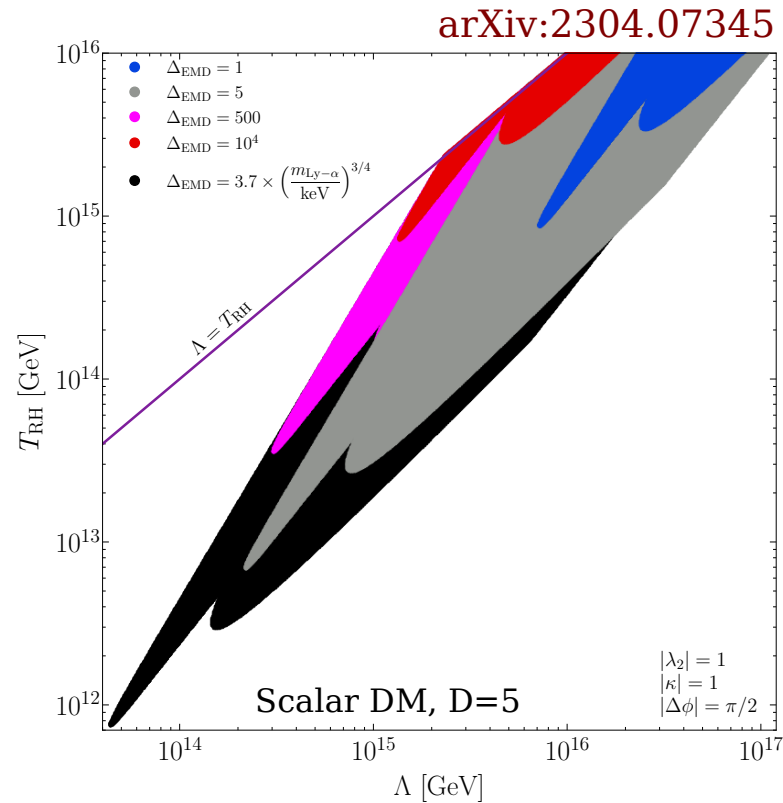
The evolution of the cosmological background is, in this case, described by the system of equations

$$\begin{aligned}\frac{d\rho_X}{d\tilde{N}} &= -3\rho_X - \frac{\Gamma_X}{H}\rho_X \\ \frac{d\rho_{\text{rad}}}{d\tilde{N}} &= -4\rho_{\text{rad}} + (1 - B_{\text{DM}})\frac{\Gamma_X}{H}\rho_X \\ \frac{d\rho_{\text{DM}}}{d\tilde{N}} &= -4\rho_{\text{DM}} + B_{\text{DM}}\frac{\Gamma_X}{H}\rho_X \\ \frac{dH}{d\tilde{N}} &= -\frac{1}{2HM_{\text{Pl}}^2} \left( \rho_X + \frac{4}{3}\rho_{\text{DM}} + \frac{4}{3}\rho_{\text{rad}} \right)\end{aligned}$$

where  $dN = d(\ln a) = Hdt$ , and we assume  $B_{\text{DM}} = 0$  (*i.e.* no  $X$  decays into DM).

# Results

All in all, the mechanism works! Dark matter production and baryogenesis can be simultaneously achieved.



- Depending on dilution size, can live with a wide range of reheating temperatures.
- In the zero-dilution case, DM again predicted to be close to the Lyman- $\alpha$  bound, but can reach the multi-MeV range in the presence of dilution.
- However, this scenario is extremely challenging phenomenologically...



# EFTs, dilution and inflation

A modified cosmological history could have an impact on inflationary observables, most notably the spectral index  $n_s$  and the tensor-to-scalar ratio  $r$ . We will focus on the former. The general idea goes as follows :

- Each of the microscopic toy models that we have considered favours, for a given dilution size, a region of the  $(\Lambda, T_{\text{RH}}, m_{\text{DM}})$  parameter space.
- All of these quantities allow us to compute the observable number of e-folds  $N_*$

$$N_* \equiv \int_{t_*}^{t_{\text{end}}} dt H = \ln(a_{\text{end}}/a_*),$$

Which, in presence of and EMD phase is shifted wrt its “thermal” value as

$$N_* = N^{(\text{th})} - \delta N_* \approx N^{(\text{th})} - \frac{1}{3} \ln(\Delta_{\text{EMD}})$$

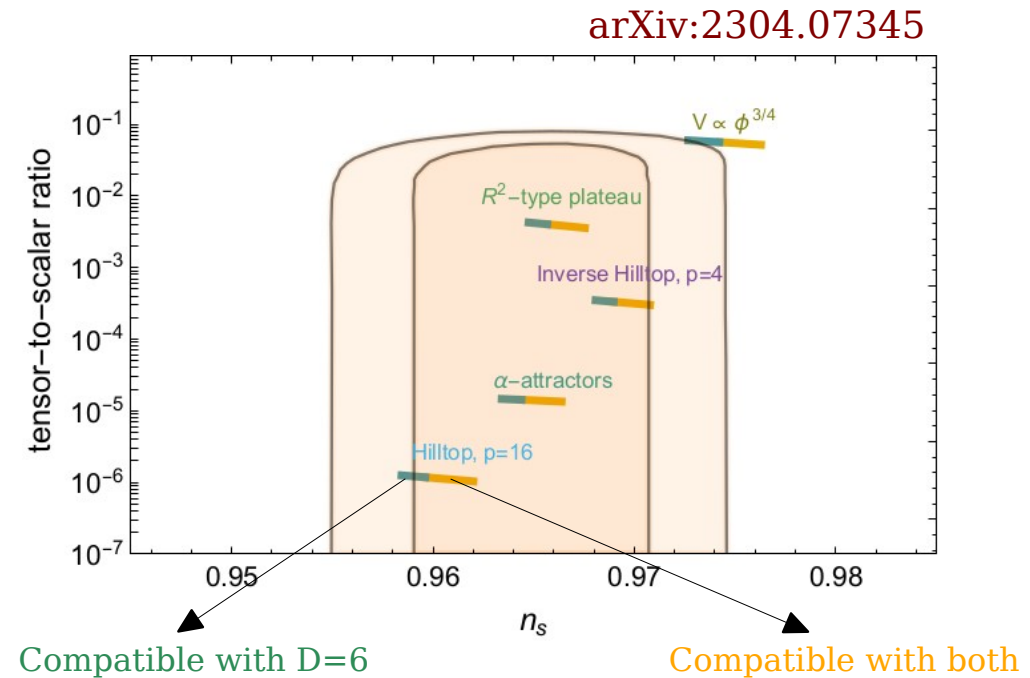
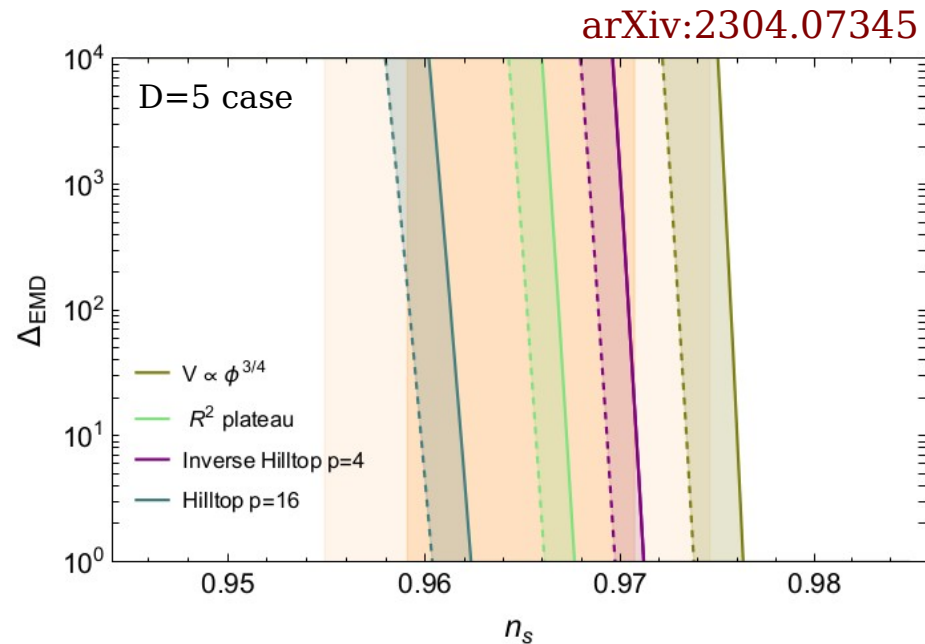
which means that our viable parameter space can be recasted in terms of  $N_*$ .

- Lastly,  $N_*$  can be related, within specific models of inflation, with the spectral index through a relation of the type  $n_s = n_s(N)$ .

In other words, given an inflationary model, we can draw conclusions on the microscopic scenarios which are favoured and vice-versa

# Potential imprints on the CMB ?

A dedicated analysis for a set of representative inflationary models yields the following results



- EUCLID/21-cm surveys could reach a  $10^{-3}$  precision in the measurement of  $n_s$ .
- CMB Stage-4 experiments could detect  $r > 0.003$ .
- Is it possible to fully test the UV FIMPy baryogenesis scenario ?  $\rightarrow$  No

$\rightarrow$  But it may become possible to use CMB observables for model selection.

# Summary and outlook

- There is no *a priori* reason why the observed dark matter abundance and the matter-antimatter asymmetry of the Universe should admit a common explanation.
- However, it is a possibility. And a much welcome one! This is the reason why such an option has been entertained since quite a few years and in the context of different DM generation mechanisms (asymmetric DM, freeze-out, freeze-in).
- Freeze-in production of DM, in particular, constitutes an interesting playground for baryogenesis, since it incorporates from the start one of the three Sakharov conditions: out-of-equilibrium dynamics.
- “Freeze-in baryogenesis” can work in wildly different contexts: asymmetric dark matter, symmetric dark matter that is mostly produced in the IR, UV freeze-in. It can give rise to interesting signals at the LHC and Cosmology.
- Once embedded within concrete inflationary scenarios, models which are otherwise extremely hard to test can give rise to observable predictions.

Thank you!