



Winter school analogue gravity/cosmology in Benasque in January 2026!

Experiments with fluids of light

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LKB 08/11/2023

In a (quantum) fluid

Fluid velocity $\mathbf{v} = (\hbar/m)\nabla\phi$

Speed of sound $c_s \propto \sqrt{\frac{g\rho_0}{m}}$

m – mass
 g – interaction constant
 ρ_0 – density

Wave eq for collective excitations of (super)fluid $\psi = \psi_0 + \epsilon_1\psi_1$

$$-\partial_t\left(\frac{\rho_0}{c_s^2}(\partial_t\rho_1 + \mathbf{v}_0\nabla\rho_1)\right) + \nabla\left(\rho_0\nabla\rho_1 - \frac{\rho_0\mathbf{v}_0}{c_s^2}\partial_t\rho_1 + \mathbf{v}_0\nabla\rho_1\right) = 0$$

Relativistic form of wave eq for collective excitations: $\Delta\rho_1 = \frac{1}{\sqrt{-\eta}}\partial_\mu(\sqrt{-\eta}\eta^{\mu\nu}\partial_\nu\rho_1) = 0$

with $\eta_{\mu\nu} = \begin{pmatrix} -(c_s^2 - \mathbf{v}_0^2) & -v_o^x & -v_o^y \\ -v_o^x & 1 & 0 \\ -v_o^y & 0 & 1 \end{pmatrix}$

Motion of collective excitations in inhomogeneous fluid flow \leftrightarrow scalar field on curved spacetime

Control parameters: $\mathbf{v}_0, \mathbf{c}_s$

In a (quantum) fluid

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m – mass
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Possible geometries with $\eta_{\mu\nu} = \begin{pmatrix} -(c_s^2 - \mathbf{v}_0^2) & -v_o^x & -v_o^y \\ -v_o^x & 1 & 0 \\ -v_o^y & 0 & 1 \end{pmatrix}$

(i) accelerating flow along 1 spatial dimension → static 1D spacetime

Horizon where $v_0 = c_s$

(ii) radially accelerating flow in 2 spatial dimensions → static spherically symmetric 2D spacetime

Horizon where $v_r = c_s$

(iii) radially and azimuthally accelerating flow in 2 spatial dimensions → static rotating spacetime

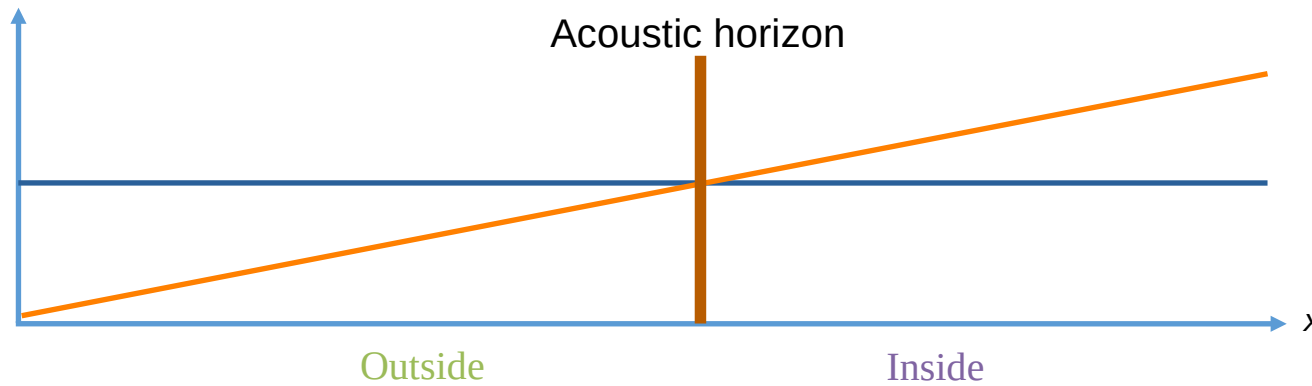
Horizon where $v_r = c_s$

Ergosurface where $|\mathbf{v}_0| = c_s$

Static 1D geometry \leftrightarrow waterfall geometry

Flow velocity of fluid v

Speed of sound c



Quantised acoustic field:

$$\text{in: } \phi = \int d\omega (a_\omega f_\omega + a_\omega^\dagger f_\omega^*) \quad a |0\rangle = 0$$

Express out modes in terms of in modes:

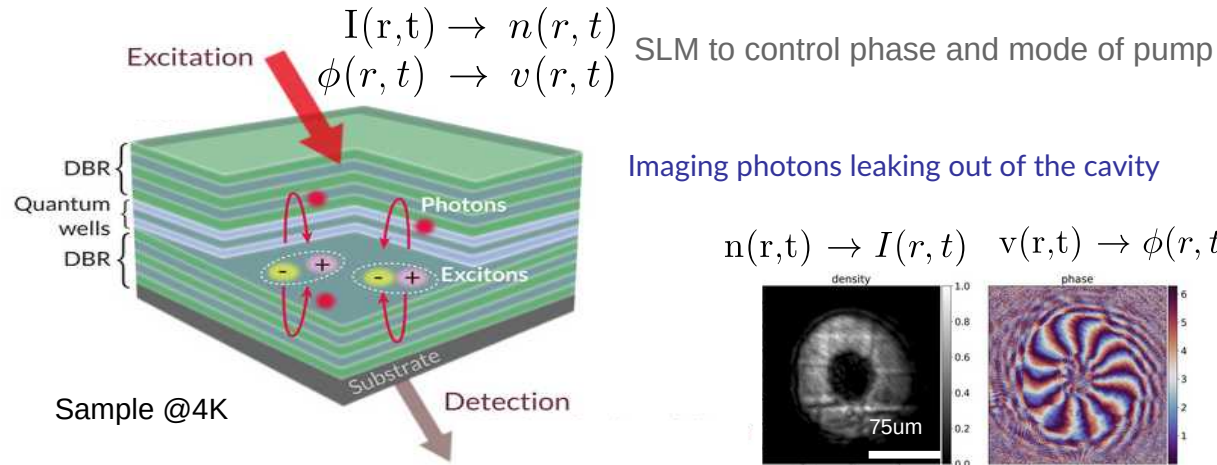
$$\text{out: } \phi = \int d\omega (\bar{a}_\omega F_\omega + \bar{a}_\omega^\dagger F_\omega^*) \quad \bar{a} |\bar{0}\rangle = 0$$

$$F_\omega = \int d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^*)$$

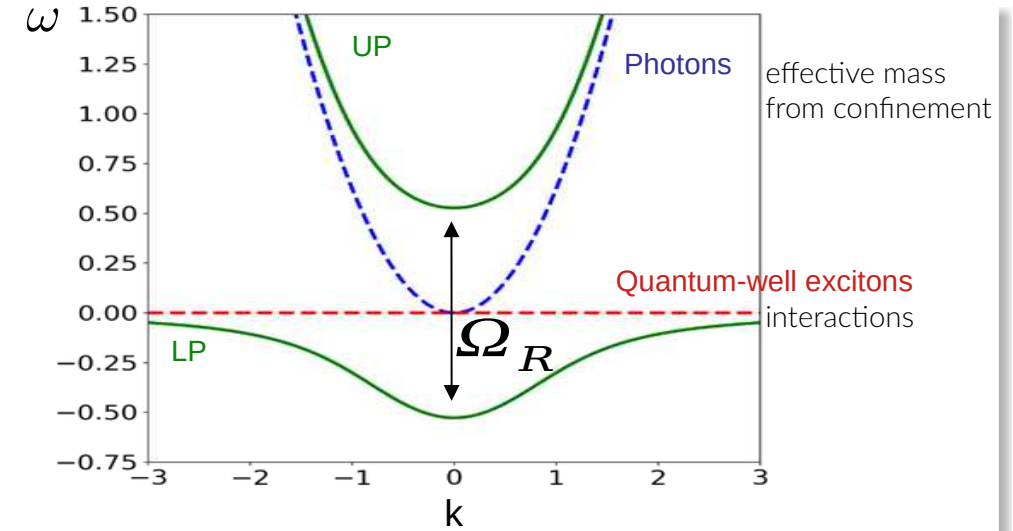
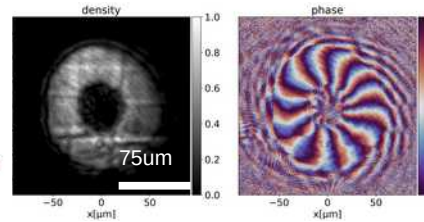
Different speeds on either side of the horizon $\Rightarrow |\bar{0}\rangle \neq |0\rangle \Rightarrow \beta_{\omega\omega'} \neq 0$

Mixing of positive and negative frequency waves \Rightarrow mixing of creation and annihilation operators

$$a |\bar{0}\rangle = \sum_{\omega'} \beta_{\omega\omega'} |\bar{1}\rangle > 0$$



$$n(r,t) \rightarrow I(r,t) \quad v(r,t) \rightarrow \phi(r,t)$$



Polaritons= photons dressed with material excitations that live in the cavity plane

Dynamics in the cavity plane described by Gross-Pitaevskii (Nonlinear Schrödinger) equation:

$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2 \nabla^2}{2m_{LP}^*} + gn \right) \psi - \frac{i\hbar\gamma}{2} \psi + P(r,t)$$

Driven-dissipative dynamics \rightarrow Out-of-equilibrium system

- g polariton-polariton interaction constant
- γ Losses
- P pump

$$\text{GPE: } i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2 \nabla^2}{2m_{LP}^*} + gn \right) \psi - \frac{i\hbar\gamma}{2} \psi + P(r, t)$$

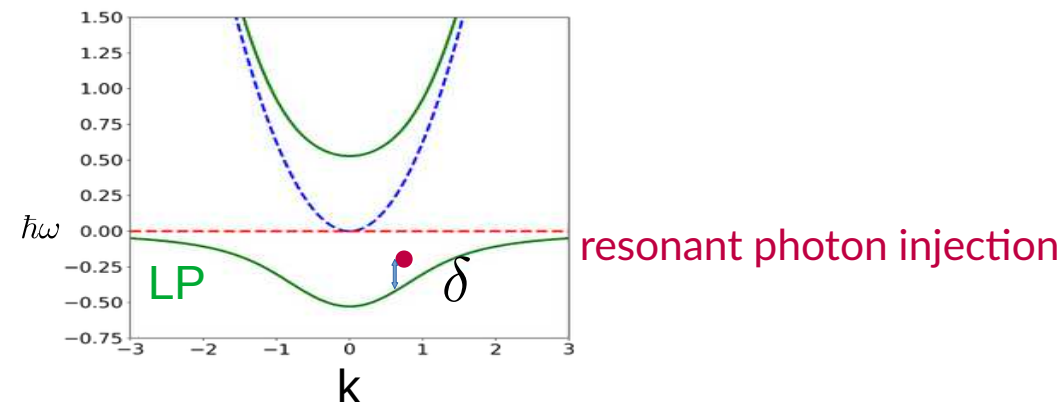
Bogoliubov theory:

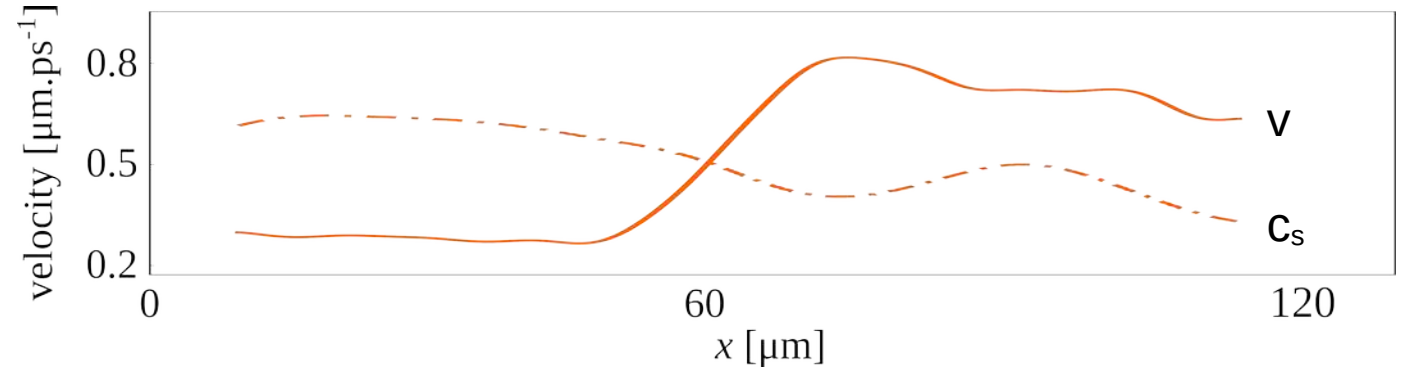
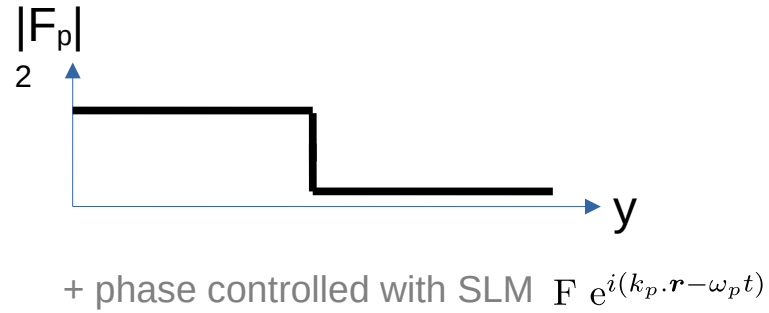
1. Linearise GPE around steady-state solution $\psi(r, t) = \psi_0(r, t) + \delta\psi(r, t)$

2. Equation of motion of weak perturbations
$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} \delta\psi(r, t) \\ \delta\psi^*(r, t) \end{pmatrix} = L_{\text{Bog}} \begin{pmatrix} \delta\psi(r, t) \\ \delta\psi^*(r, t) \end{pmatrix}$$

3. Eigenvalues of L_{Bog} == dispersion relation

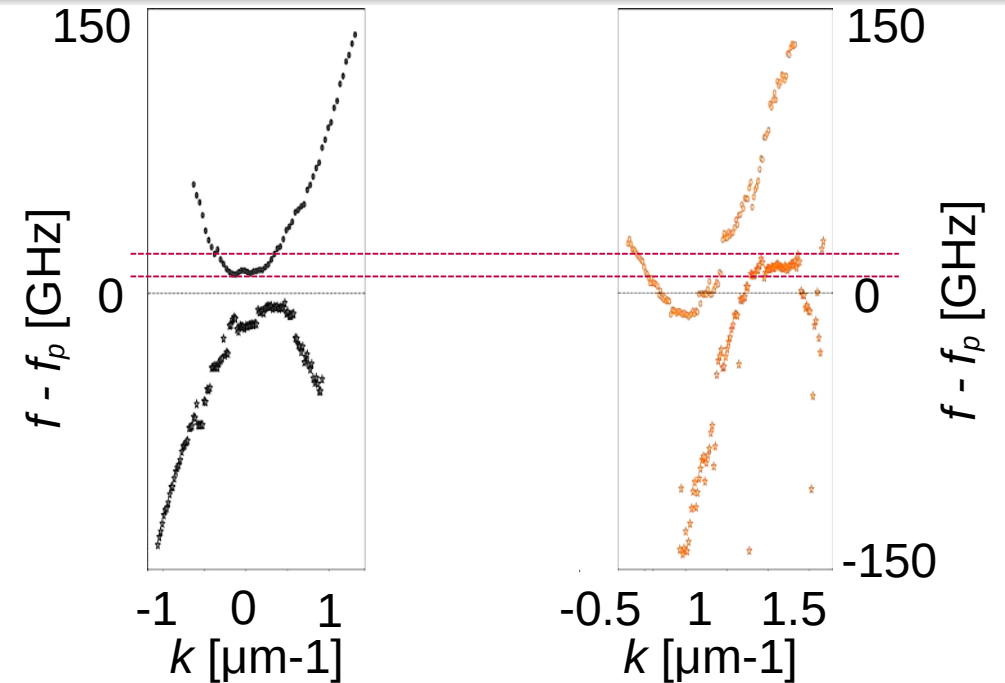
$$\hbar\omega(k) = \pm \sqrt{\left(\frac{\hbar^2 k^2}{2m} - \delta + 2gn \right)^2 - (gn)^2} - i\gamma$$

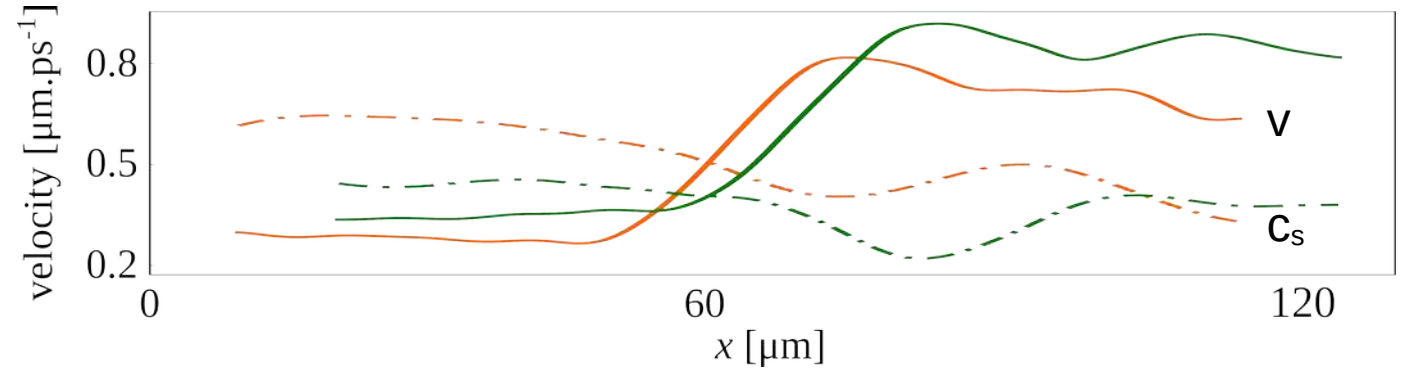
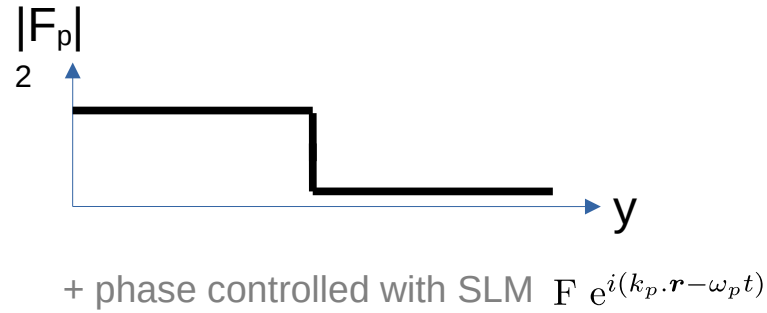




Spectrum of collective excitations

$$\omega^\pm(k) - v(x) \cdot k(x) = \frac{i\gamma}{2} \pm \sqrt{\left(\frac{\hbar k(x)^2}{2m} - \delta(k_p) + 2gn(x)\right)^2 - (gn(x))^2}$$





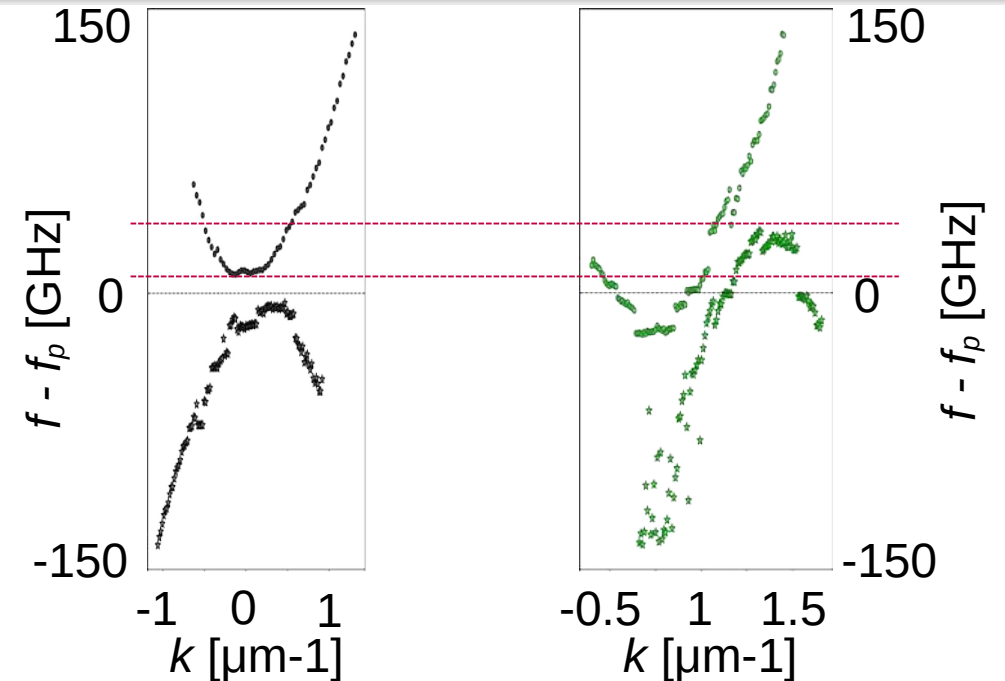
Spectrum of collective excitations

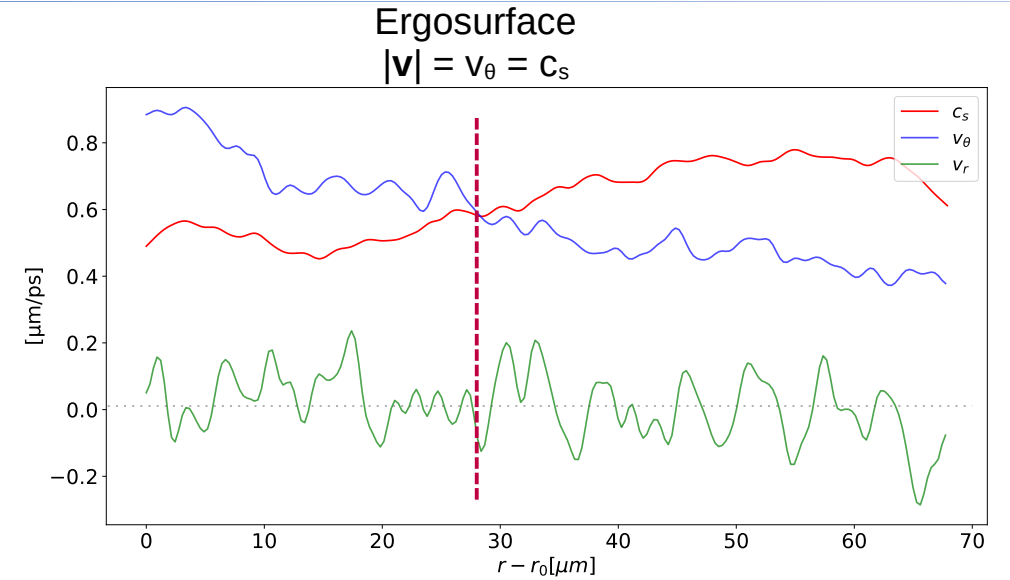
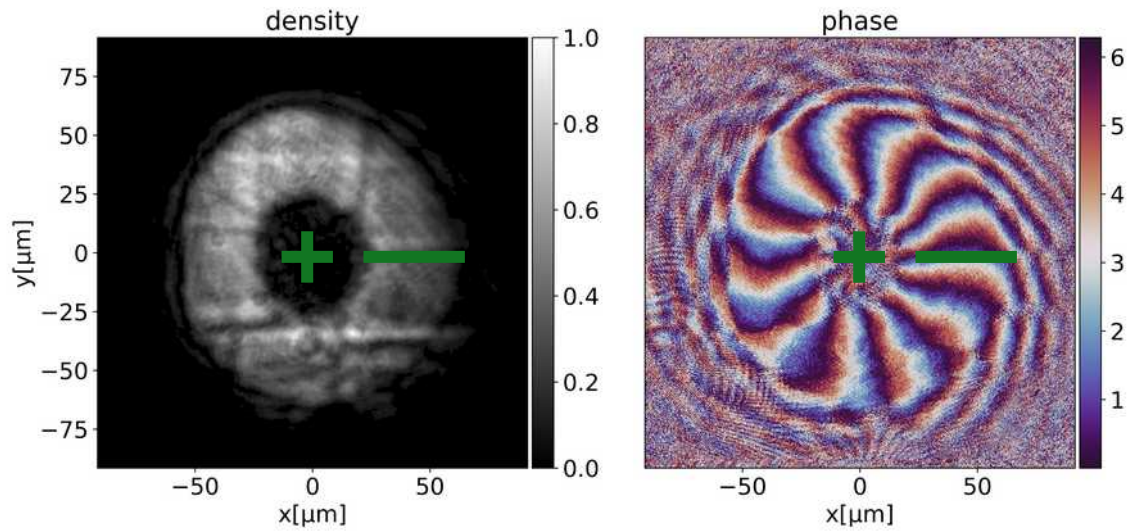
$$\omega^\pm(k) - v(x) \cdot k(x) = \frac{i\gamma}{2} \pm \sqrt{\left(\frac{\hbar k(x)^2}{2m} - \delta(k_p) + 2gn(x)\right)^2 - (gn(x))^2}$$

Mach number: $M = c_s / v$

Horizon steepness: $\partial M / \partial x|_{x_H}$

orange $0.08 \mu\text{m}^{-1}$
green $0.04 \mu\text{m}^{-1}$



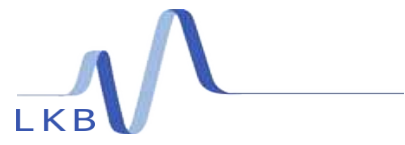


$$\nabla \phi_{SLM} = \frac{C}{r} \mathbf{u}_\theta - \frac{D}{r} \mathbf{u}_r$$



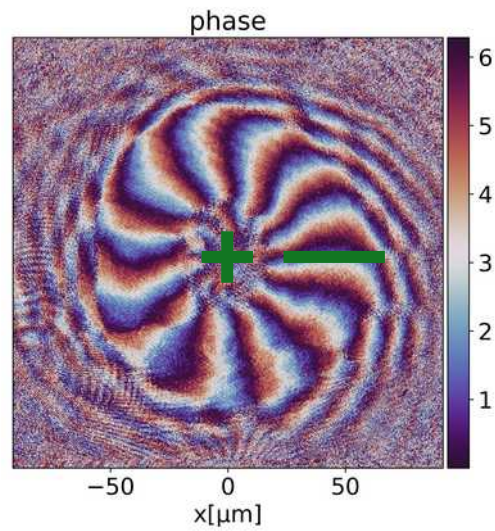
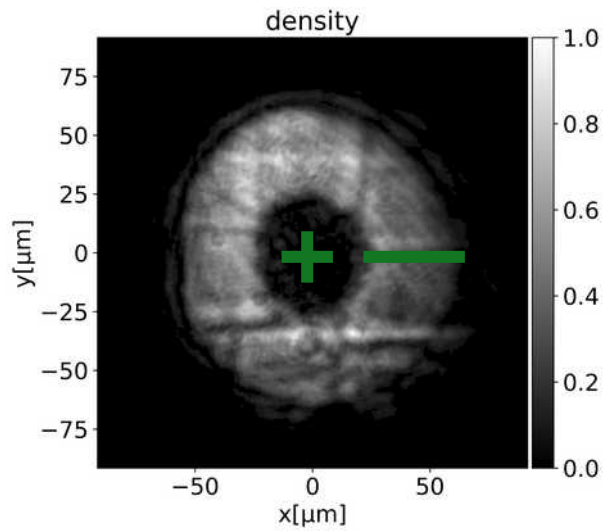
$$\phi_{SLM} = C\theta - D \ln(r)$$

$$C = 12, D = 0$$



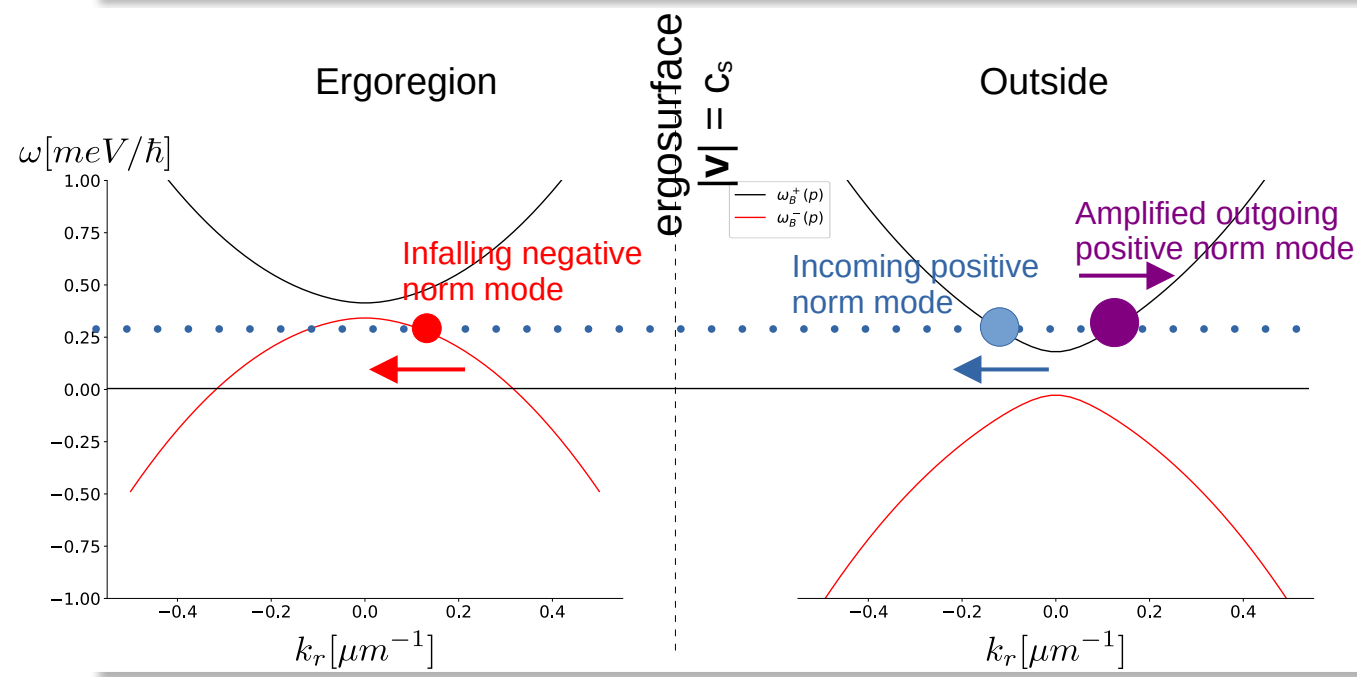
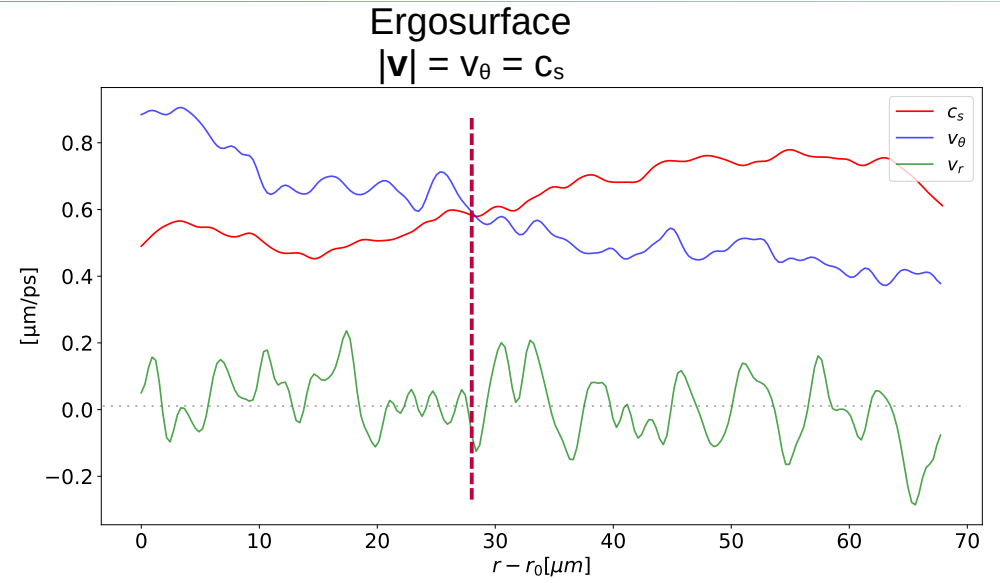
Rotating geometry $D = 0, C = 12$

Data by Killian Guerrero arXiv:2310.16031

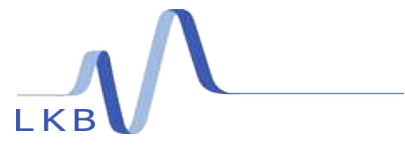


$$c_s \propto \sqrt{gn}$$

$$v_{fluid} \propto \nabla \phi$$

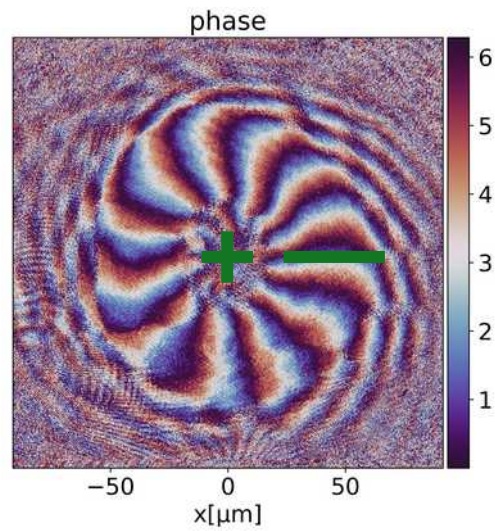
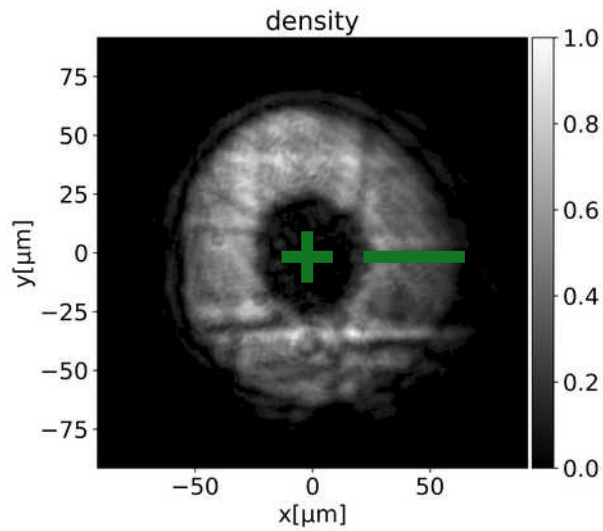


$$Q(\phi) = i \int_0^\infty \int_0^{2\pi} d\theta \frac{r}{c^2} (\phi^* \partial_t \phi - \partial_t \phi^* \phi)$$



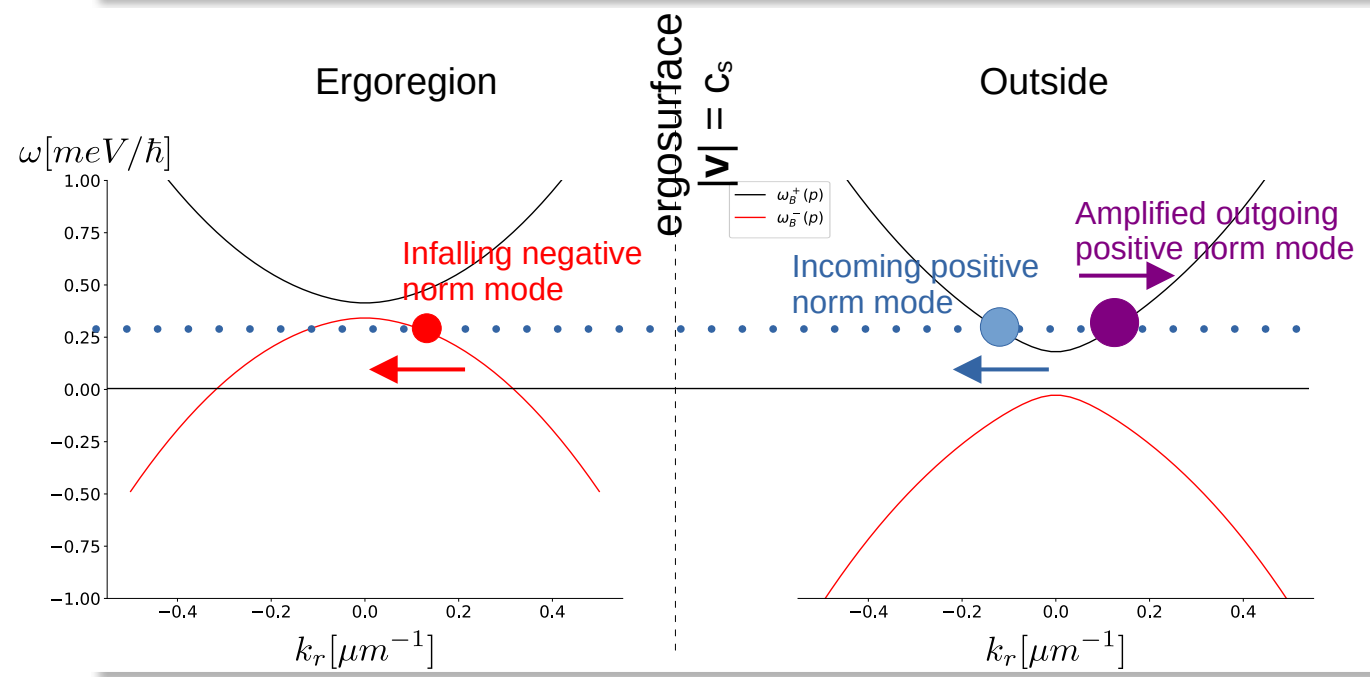
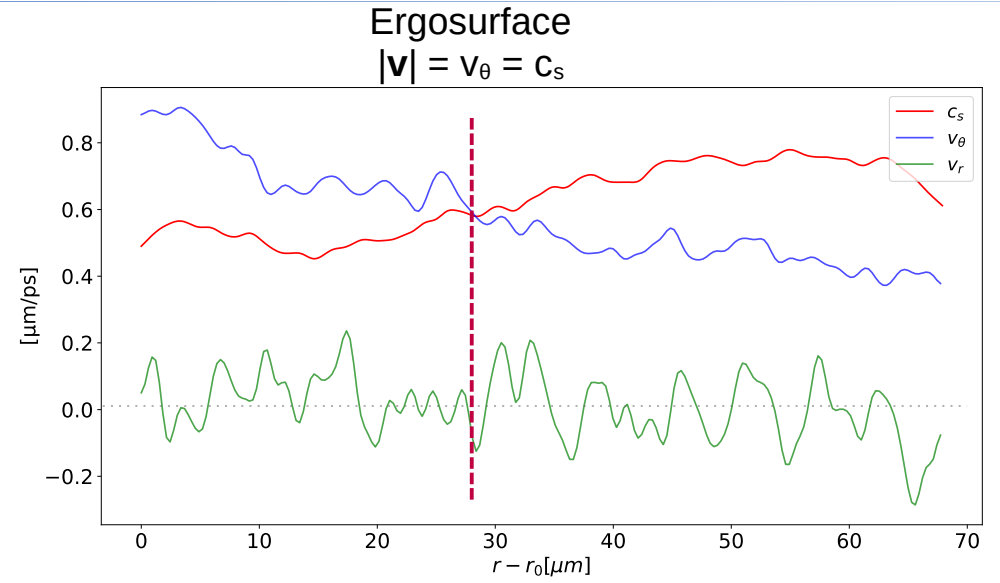
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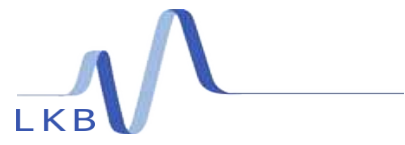
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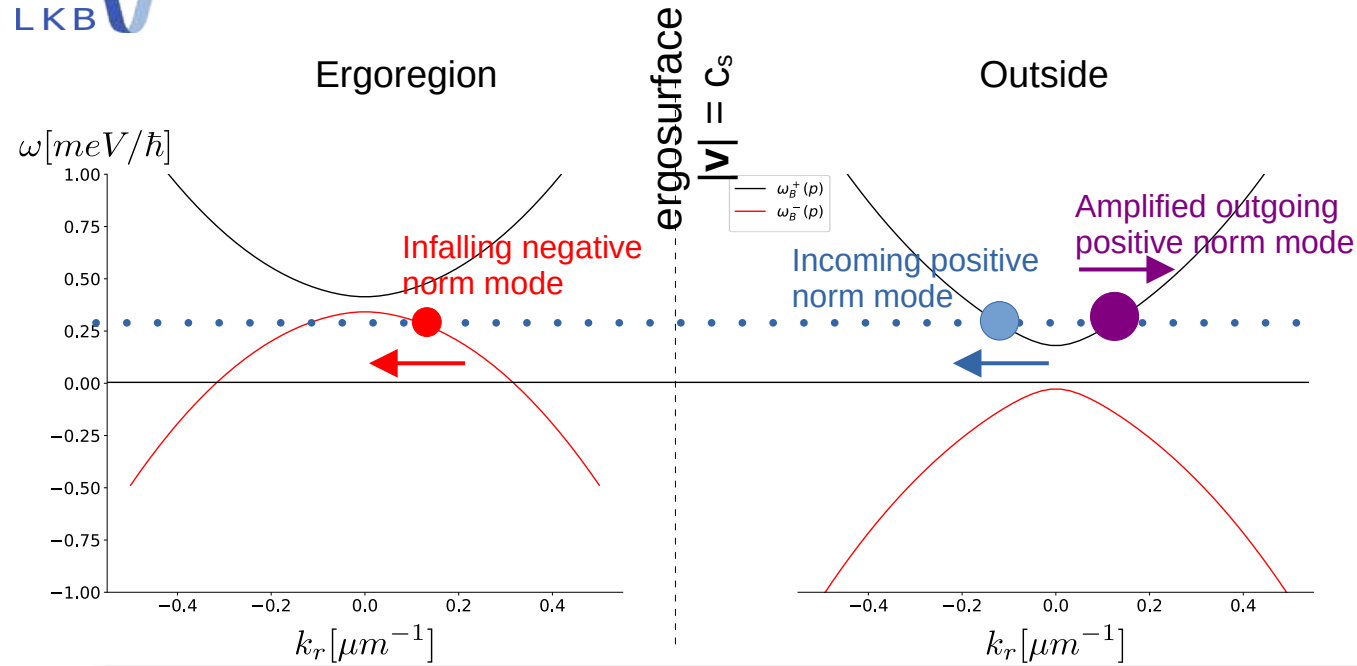
Stable ergosurface without a horizon?

$$v_g = 1 \mu m \cdot ps^{-1} \quad \tau = 15 ps$$



Entanglement in rotating geometry

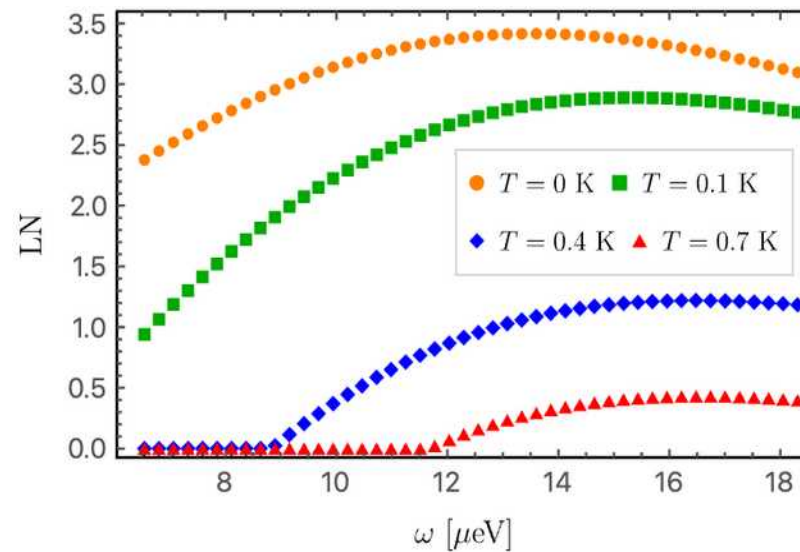
Calculations by Aria Delhom arXiv:2310.16031

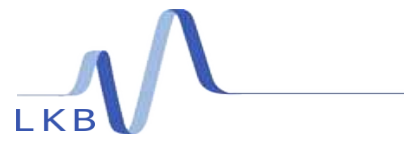


$$(W_r^{in}, W_l^{in}) \rightarrow (W_r^{out}, W_r^{out})$$

$$B_{\omega l} = \begin{pmatrix} T_{\omega l} & r_{\omega l} \\ R_{\omega l} & t_{\omega l} \end{pmatrix}$$

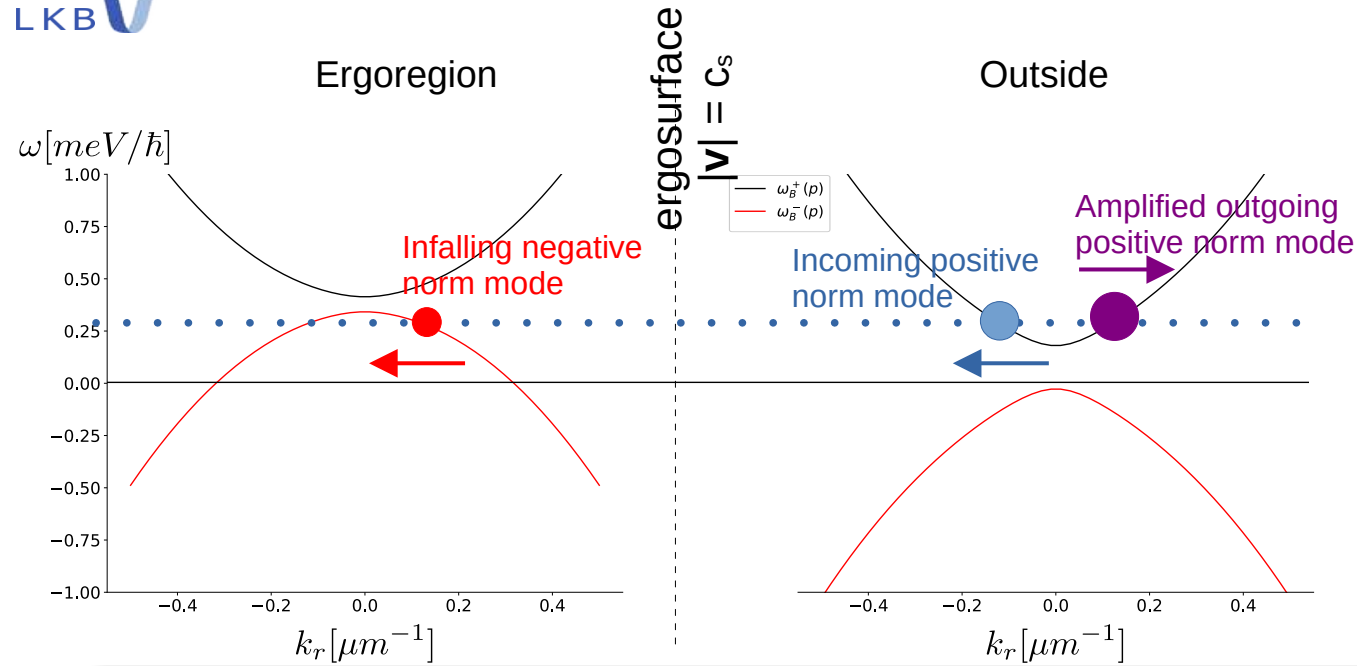
Vacuum + thermal $|0\rangle \otimes |Th\rangle$
noise at input





Entanglement in rotating geometry

Calculations by Aria Delhom arXiv:2310.16031

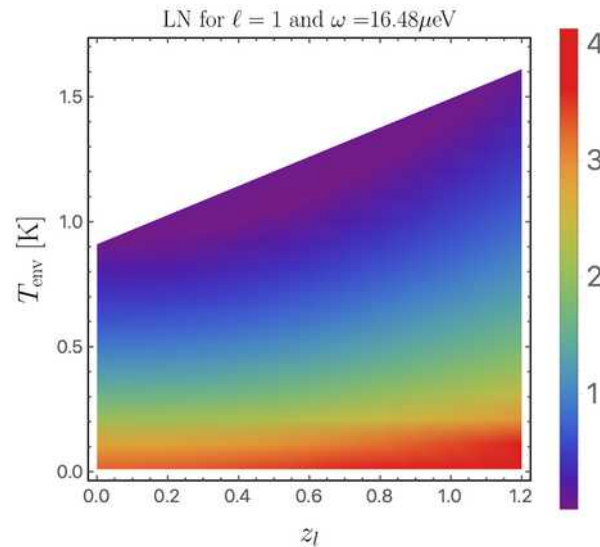


$$(W_r^{in}, W_l^{in}) \rightarrow (W_r^{out}, W_r^{out})$$

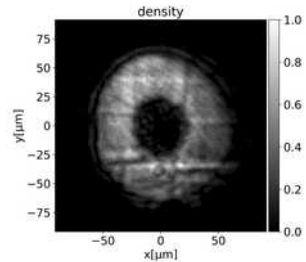
$$B_{\omega l} = \begin{pmatrix} T_{\omega l} & r_{\omega l} \\ R_{\omega l} & t_{\omega l} \end{pmatrix}$$

Vacuum + thermal noise + single mode squeezed state at input

$$|0\rangle \otimes |Th\rangle \otimes |SMST\rangle$$



Where do we go from here?

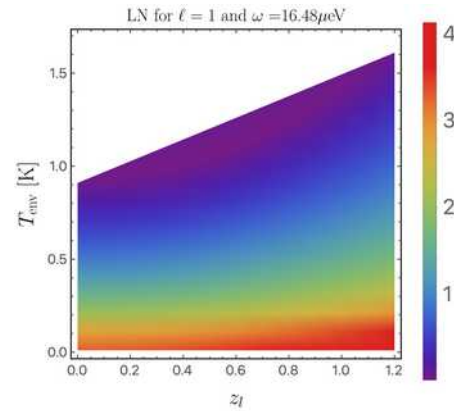


Experiments with polaritons

- Measure Hawking radiation and rotational superradiance independently from one another
- Measure interplay between the two effects → modification of correlations?

Static 1D geometry: arXiv:2311.01392

Rotating geometry: arXiv:2310.16031



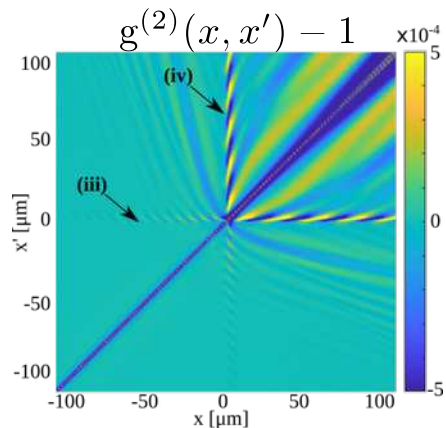
All optical experiments

- Measure phase and density → access full field statistics and dynamics
- High resolution spectroscopy in 1 and 2D with and without rotation
- Homodyne detection to enhance signal strength and measure quantum correlations
- Enhance strength of emission and degree of entanglement by probing with squeezed state

F Claude *et al* PRL **129** 103601 2022,

PRB **107** 174507 2023

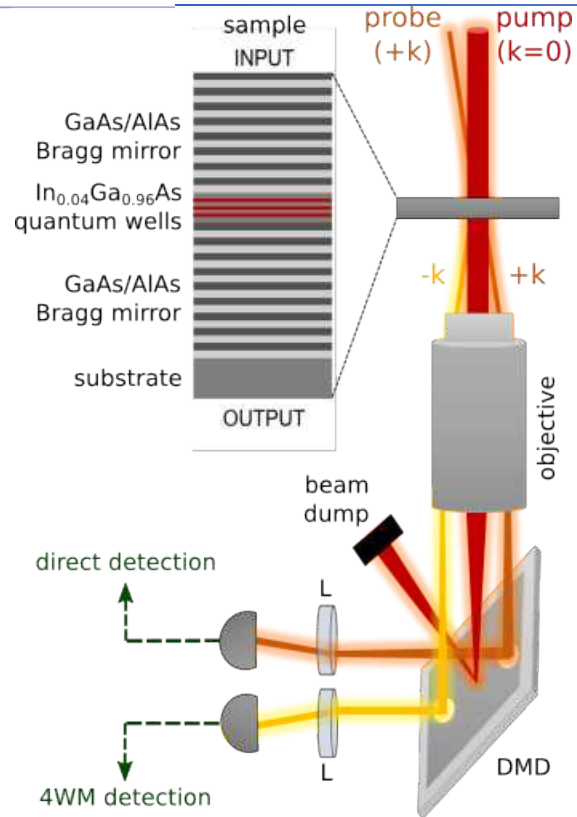
I Agullo *et al* PRL **128** 091301 2022



Numerical simulations

- New effect of quantum fields predicted: vacuum excitation of quasi-normal mode of acoustic field
- Good experimental configuration to observe strong correlations

Jacquet *et al.* PRL **130** 2023, EPJD **76** 2022

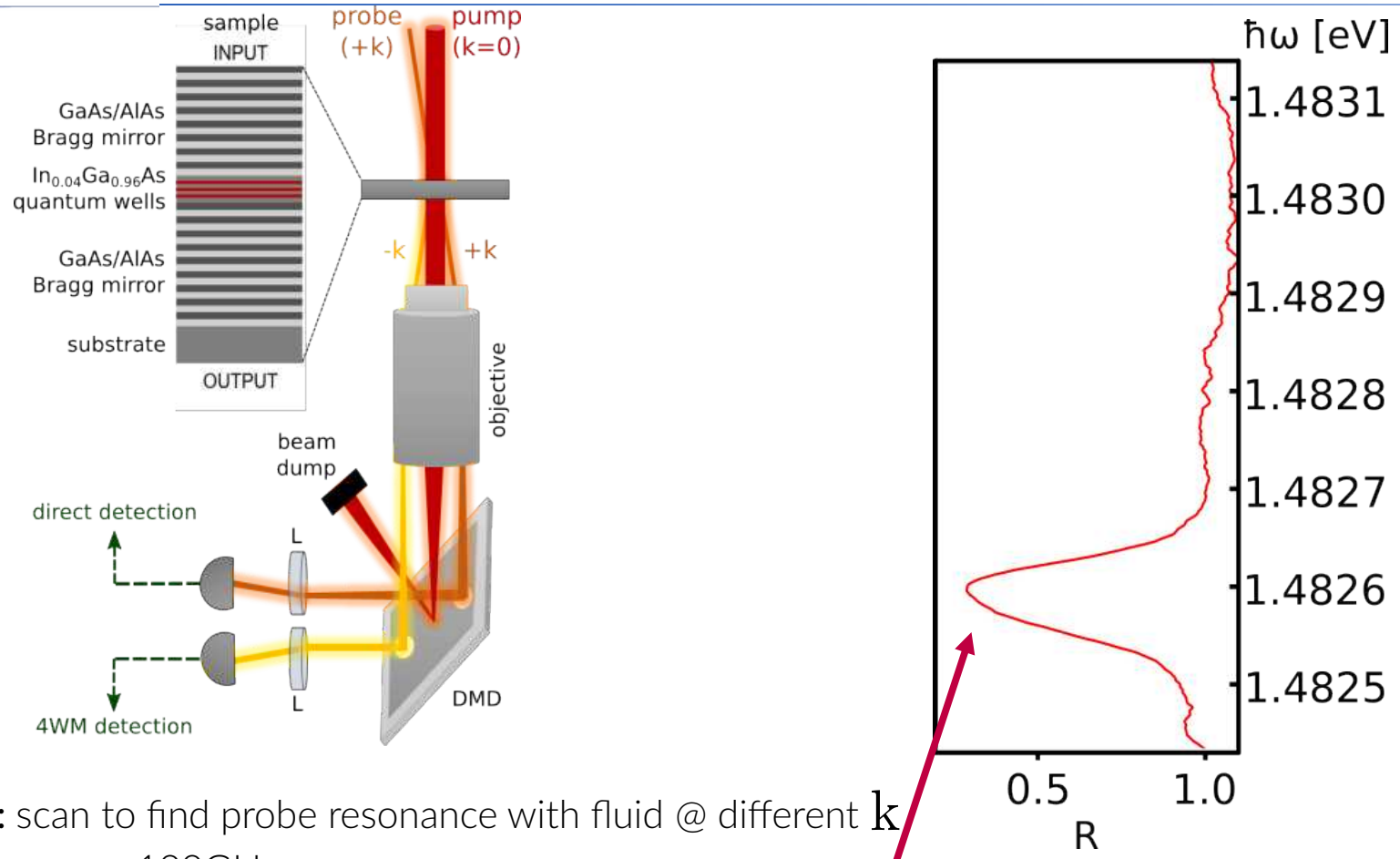


Probe: scan to find probe resonance with fluid @ different \mathbf{k}
 energy scan: $\sim 100\text{GHz}$

Detection: probe resonates with fluid \rightarrow transmission = dip in reflectivity

ω resolution fixed by the probe laser linewidth ($< 250\text{kHz}$)

\mathbf{k} resolution fixed by the k -space filtering ($0.02\mu\text{m}^{-1}$)

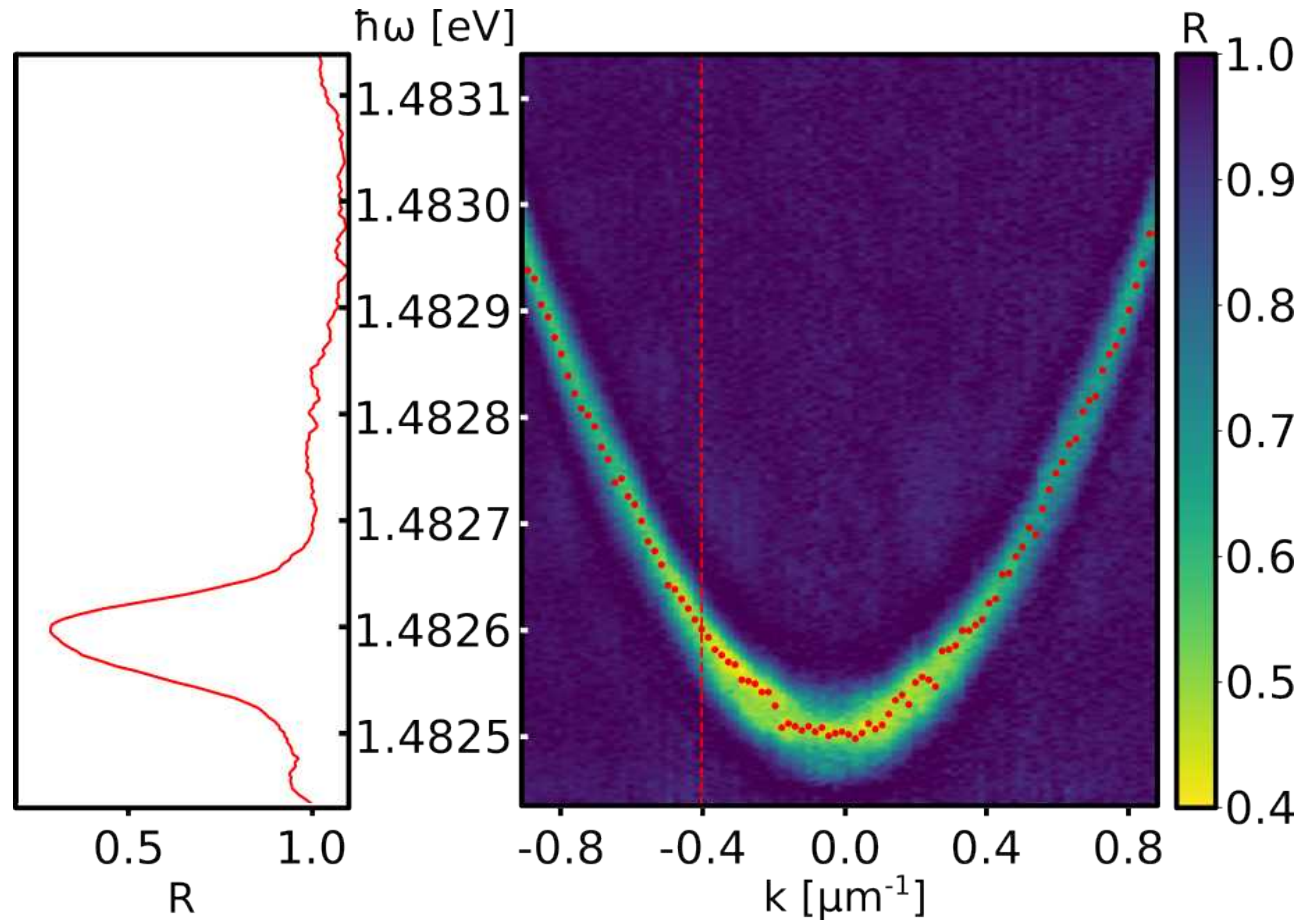
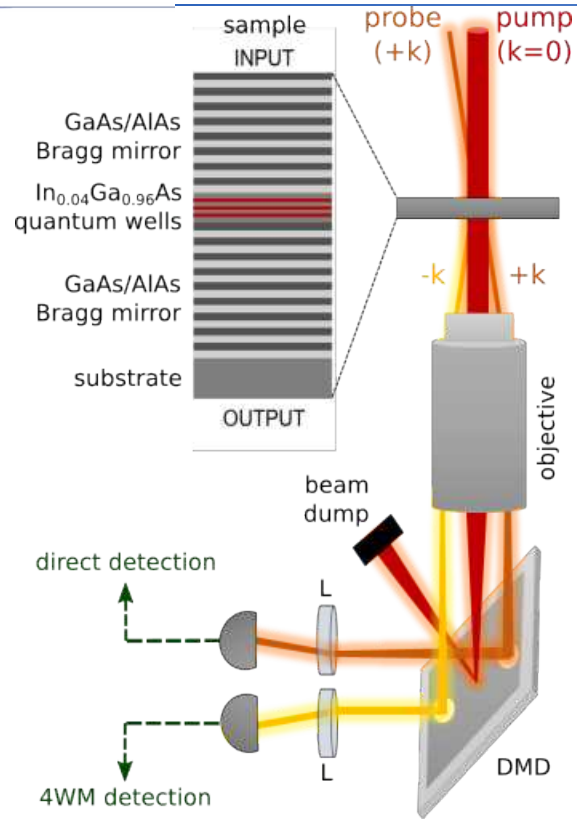


Reflectivity of probe @ k

Probe: scan to find probe resonance with fluid @ different k
 energy scan: $\sim 100\text{GHz}$

Detection: probe resonates with fluid \rightarrow transmission = dip in reflectivity

- ω resolution fixed by the probe laser linewidth ($< 250\text{kHz}$)
- k resolution fixed by the k -space filtering ($0.02\mu\text{m}^{-1}$)



Probe: scan to find probe resonance with fluid @ different k
 energy scan: $\sim 100\text{GHz}$

Detection: probe resonates with fluid \rightarrow transmission = dip in reflectivity

Coherent probe spectroscopy: Reflectivity map of probe vs (k, ω) \leftrightarrow spectrum of collective excitations

- ω resolution fixed by the probe laser linewidth ($< 250\text{kHz}$)
- k resolution fixed by the k -space filtering ($0.02\mu\text{m}^{-1}$)