

Neutrino Cosmology in 2023

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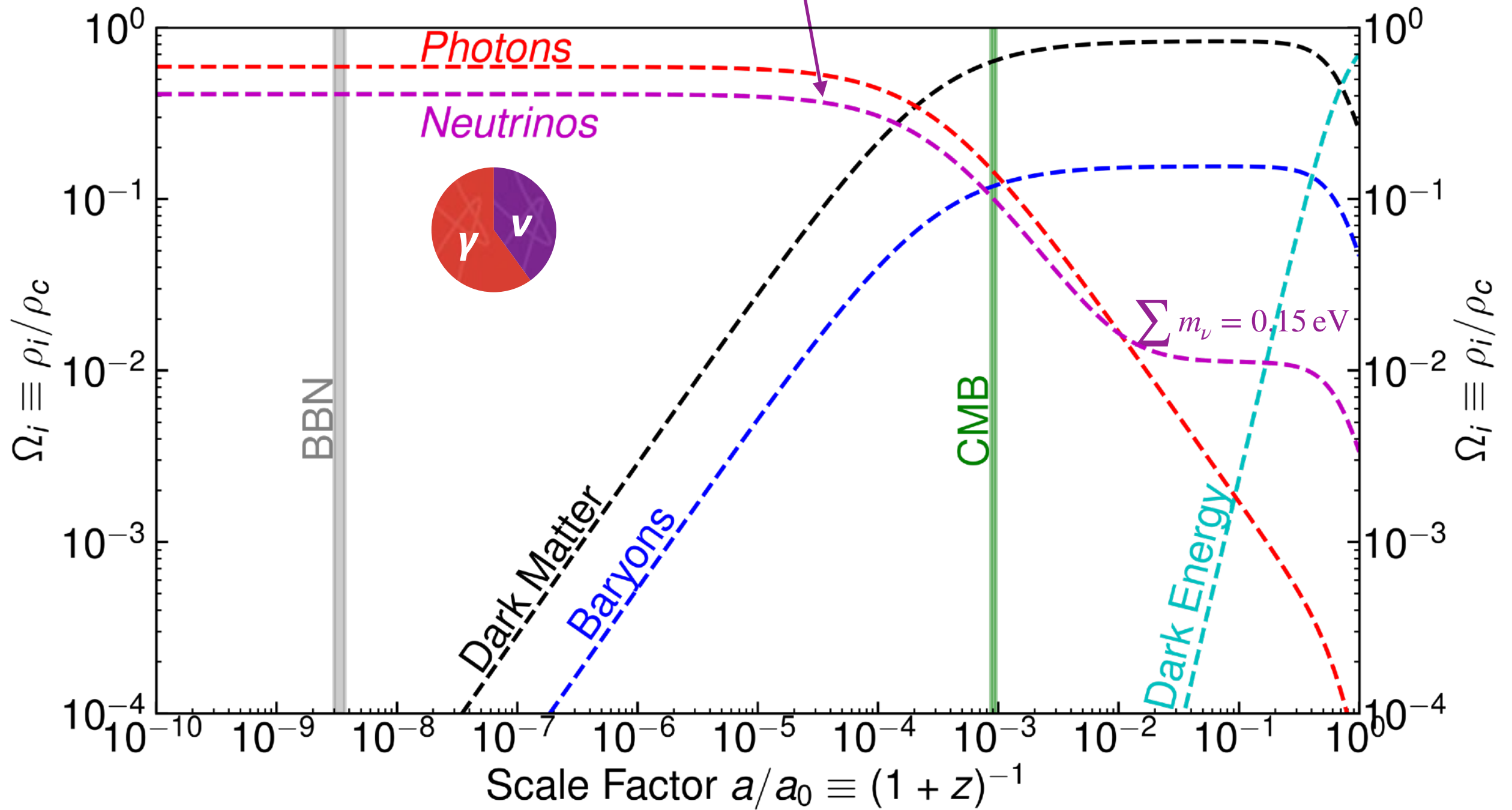


IJCLab Orsay

05-12-2023

Neutrino Evolution

Neutrinos are always a relevant species in the Universe's evolution

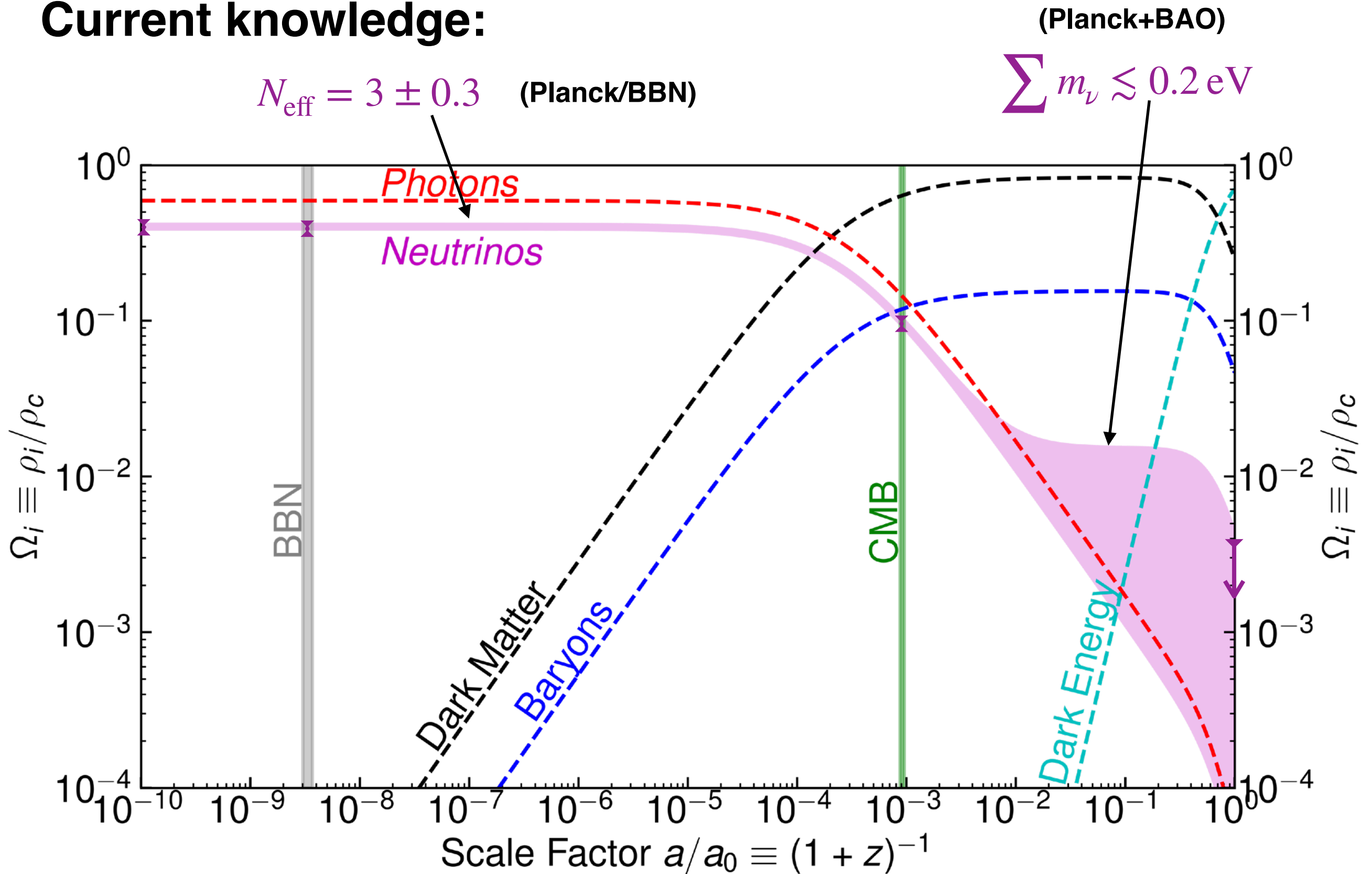


Non-Rel: $z_\nu^{\text{non-rel}} \simeq 110 \frac{m_\nu}{0.06 \text{ eV}}$

Hot DM: $\Omega_\nu h^2 = \sum m_\nu / (93.14 \text{ eV})$

Global Perspective

Current knowledge:



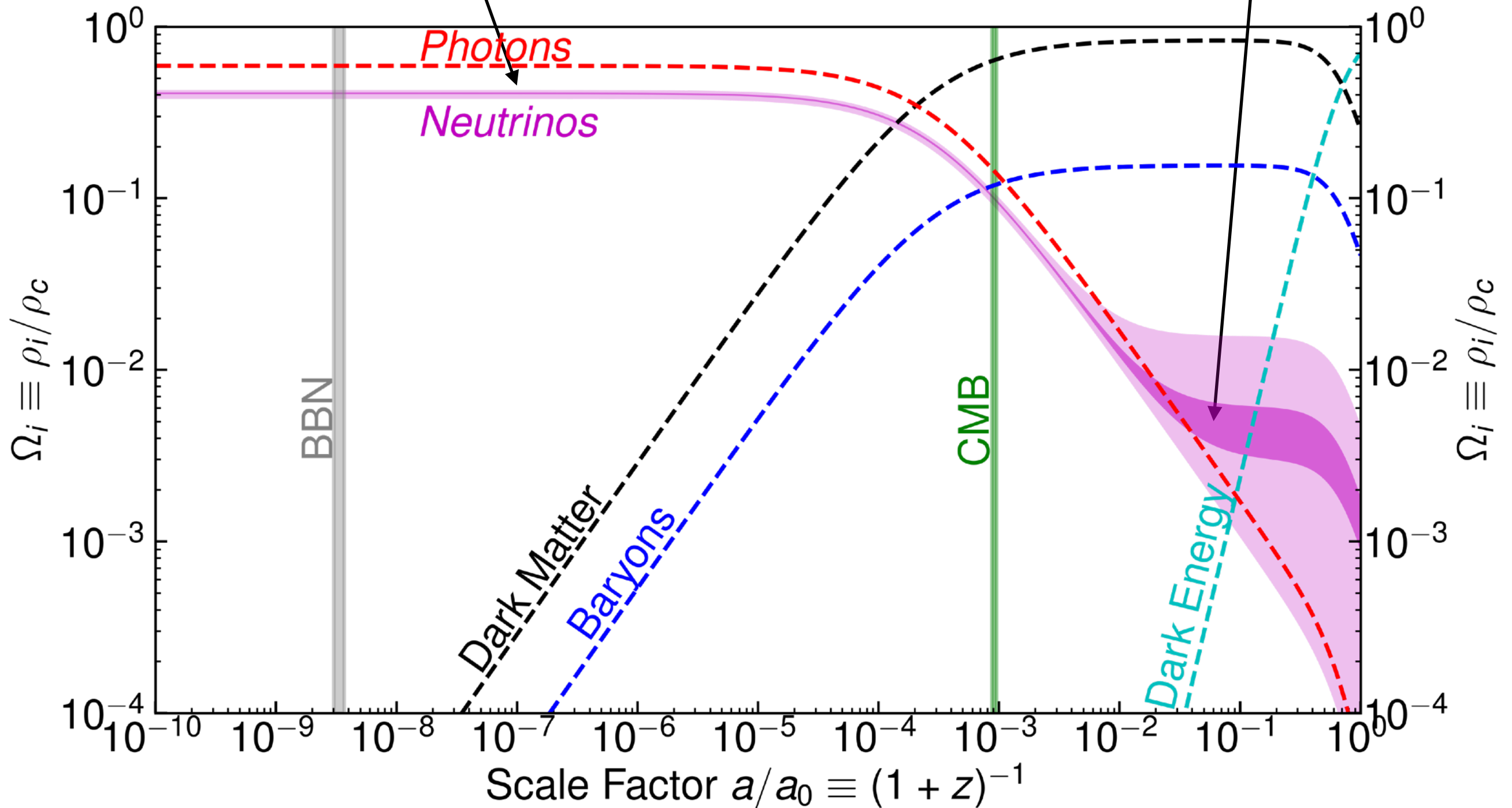
Global Perspective

In the next 5-6 years:

(DESI/Euclid + Planck)

$$N_{\text{eff}} = 3.043 \pm 0.06 \quad (\text{Simons Observatory})$$

$$\sum m_\nu = 0.06 \pm 0.02 \text{ eV}$$



Motivation/Outline

Neutrino Cosmology is about to enter the ultrahigh precision regime

1) Understand with high accuracy N_{eff} in the Standard Model

M.E.A. [1812.05605](#) & [2001.04466](#) [JCAP]

with Cielo, Mangano & Pisanti [2306.05460](#) [PRD]

2) Understand the model dependence of the bounds on $\sum m_\nu$ from Cosmology

Particularly given the strong complementarity with laboratory experiments:

$$\text{Planck: } \sum m_\nu < 0.12 \text{ eV}$$

$$\text{KATRIN 2023: } \sum m_\nu < 2.4 \text{ eV}$$

$$\text{KATRIN 2027: } \sum m_\nu < 0.6 \text{ eV?}$$

$$\sum m_\nu \Big|_{\text{NO}} \geq 0.06 \text{ eV} \quad \sum m_\nu \Big|_{\text{IO}} \geq 0.1 \text{ eV}$$

Exciting $0\nu\beta\beta$ program

with Alvey & Sabti [2111.12726](#) [JCAP]

with Schwetz & Terol-Calvo [2211.01729](#) [JHEP]

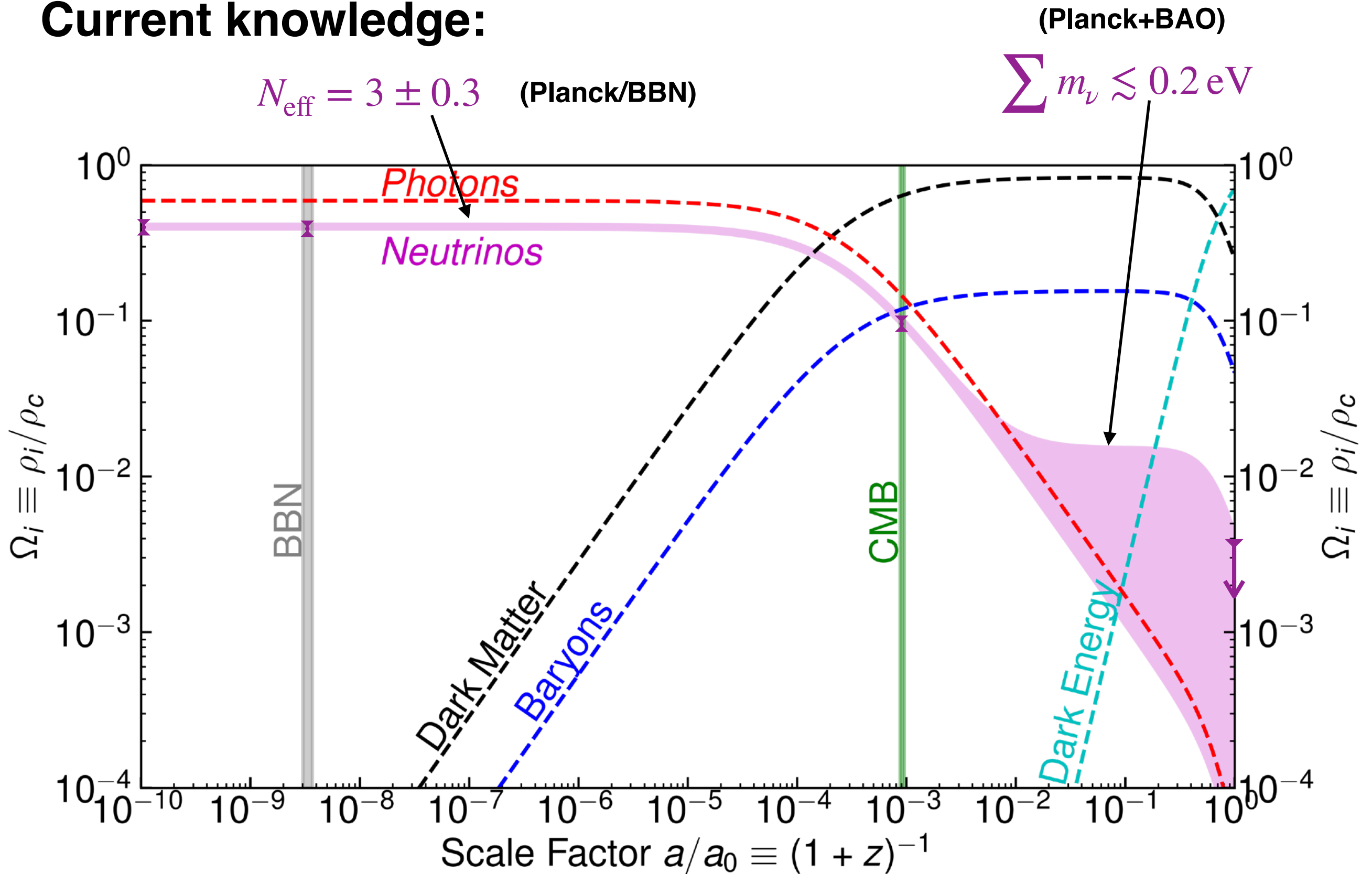
3) Understand to which degree can neutrinos have BSM interactions

with Witte [1909.04044](#), [2103.03249](#) & with Sandner [2305.01692](#) [EPJC]

with Taule & Garny [2207.04062](#) [PRD]

Outline

Current knowledge:



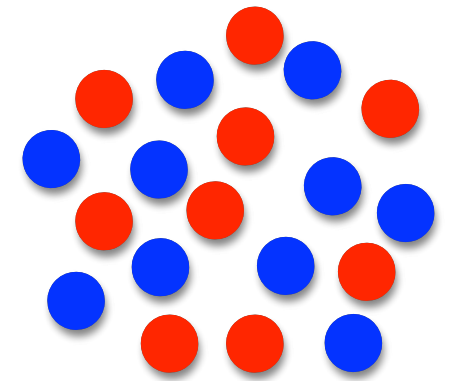
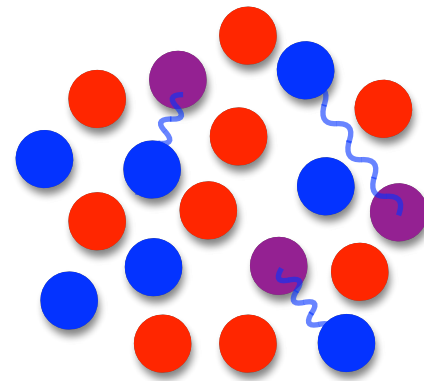
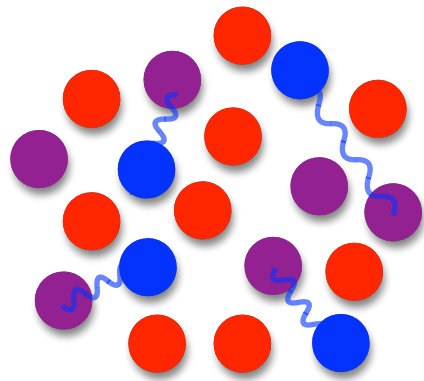
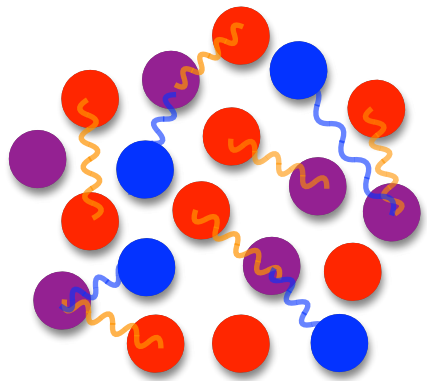
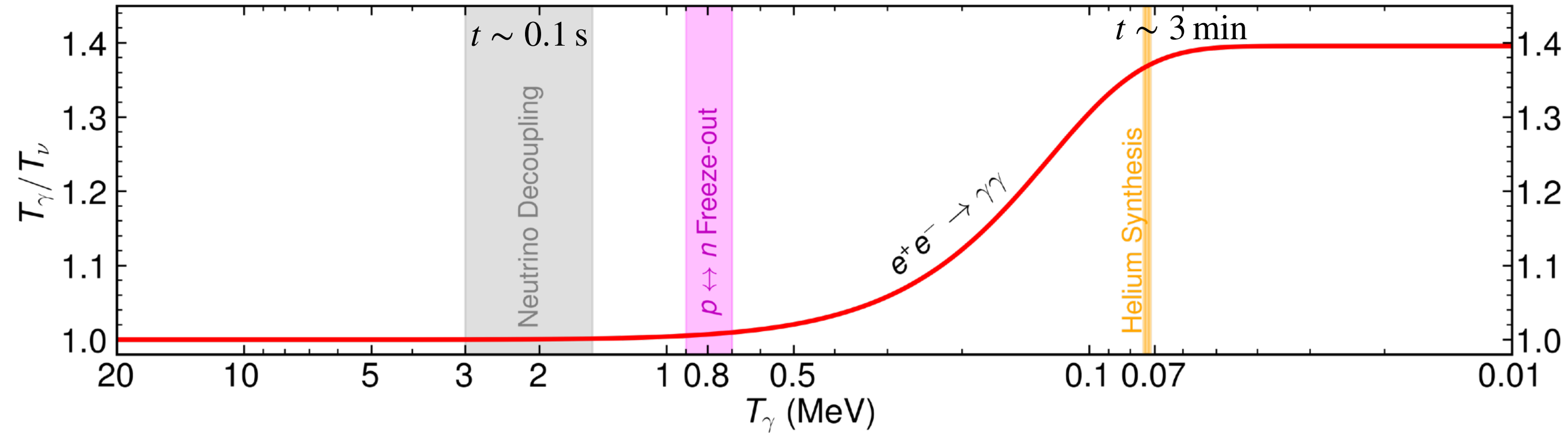
Set Up

Unlike neutrinos, I do like to interact 😊

**Questions, Comments and
Criticism are most
welcome, at any time!!!!**

Neutrino Decoupling

Evolution in the Standard Model



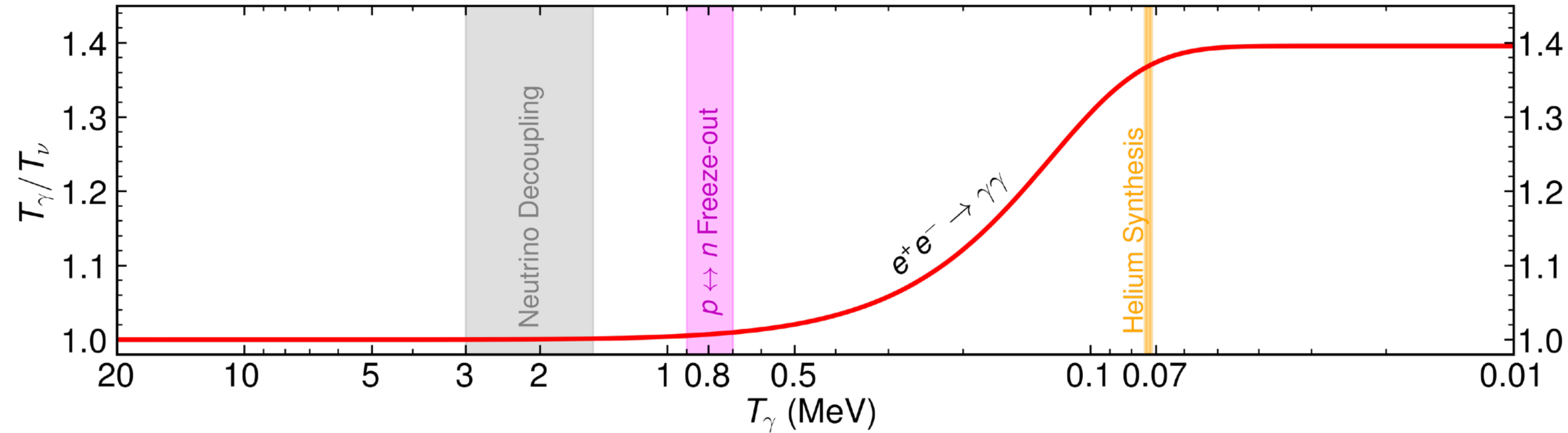
$$e^+e^- \leftrightarrow \bar{\nu}_i\nu_i$$

$$e^\pm\nu_i \leftrightarrow e^\pm\nu_i$$

- **Neutrinos**
- **Electrons**
- **Photons**
- ~ **Z-W (off-shell)**

Cosmic Neutrino Background

Evolution in the Standard Model



- $$N_{\text{eff}} \equiv \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \left(\frac{\rho_{\text{rad}} - \rho_\gamma}{\rho_\gamma} \right)$$

$$N_{\text{eff}} = 3 \left(\frac{1.4T_\nu}{T_\gamma} \right)^4$$

- $$N_{\text{eff}}^{\text{SM}} = 3.044(1)$$

Mangano et al. hep-ph/0506164

de Salas & Pastor 1606.06986

Bennett, Buldgen, Drewes & Wong 1911.04504

Escudero Abenza 2001.04466

Akita & Yamaguchi 2005.07047

Froustey, Pitrou & Volpe 2008.01074

Gariazzo, de Salas, Pastor et al. 2012.02726

Hansen, Shalgar & Tamborra 2012.03948

Why is N_{eff} in the SM not 3?

Recently reviewed by Akita & Yamaguchi, 2210.10307, see also the nice review by Dolgov hep-ph/0202122

Relic Neutrino Decoupling $t \sim 0.1 \text{ s}$ $T_\nu \sim 2 \text{ MeV}$

- 1) Some $e^+e^- \rightarrow \bar{\nu}\nu$ heating because $T_\nu^{\text{dec}} \sim 4 \times m_e$ $\Delta N_{\text{eff}} \simeq +0.03$ Kolb et al. '82
Dolgov et al. '97
- 2) Finite temperature corrections to $\delta m_\gamma^2(T)$ and $\delta m_e^2(T)$ $\Delta N_{\text{eff}} \simeq +0.01$ Heckler '94
Bennet et al. '21
- 3) Neutrino oscillations $\Delta N_{\text{eff}} \simeq +0.0007$ Mangano et al. '05
de Salas & Pastor '16

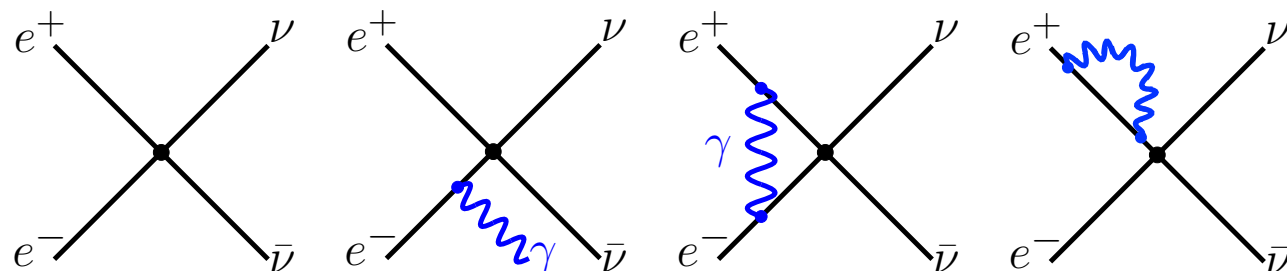
Standard Model prediction as of 2022: $N_{\text{eff}}^{\text{SM}} = 3.0440(2)$ Akita & Yamaguchi 2005.07047
Froustey, Pitrou & Volpe 2008.01074
Bennett, Buldgen, de Salas, Drewes, Gariazzo, Pastor & Wong 2012.02726

4) Finite temperature QED corrections to $e^+e^- \rightarrow \bar{\nu}\nu$ processes

QED corrections are well known to be sizable ($\sim 5\%$) for $\nu e \rightarrow \nu e$ scatterings for solar neutrinos, see e.g. Bahcall, Kamionkowski & Sirlin [astro-ph/9502003]

Estimate of this effect was made in Escudero Abenza 2001.04466 using an interpolation of the NLO rates from Esposito et al. [astro-ph/0301438]

Together with Gianpiero Mangano, Ofelia Pisanti and Mattia Cielo we have actually accurately accounted for the correction to the energy transfer rates which is $\sim -4\%$ at $T = 1 \text{ MeV}$



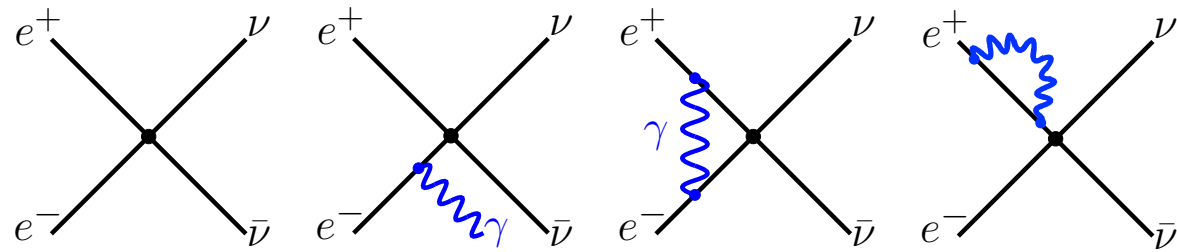
At NLO: $\Delta N_{\text{eff}} \simeq -0.0007$

$$N_{\text{eff}}^{\text{SM}} = 3.0432(2) = 3.043$$

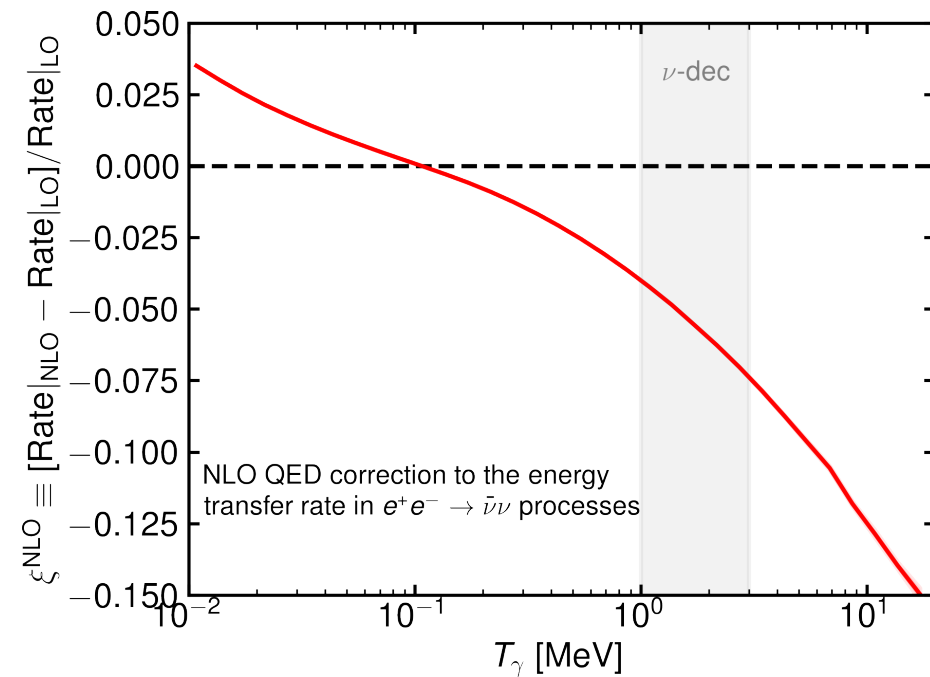
Cielo, Escudero, Mangano & Pisanti 2306.05460

The essence in 1 slide

1) Calculation performed following the real time formalism in thermal field theory



See Esposito, Mangano, Miele, Picard & Pisanti astro-ph/0301438 & astro-ph/0112384



2) Solve for the process of neutrino decoupling:

exact:
$$\frac{df_\nu}{dt} - Hp \frac{\partial f_\nu}{\partial p} = C[f_\nu]$$

approximate but very accurate:
$$\frac{dT_\nu}{dt} = -HT_\nu + \frac{\delta\rho_{\nu e}}{\delta t} + 2 \frac{\delta\rho_{\nu\mu}}{\delta t} \Big/ 3 \frac{\partial\rho_\nu}{\partial T_\nu}$$

Escudero 1812.05605 & 2001.04466 [JCAP]

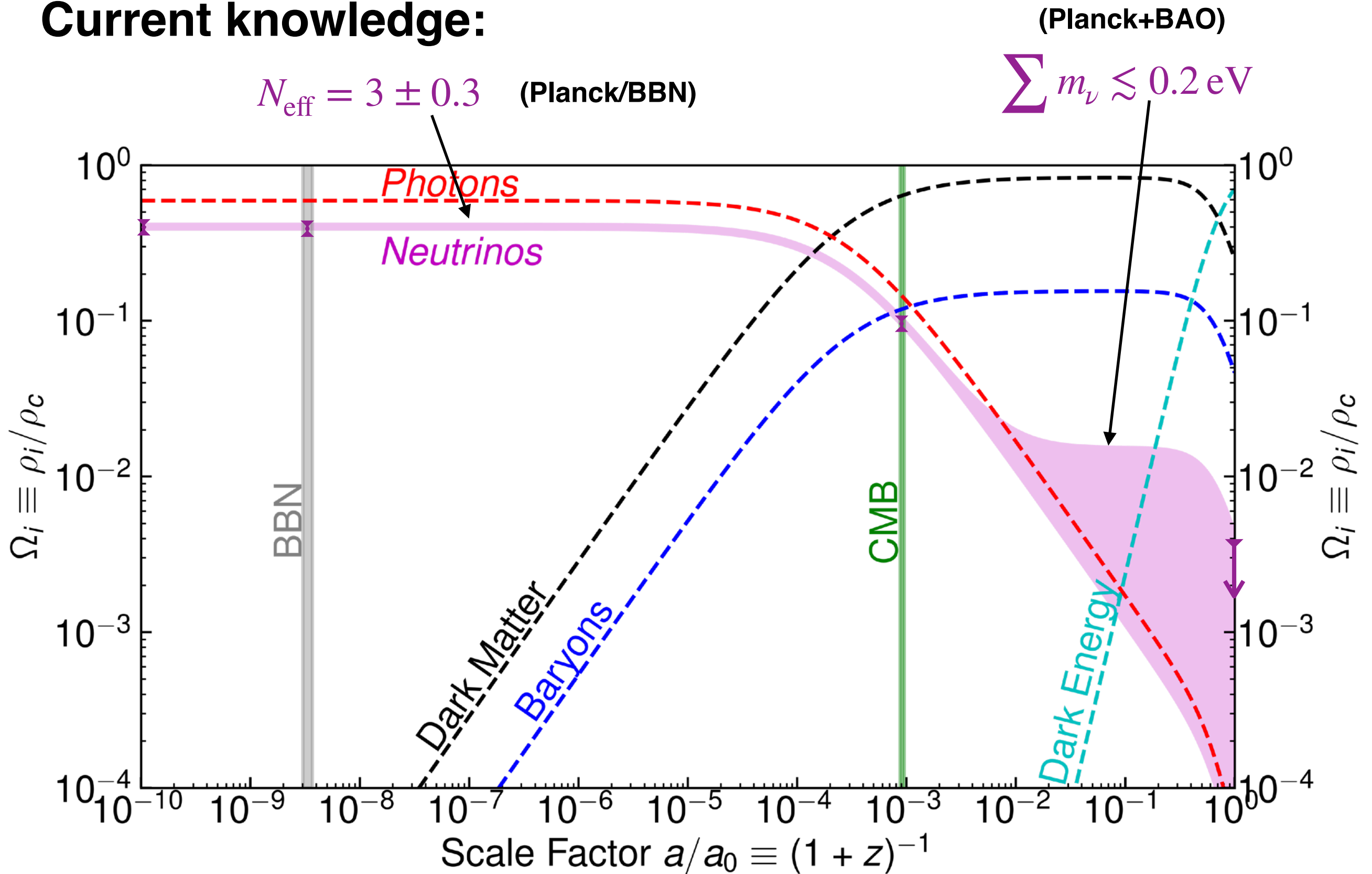
3) Result at NLO: $\Delta N_{\text{eff}} \simeq -0.0007$

$$N_{\text{eff}}^{\text{SM}} = 3.0432(2) = 3.043$$

Cielo, Escudero, Mangano & Pisanti 2306.05460

Global Perspective




Current knowledge:



Evidence for Cosmic Neutrinos

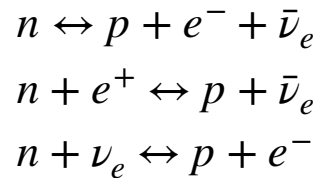
Big Bang Nucleosynthesis

Current measurements are consistent with the SM picture

 H ~ 75%  ${}^4\text{He}$ ~ 25%  D ~ 0.005%

This implies that neutrinos should have been present:

1) It is impossible to have successful BBN without neutrinos. They participate in $p \leftrightarrow n$ conversions up to $T \gtrsim 0.7 \text{ MeV}$



2) Neutrinos contribute to the expansion rate $H \propto \sqrt{\rho}$

By comparing predictions against observations, we know:

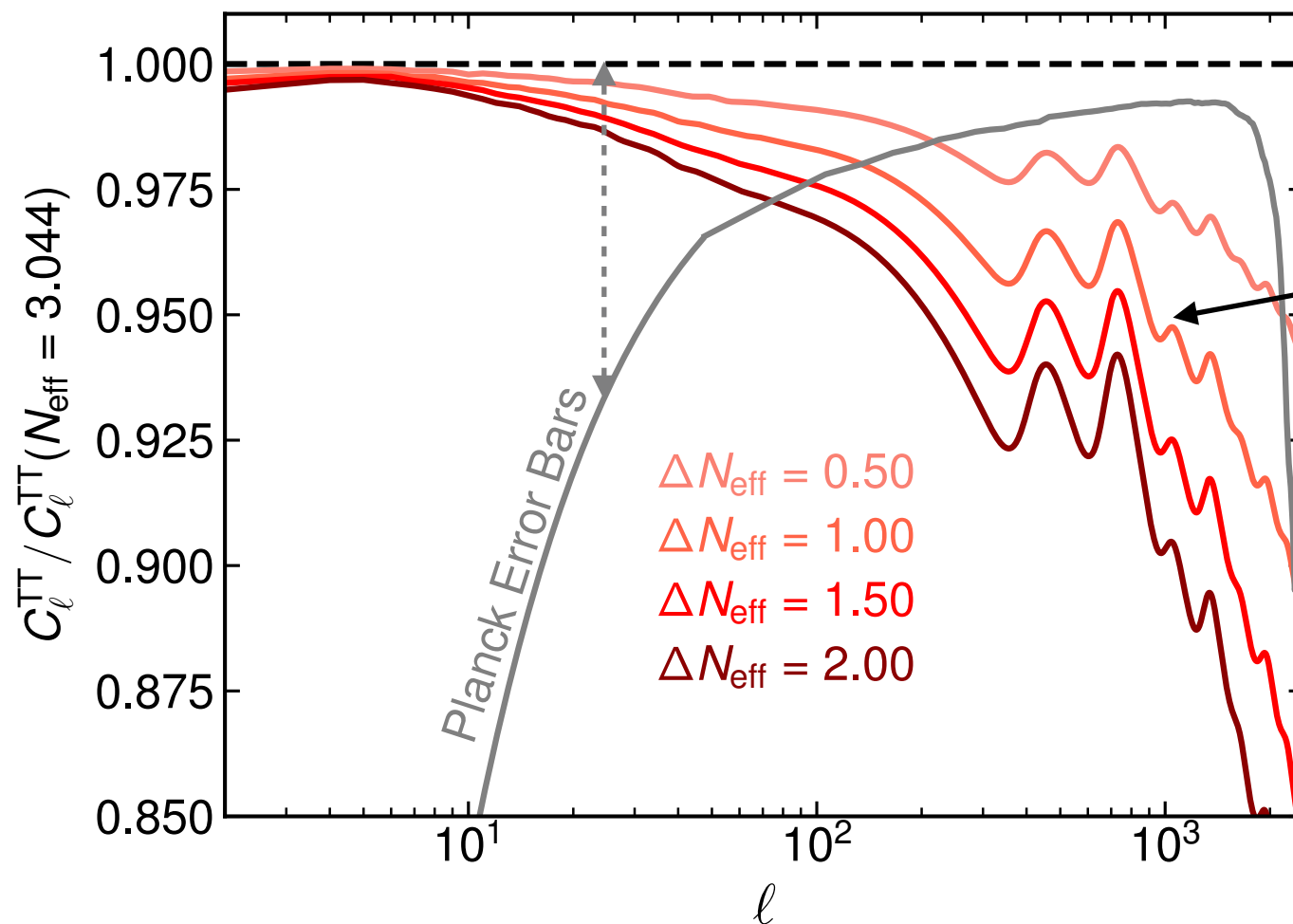
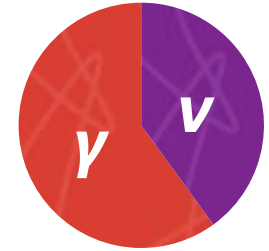
$$N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$$

see e.g., Pisanti et al. 2011.11537
and Yeh et al. 2207.13133

Evidence for Cosmic Neutrinos

Cosmic Microwave Background Why?

Ultra-relativistic neutrinos represent a large fraction of the energy density of the Universe, $H \propto \sqrt{\rho}$



N_{eff} is constrained by the high- ℓ multipoles, i.e. Silk damping

$$N_{\text{eff}}^{\text{CMB+BAO}} = 2.99 \pm 0.17$$

Planck 2018 1807.06209

Evidence for Cosmic Neutrinos

- **Current constraints**

BBN

$$N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$$

Pisanti et al. 2011.11537
Yeh et al. 2207.13133

Planck+BAO

$$N_{\text{eff}}^{\text{CMB}} = 2.99 \pm 0.17$$

Planck 2018, 1807.06209

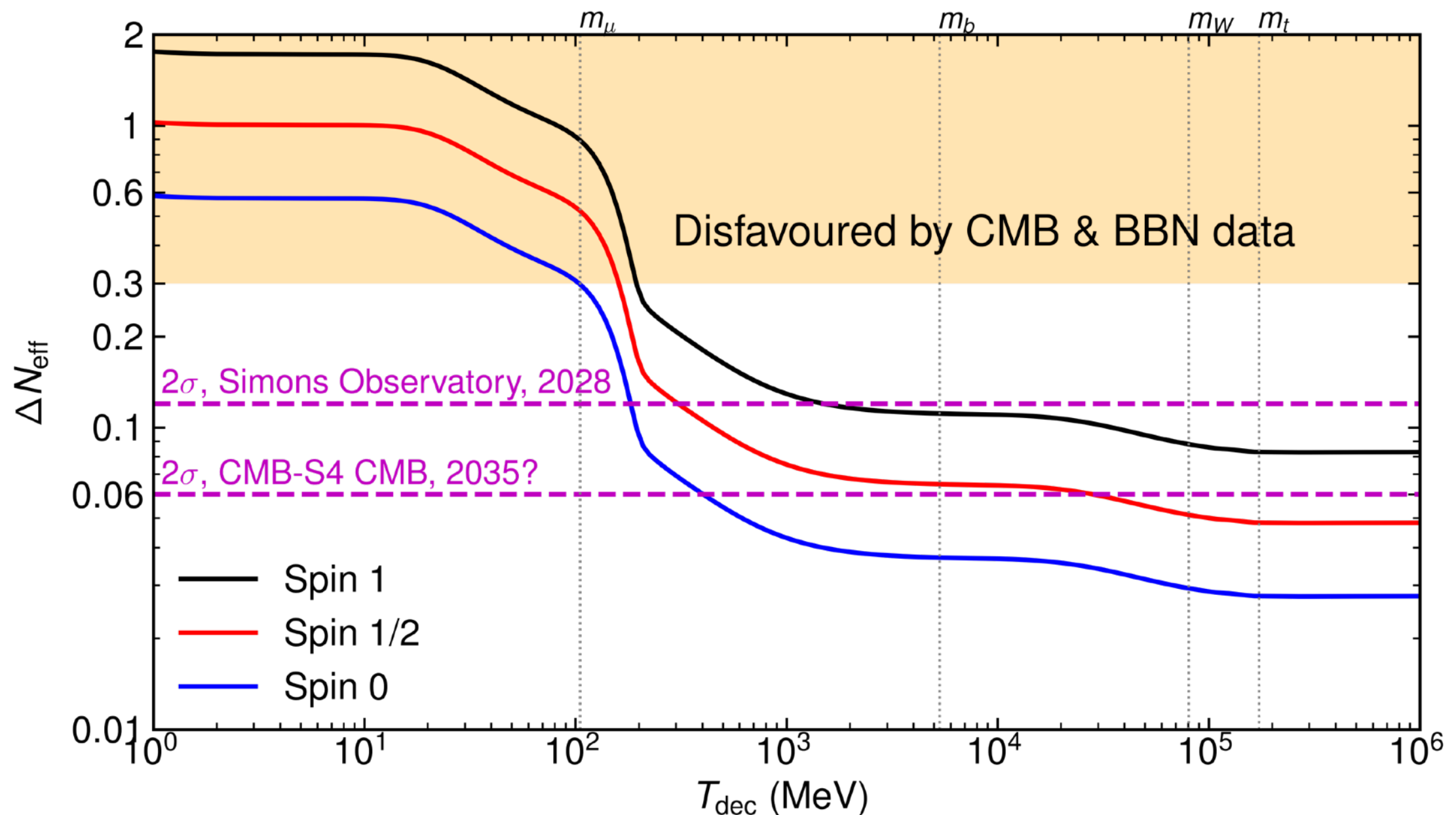
- **Standard Model prediction:** $N_{\text{eff}}^{\text{SM}} = 3.043$
- **Data is in excellent agreement with the Standard Model prediction**
- **This provides strong (albeit indirect) evidence for the Cosmic Neutrino Background**

Implications:

- 1) **Stringent constraint on many BSM scenarios**
- 2) **We can use cosmological data to test neutrino properties**

Constraints on N_{eff}

- **Sterile Neutrino** $m_N \sim \text{eV}$ $\Delta N_{\text{eff}} = 1$ (e.g. Gariazzo, de Salas & Pastor 1905.11290)
- **Goldstone Bosons** Weinberg 1305.1971
- **Other sterile long-lived particles** Axion, gravitino, axino, hidden sector particles ...



Stringent constraint on BSM physics

Constraints are relevant in many other BSM settings:

- **WIMPs** $m_{\text{WIMP}} > (4 - 10) \text{ MeV}$ Sabti et. al. 1910.01649
Boehm et. al. 1303.6270
- **GeV-Sterile Neutrinos** $\tau_N \lesssim 0.05 \text{ s}$ Sabti et. al. 2006.07387
Dolgov et. al. hep-ph/0008138
- **Vector Bosons** $g \lesssim 10^{-10} \quad m \lesssim 10 \text{ MeV}$ Escudero et. al. 1901.02010
Kamada & Yu 1504.00711
- **Axions** Raffelt et. al. 1011.3694
Blum et.al. 1401.6460
- **Low Reheating** $T_{\text{RH}} > (2 - 5) \text{ MeV}$ de Salas et. al. 1511.00672
Hasegawa et. al. 1908.10189
- **Variations of GN** $G_{\text{BBN}}/G_0 = 0.98 \pm 0.03$ Alvey et. al. 1910.10730
Copi et.al. astro-ph/0311334
- **PBHs** $6 \times 10^8 \text{ g} < M_{\text{PBH}} < 2 \times 10^{13} \text{ g}$ Carr et. al. 0912.5297
Keith et.al. 2006.03608

Check out a review on non-standard expansion histories:

[2006.16182](#) Vaskonen et al. (Escudero & Poulin)

Evidence for Cosmic Neutrinos

- **Current constraints**

BBN

$$N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$$

Pisanti et al. 2011.11537
Yeh et al. 2207.13133

Planck+BAO

$$N_{\text{eff}}^{\text{CMB}} = 2.99 \pm 0.17$$

Planck 2018, 1807.06209

- **Standard Model prediction:** $N_{\text{eff}}^{\text{SM}} = 3.043$
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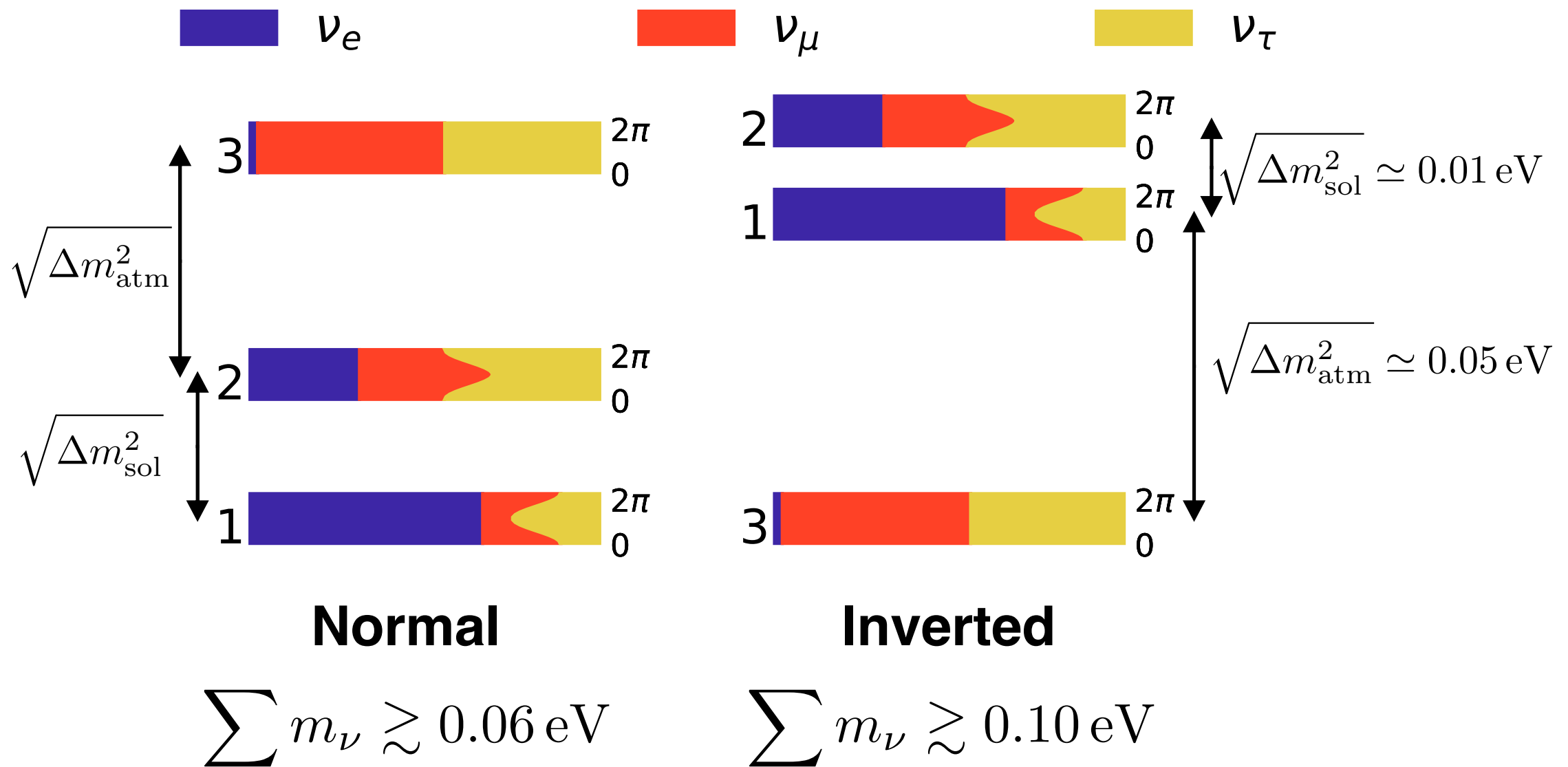
Implications:

1) **Stringent constraint on many BSM scenarios**

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Neutrino Properties

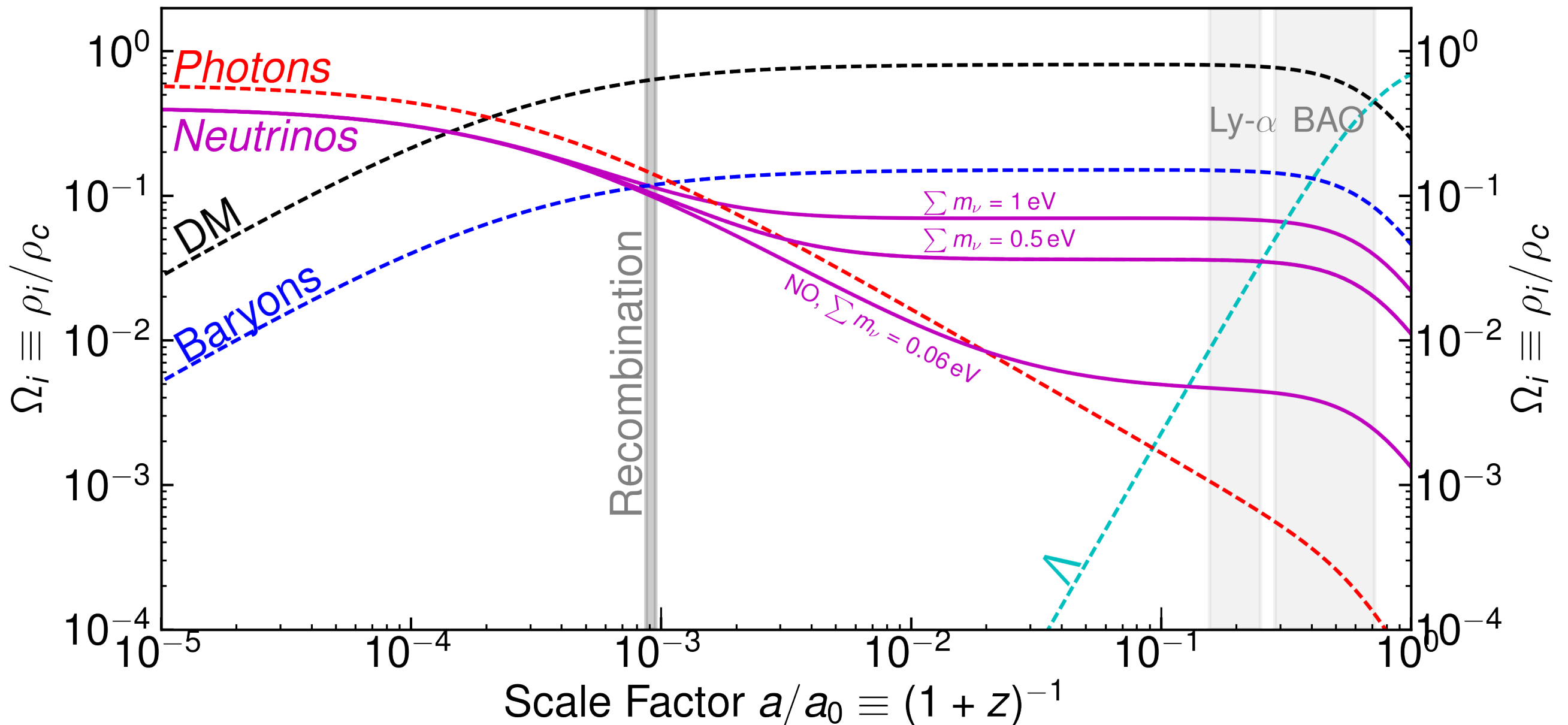
Figure from de Salas et al. 1806.11051



- Mass differences and mixings measured with high precision
- What is δ_{CP} and what is the mass ordering? [Neutrino Oscillations](#)
- Are Neutrinos Dirac or Majorana Particles? [\$0\nu 2\beta\$ Experiments](#)
- What is the neutrino mass scale? i.e. $\sum m_\nu$? i.e. m_{lightest} ? [Cosmology](#)

Neutrino Masses in Cosmology

- 1) Massive neutrinos modify the expansion history



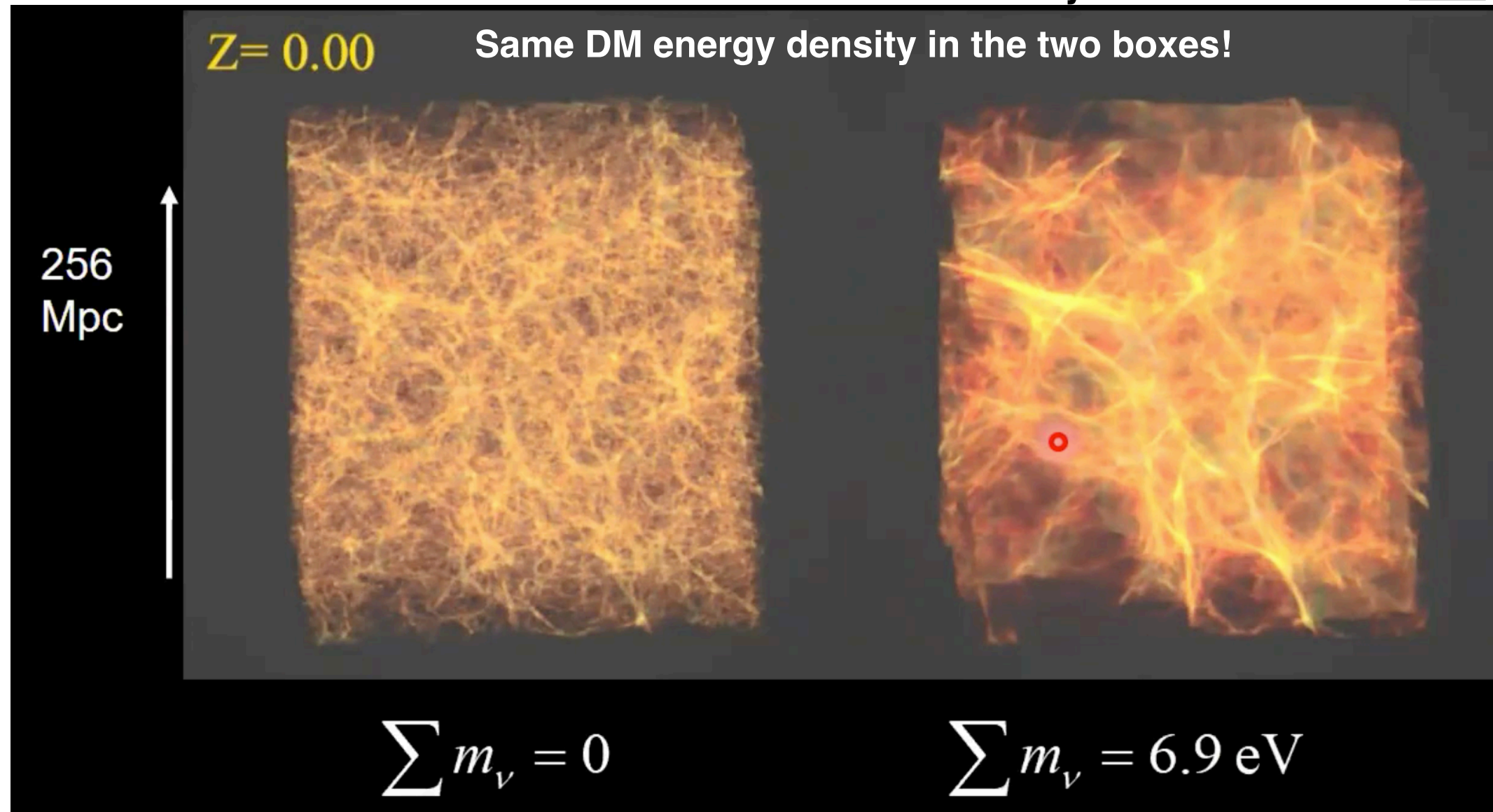
Non-Rel: $z_{\nu}^{\text{non-rel}} \simeq 200 \frac{m_{\nu}}{0.1 \text{ eV}}$

Hot DM: $\Omega_{\nu} h^2 = \sum m_{\nu} / (93.14 \text{ eV})$

Neutrino Masses in Cosmology

- 2) Massive neutrinos suppress the growth of structure

Taken from a talk by Steen Hannestad [Link](#).



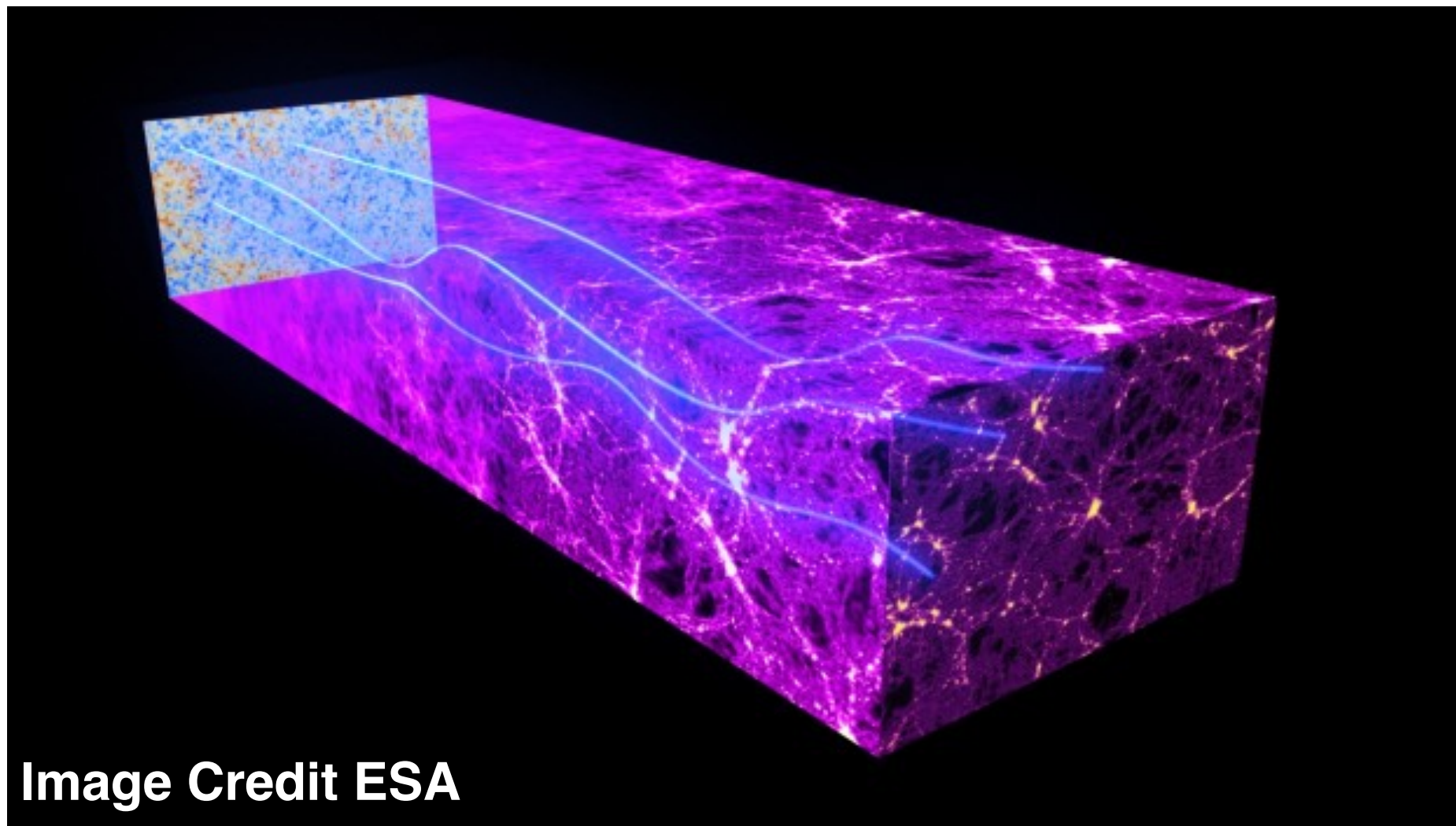
This happens because neutrinos travel very fast and therefore cannot fall in gravitational potentials. The effect of this smoothing is proportional to Ω_ν

Neutrino Masses in Cosmology

Cosmic Microwave Background Anisotropies

Neutrinos of $m_\nu < 0.5 \text{ eV}$ become non-relativistic after recombination. That means that their effect on the anisotropies is somewhat small!

The most relevant impact is through the effect of gravitational lensing:

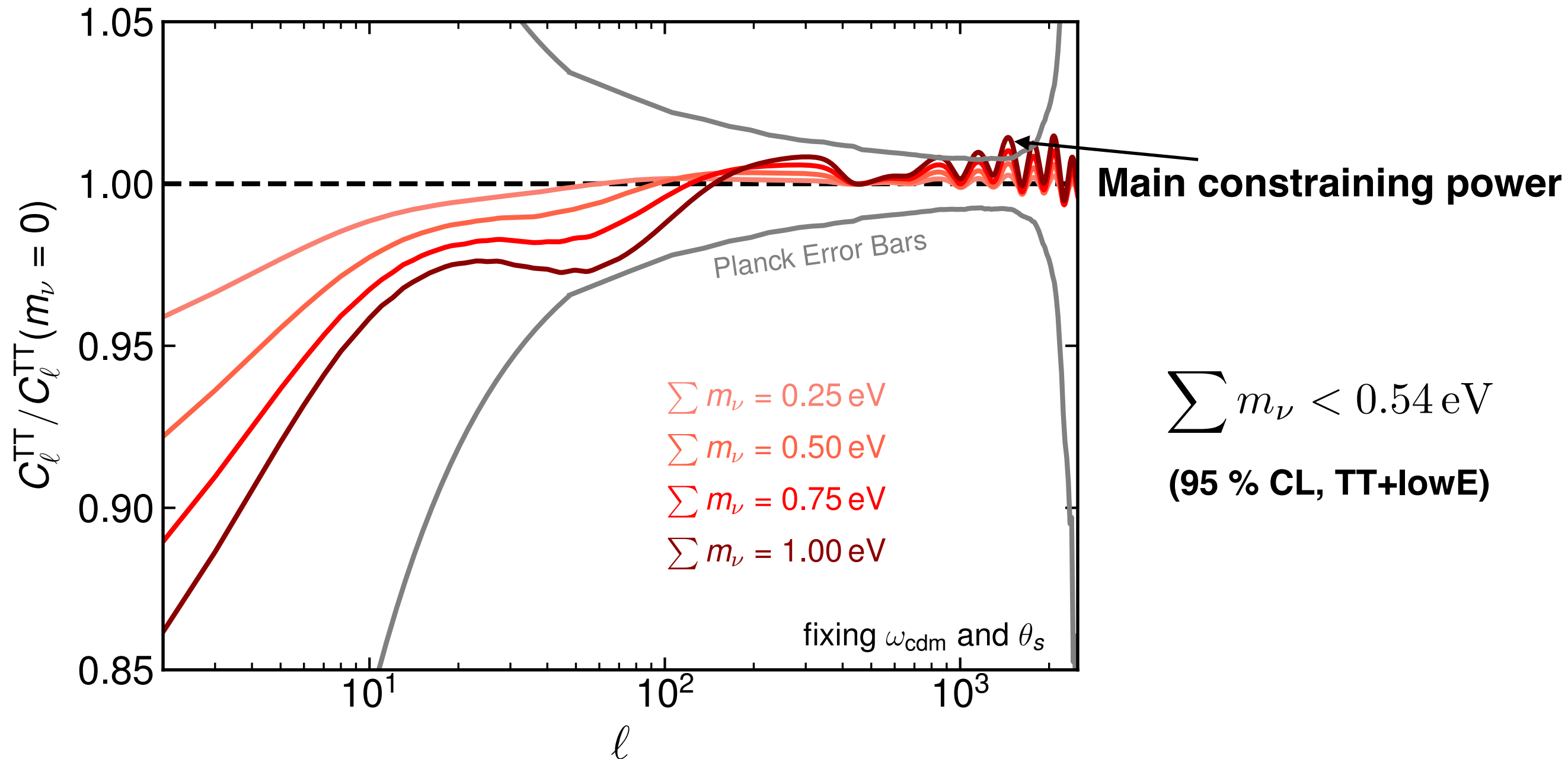


The larger the neutrino mass the less is the CMB light lensed!

Neutrino Masses in Cosmology

Cosmic Microwave Background Anisotropies

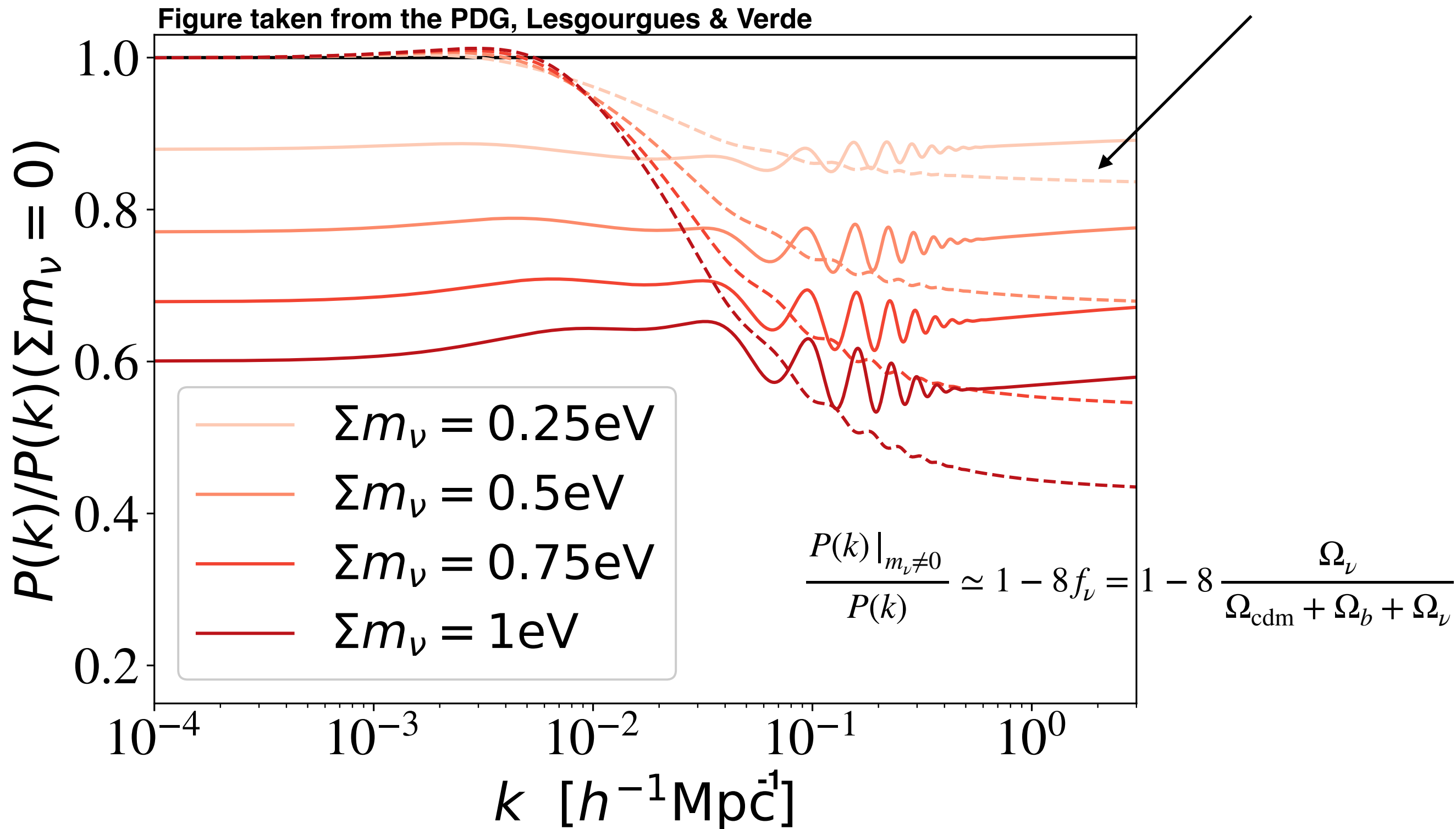
The effect of neutrino masses in the CMB:



Neutrino Masses in Cosmology

Galaxy Surveys

Suppression from $\Omega_\nu h^2$



Neutrino Masses from Cosmology

Planck 2018 for Λ CDM (1807.06209)

$$\sum m_\nu < 0.54 \text{ eV} \quad (95 \% \text{ CL, TT+lowE})$$

$$\sum m_\nu < 0.26 \text{ eV} \quad (95 \% \text{ CL, TTTEEE+lowE})$$

$$\sum m_\nu < 0.24 \text{ eV} \quad (95 \% \text{ CL, TTTEEE+lowE+lensing})$$

$$\sum m_\nu < 0.12 \text{ eV} \quad (95 \% \text{ CL, TTTEEE+lowE+lensing+BAO})$$

To be compared to the KATRIN bound: $\sum m_\nu < 2.4 \text{ eV}$

Very robust bounds from linear Cosmology $\Delta T/T \sim 10^{-5}$

What about other non-linear cosmological data?

Importantly, all cosmological bounds are cosmological model dependent

What is the dependence upon the assumed Cosmological Model?

Neutrino Masses from Cosmology

Data beyond Planck and BAO within Λ CDM

$\sum m_\nu < 0.26 \text{ eV}$	Planck	Planck 1807.06209
$\sum m_\nu < 0.12 \text{ eV}$	Planck+BAO	Planck 1807.06209
$\sum m_\nu < 0.86 \text{ eV}$	BOSS P(k)	Ivanov et al. 1909.05277
$\sum m_\nu < 0.16 \text{ eV}$	Planck+BOSS P(k)	Ivanov et al. 1912.08208
$\sum m_\nu < 0.58 \text{ eV}$	Lyman-α+H_0prior	Palanque-Delabrouille et al. 1911.09073
$\sum m_\nu < 0.10 \text{ eV}$	Planck+Lyman-α	Choudhury & Hannestad 1907.12598
$\sum m_\nu < 0.08 \text{ eV}$	Planck+BAO+H_0	di Valentino, Gariazzo & Mena 2106.15267
$\sum m_\nu < 0.09 \text{ eV}$	Planck+BAO+SN+RSD	

- **Planck is driving current cosmological constraints**
- **Non-linear or mildly non-linear data sets break degeneracies in the fit**
- **The larger H_0 is, the stronger the constraint on $\sum m_\nu$ is** (However, this comes from combining two data sets in strong tension!)

Neutrino Masses from Cosmology

Cosmological Model Dependence

Planck+BAO and 3 degenerate neutrinos

$$\sum m_\nu < 0.12 \text{ eV}$$

Standard Case

Planck 1807.06209

Λ CDM+m_ν

$$\sum m_\nu < 0.25 \text{ eV}$$

Dark Energy dynamics

Choudhury & Hannestad 19'

CDM+m_ν+ω_a+ω

$$\sum m_\nu < 0.15 \text{ eV}$$

Varying Curvature

Choudhury & Hannestad 19'

Λ CDM+m_ν+Ω_k

$$\sum m_\nu < 0.13 \text{ eV}$$

Varying N_{eff}

Planck 1807.06209

Λ CDM+m_ν+N_{eff}

$$\sum m_\nu < 0.17 \text{ eV}$$

Varying N_{eff}+ω+a_s+m_ν

di Valentino et al. 1908.01391

CDM+m_ν+N_{eff}+ω+a_s+m_ν

- **Constraints are robust upon standard modifications of Λ CDM**

Neutrino Masses from Cosmology

Cosmological Model Dependence

Non-standard Neutrino Cosmologies:

Invisible Neutrino Decay

$$\nu_i \rightarrow \nu_j \phi$$

$$\sum m_\nu \lesssim 0.2 \text{ eV}$$

Oldengott, Wong et al. 2203.09075 & 2011.01502
Escudero & Fairbairn 1907.05425

$$\nu_i \rightarrow \nu_4 \phi$$

$$\sum m_\nu \lesssim 0.42 \text{ eV}$$

Abellán, Poulin et al. 1909.05275, 2112.13862
Escudero, López-Pavón, Rius & Sandner 2007.04994

Time Dependent Neutrino Masses

Late phase transition

$$\sum m_\nu < 1.4 \text{ eV}$$

Dvali & Funcke 1602.03191
Lorenz et al. 1811.01991 & 2102.13618

Ultralight scalar field screening

$$\sum m_\nu < 3 \text{ eV}$$

Esteban & Salvadó 2101.05804
Wetterich et al. 1009.2461

Non-standard Neutrino Populations

$$T_\nu < T_\nu^{\text{SM}} + \text{DR}$$

$$\sum m_\nu < 3 \text{ eV}$$

Farzan & Hannestad 1510.02201
Escudero, Schwetz & Terol-Calvo 2211.01729

$$\langle p_\nu \rangle > 3.15 T_\nu^{\text{SM}}$$

$$\sum m_\nu < 3 \text{ eV}$$

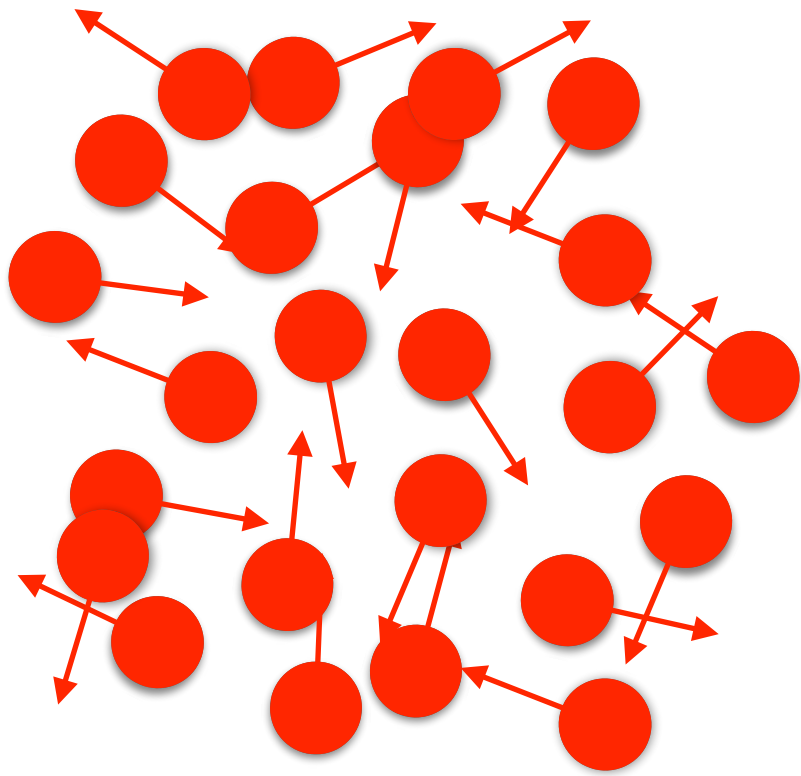
Oldengott et al. 1901.04352
Alvey, Escudero & Sabti 2111.12726

- **Bounds can be significantly relaxed in some extensions of Λ CDM. They require modifications to the neutrino sector.**

But Why? and How?

The Cosmic Neutrino Background

SM

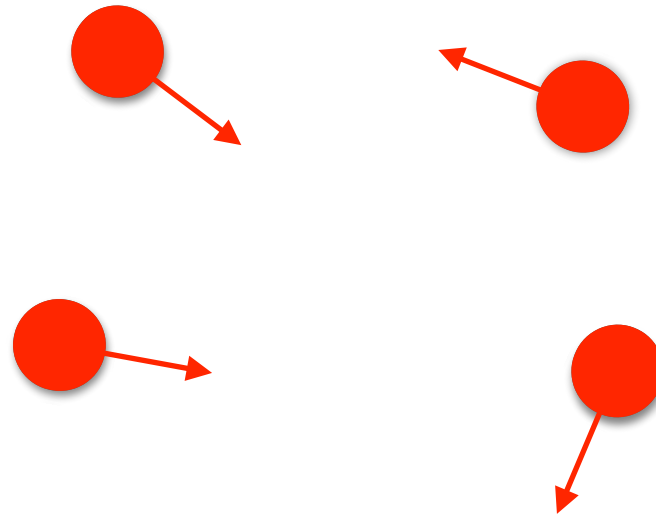


$$n_\nu^{\text{SM}} \sim 300 \text{ cm}^{-3}$$

$$\sum m_\nu < 0.12 \text{ eV}$$

$$n_\nu \propto T_\nu^3$$

$$\text{😞 } N_{\text{eff}} \propto T_\nu^4 \ll 3$$

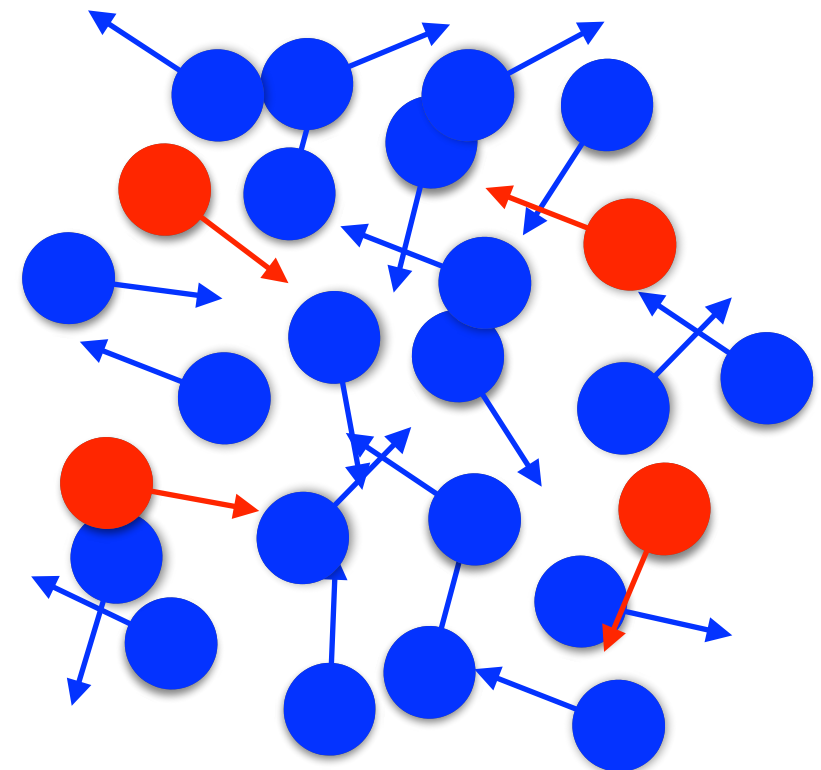


$$n_\nu^{\text{SM}} \sim 30 \text{ cm}^{-3}$$

$$\text{😄 } \sum m_\nu < 1.2 \text{ eV}$$

solution:

add Dark Radiation



$$\text{😄 } N_{\text{eff}} \simeq 3$$

$$\text{😄 } \sum m_\nu < 1.2 \text{ eV}$$

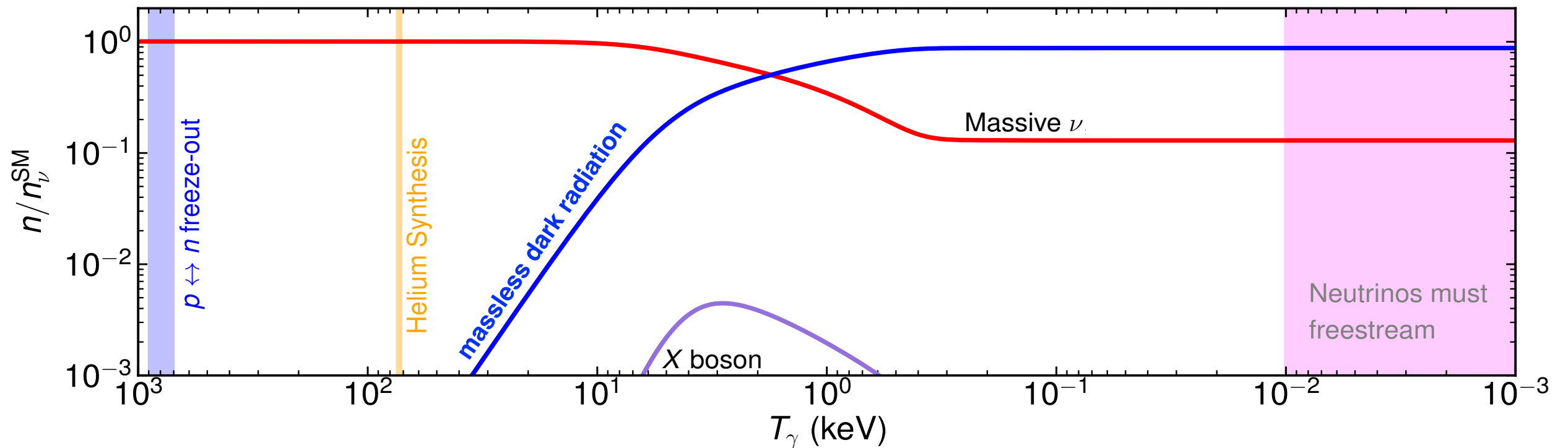
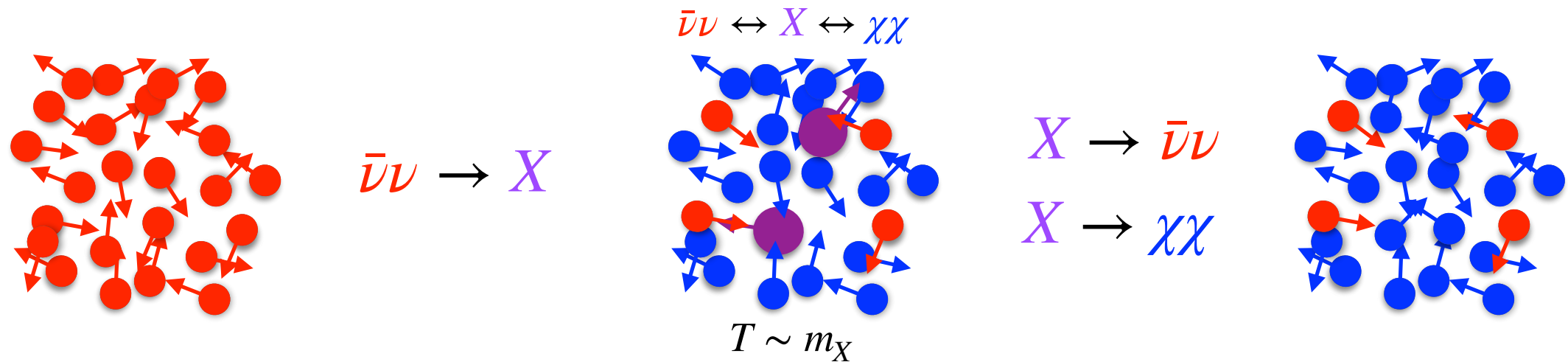
CMB observations can only constrain: $\rho_\nu^0 \equiv \sum m_\nu n_\nu$

see Alvey, M.E.A. & Sabti [2111.12726](#) and also Oldengott et al. [1901.04352](#) and Renk et al. [2009.03286](#)

The Picture

Farzan & Hannestad 1510.02201

Escudero, Schwetz & Terol-Calvo 2211.01729



$$m_{Z'} \sim 1 - 100 \text{ keV}$$

$$g_{\nu\nu} \gtrsim 10^{-10}$$

Summary of main features

Escudero, Schwetz & Terol-Calvo 2211.01729

● Requirements:

- 1) Features a neutrino mass mechanism (type-I seesaw)
- 2) Have a large number of massless sterile states coupled to neutrinos

$$\sum m_\nu < 0.12 \text{ eV} \left[1 + N_\chi/3 \right] \quad N_\chi \sim 10 - 20$$

- 3) These states are nevertheless coupled enough so that they can thermalize in the early Universe between BBN and recombination
- 4) A new interacting boson at the keV scale

● The essence of the models:

- Add a $U(1)_X$ symmetry with a scalar field and a singlet left-handed state S_L

$$\mathcal{L} = y\Phi \bar{N}_R S_L \quad M_\nu|^{7 \times 7} = \begin{pmatrix} 0 & m_D & 0 \\ m_D^t & M_R & y_\alpha v_\Phi \\ 0 & (y_\alpha v_\Phi)^t & 0 \end{pmatrix}$$

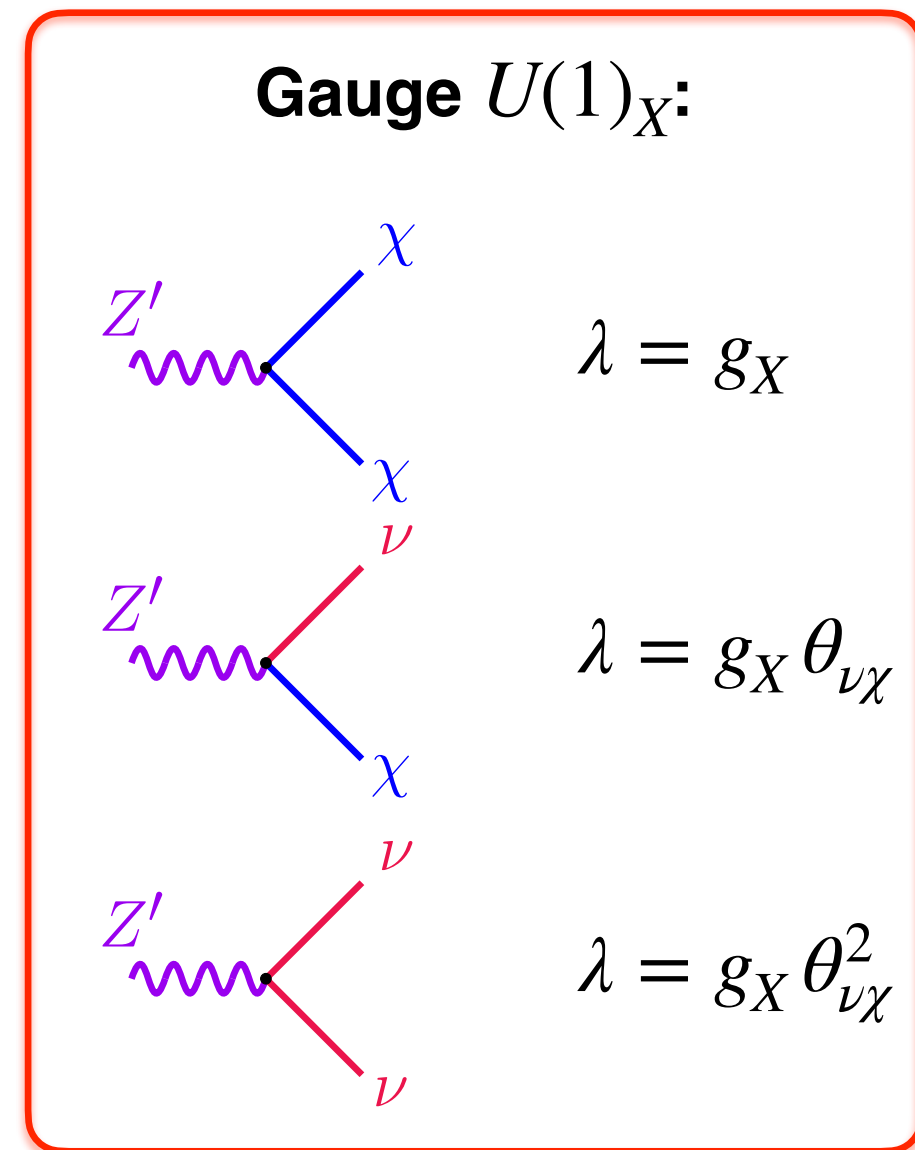
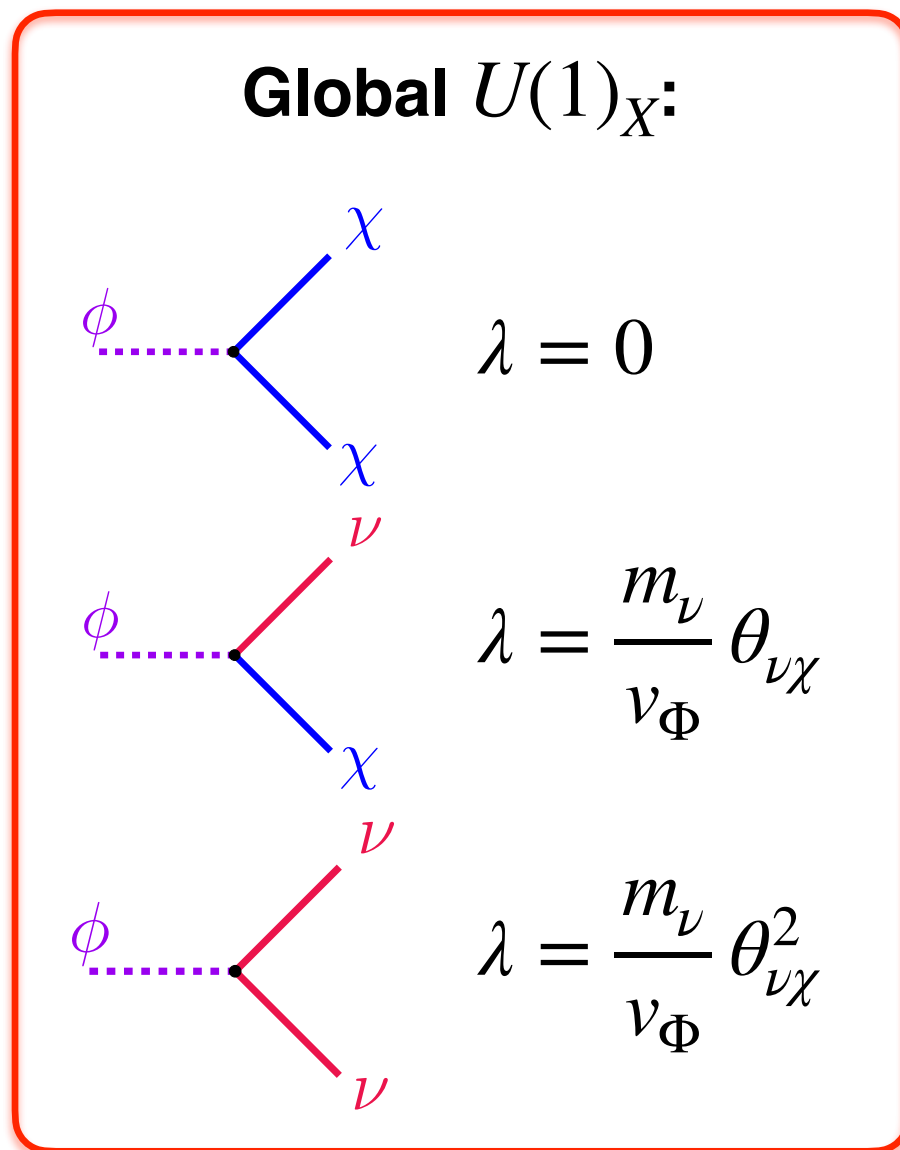
Provided $y_\alpha v_\Phi \ll m_D$ ● Seesaw mechanism at play $m_\nu \simeq m_D^2 / M_R$

● Right ν_4 properties: $m_\chi \simeq 0 \quad U_{\alpha 4} \sim \frac{y_\alpha v_\Phi}{m_D} \ll 1$

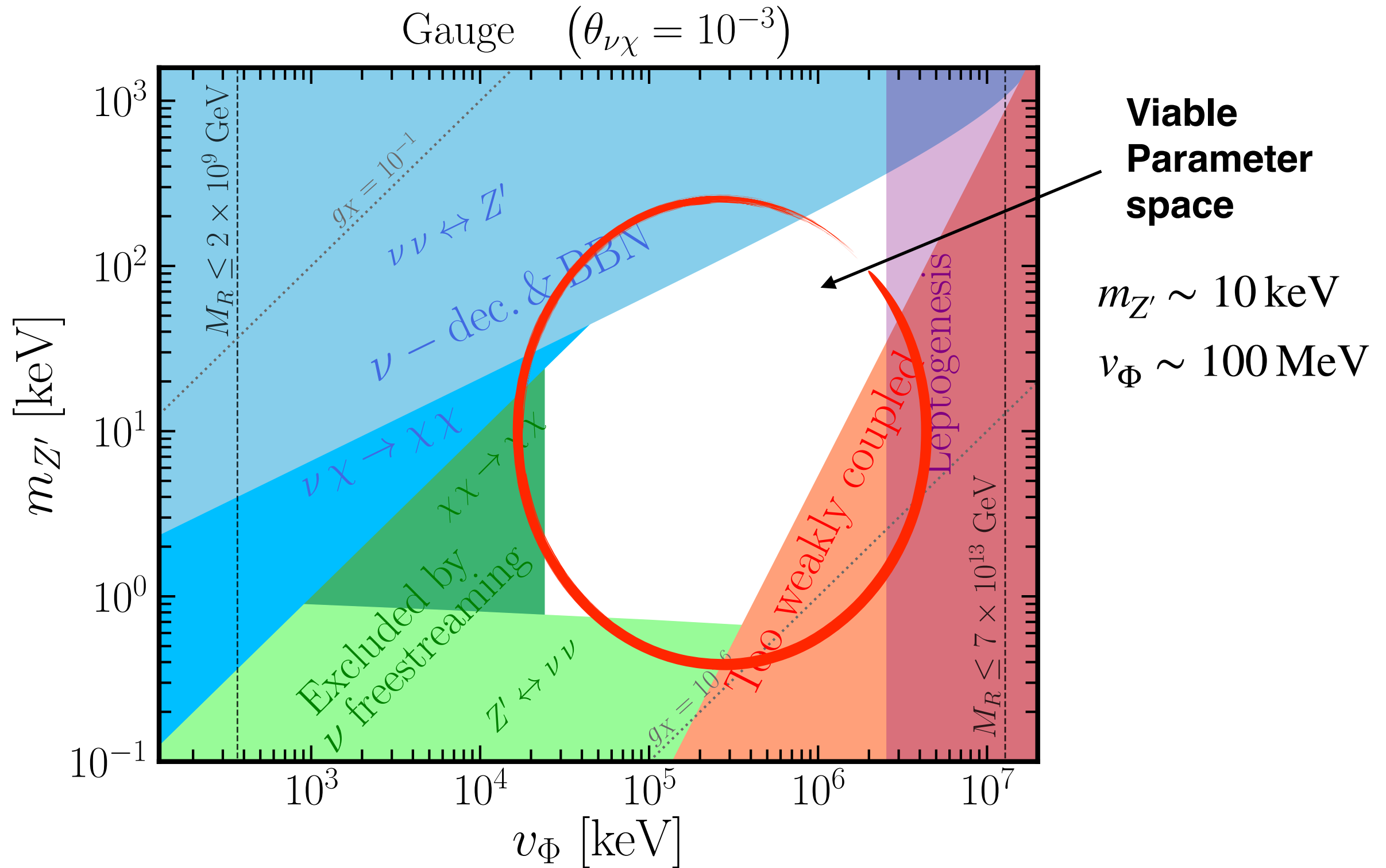
Trivial to generalize to the case of $N_\chi \sim 10 - 20$. Additional Z_2 needed for Gauge case

Complete UV models

- **Parameters:**
 - $\theta_{\nu\chi}$ Active-sterile mixing
 - v_Φ Scale at which $U(1)_X$ is spontaneously broken
 - m_X Mass of mediator
- **Two cases:**



Parameter Space



Neutrino Masses from Cosmology

Non-standard Neutrino Cosmologies:

Invisible Neutrino Decay

$$\nu_i \rightarrow \nu_j \phi$$
$$\sum m_\nu < 0.2 \text{ eV}$$

Oldengott, Wong et al. 2203.09075 & 2011.01502
Escudero & Fairbairn 1907.05425
Archidiacono & Hannestad 1311.3873

Time Dependent Neutrino Masses

Late phase transition

$$\sum m_\nu < 1.4 \text{ eV}$$

Dvali & Funcke 1602.03191
Lorenz et al. 1811.01991 & 2102.13618

Non-standard Neutrino Populations

$$T_\nu < T_\nu^{\text{SM}}$$

$$\sum m_\nu < 3 \text{ eV}$$

Farzan & Hannestad 1510.02201
Escudero, Schwetz & Terol-Calvo 2211.01729

$$\nu_i \rightarrow \nu_4 \phi$$

at least: $\sum m_\nu \lesssim 0.42 \text{ eV}$

Abellán, Poulin et al. 1909.05275, 2112.13862
Escudero, López-Pavón, Rius & Sandner 2007.04994

Ultralight scalar field screening

$$\sum m_\nu < 3 \text{ eV}$$

Esteban & Salvadó 2101.05804
Wetterich et al. 1009.2461

$$\langle p_\nu \rangle > 3.15 T_\nu^{\text{SM}}$$

$$\sum m_\nu < 3 \text{ eV}$$

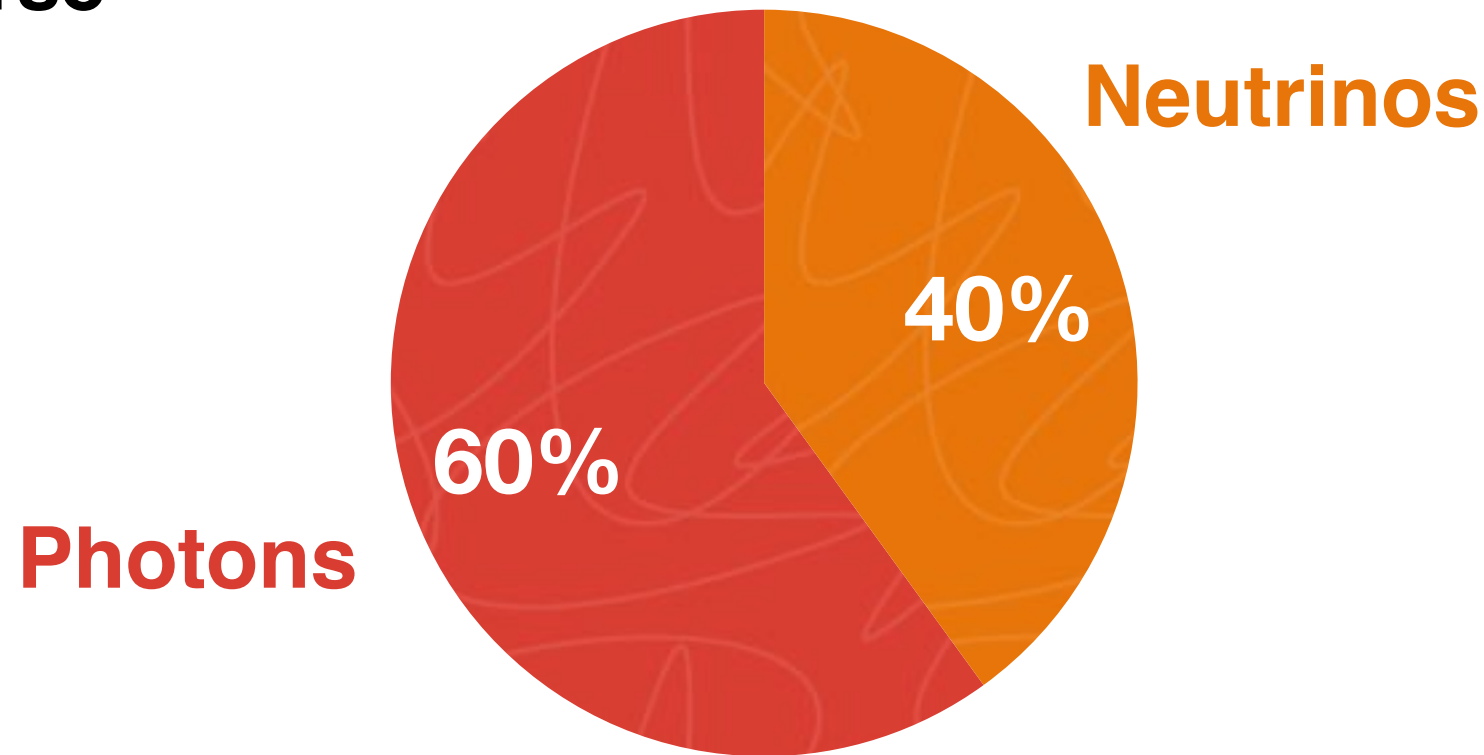
Oldengott et al. 1901.04352
Alvey, Escudero & Sabti 2111.12726

Take Away Messages:

- Cosmology can only constrain $\Omega_\nu(z)$ and not directly m_ν
- Of course, in Λ CDM there is a direct link between $\Omega_\nu(z)$ and m_ν
- All these models reduce $\Omega_\nu(z)$ with respect to the one in Λ CDM and are in excellent agreement with all known cosmological data
- Importantly, they entail non-standard neutrino properties

Neutrino Interactions

- Neutrinos represent a large component of the energy density of the Universe



- Neutrinos have very special cosmological perturbations
 - 1) They are ultrarelativistic until $z \sim 200 m_\nu / 0.1 \text{ eV}$
 - 2) In the SM: since $t_U \sim 0.1 \text{ s}$ ($T \sim 2 \text{ MeV}$), they are free streaming i.e. do not interact with anything

These together actually mean that CMB observations can probe potential neutrino interactions!

Why?

First discussed by Bashinsky & Seljak in [astro-ph/0310198] and applied by Chacko, Hall, Okui & Oliver [hep-ph/0312267]

- The key is in Einstein's equations

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad \text{Background expansion: } N_{\text{eff}}$$

$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu} \quad \text{Perturbations: can tell about interactions}$$

Neutrino anisotropic stress

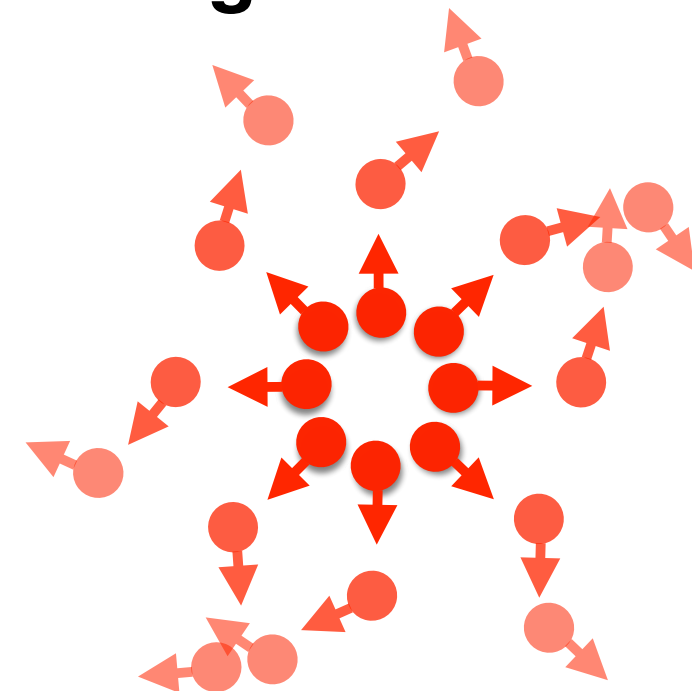
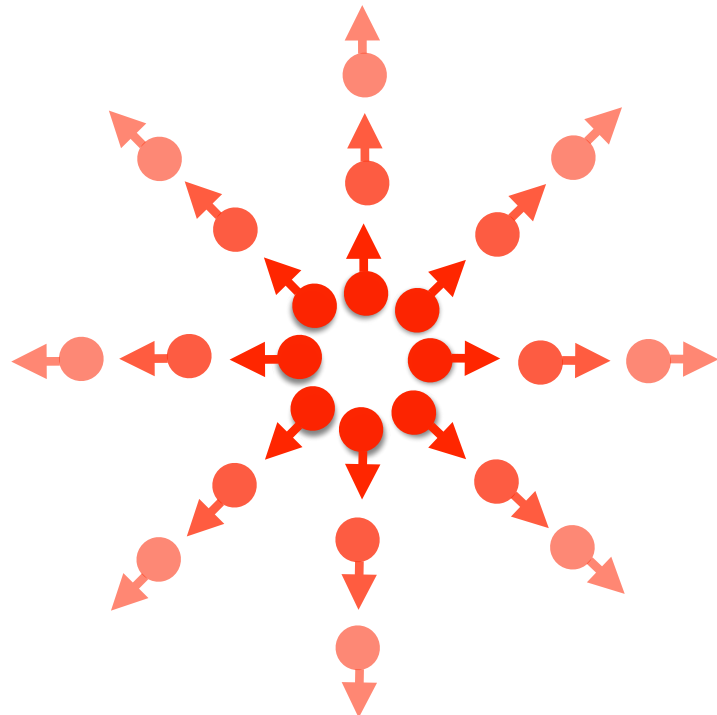
Metric

CMB spectra

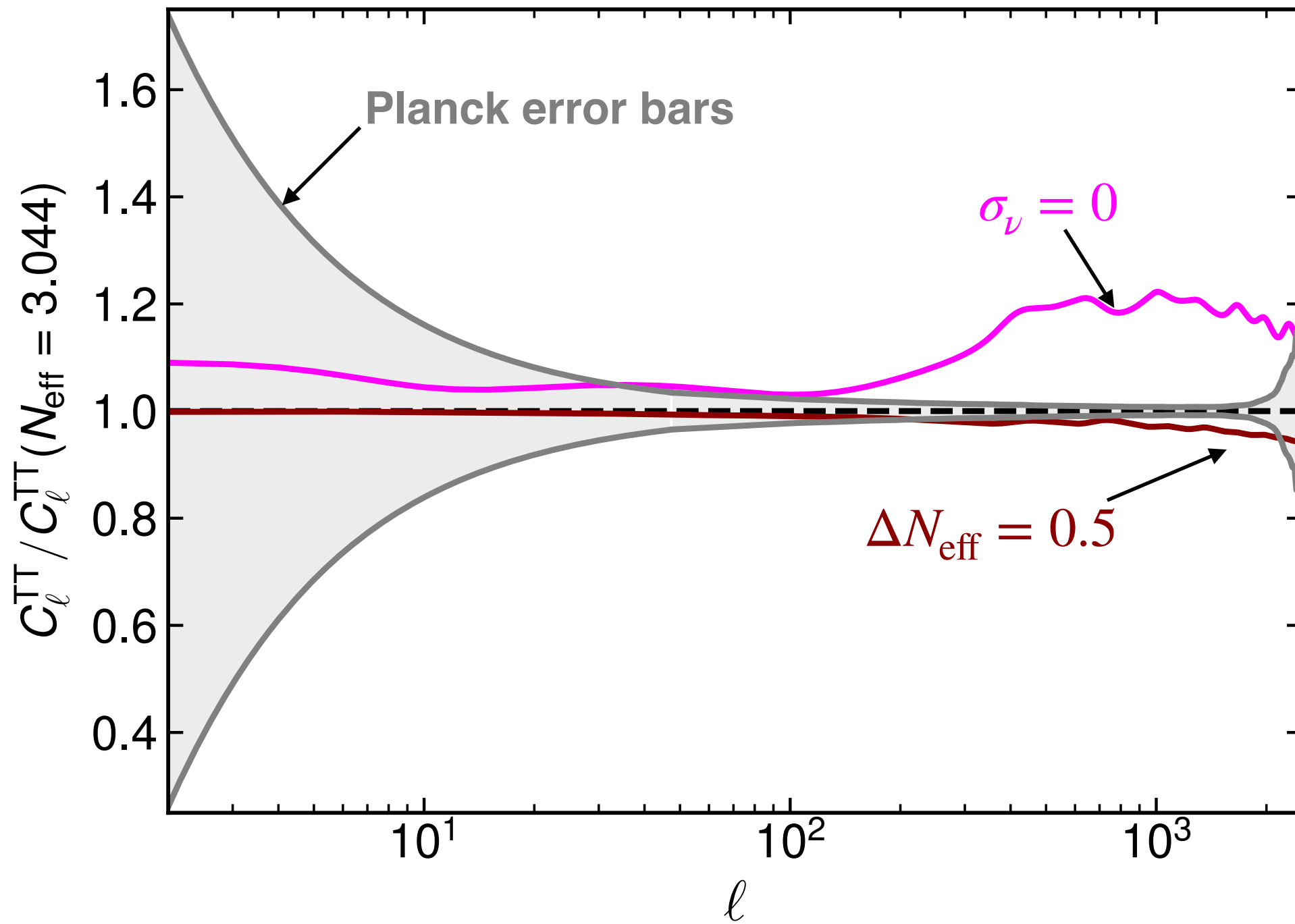
$$\sigma_\nu \longrightarrow \delta g_{\mu\nu} \longrightarrow \Delta T_\gamma$$

Free Streaming Neutrinos $\sigma_\nu \neq 0$

Interacting Neutrinos $\sigma_\nu \rightarrow 0$



Effect of Neutrino Free-streaming in the CMB



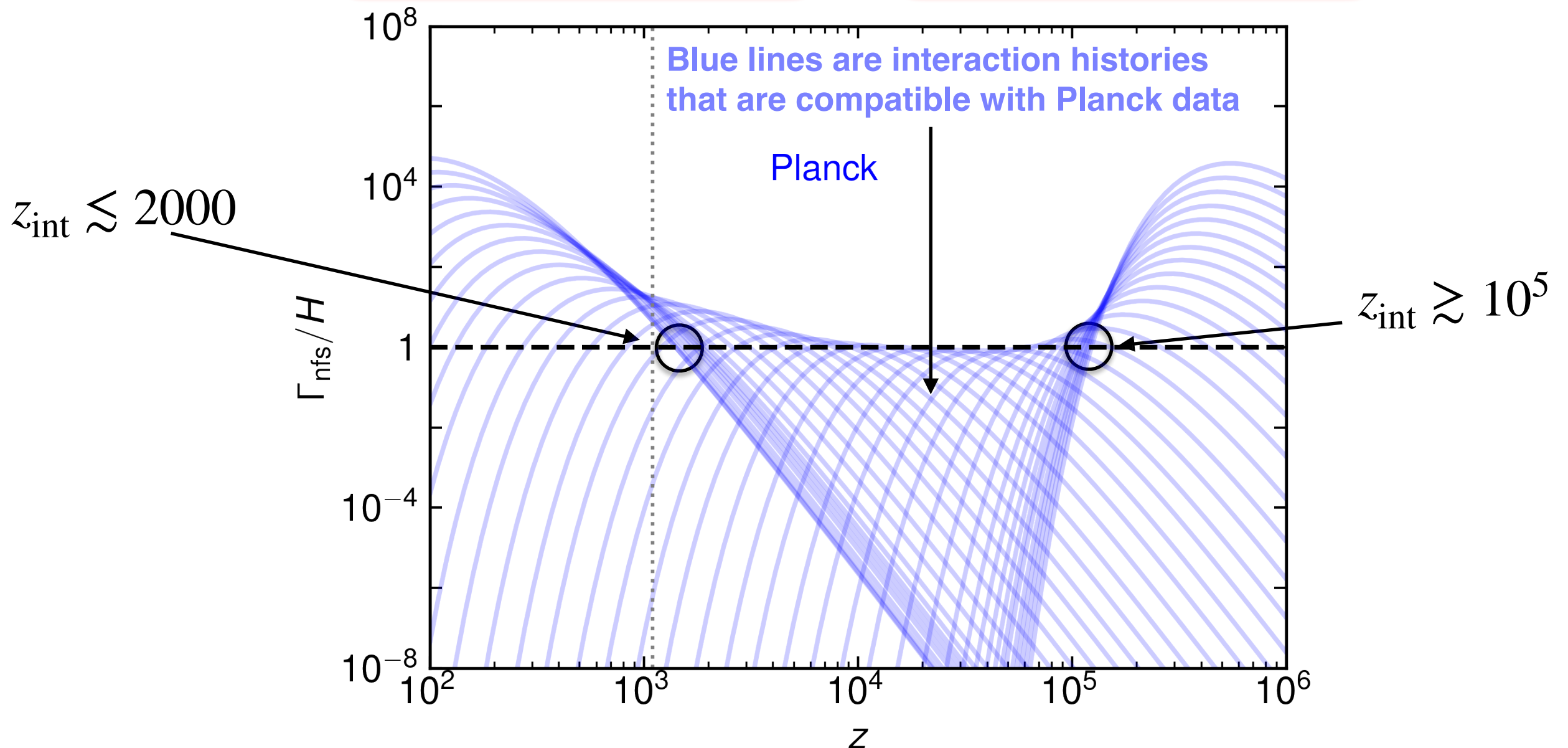
The Neutrino Freestreaming window

Together with Petter Taule and Mathias Garny in 2207.04062 we have recently established the presence of a neutrino free streaming window.

We have demonstrated in a model independent way that neutrinos cannot interact efficiently between themselves or other light particles in the range:

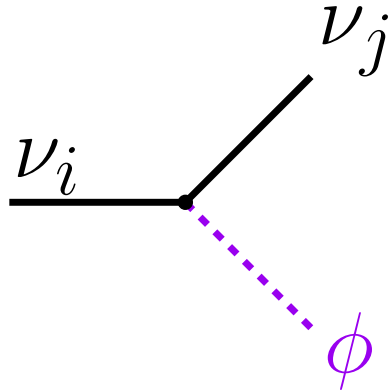
$$2000 \lesssim z_{\text{int}} \lesssim 10^5$$

$$0.3 \text{ eV} \lesssim T_\nu \lesssim 15 \text{ eV}$$



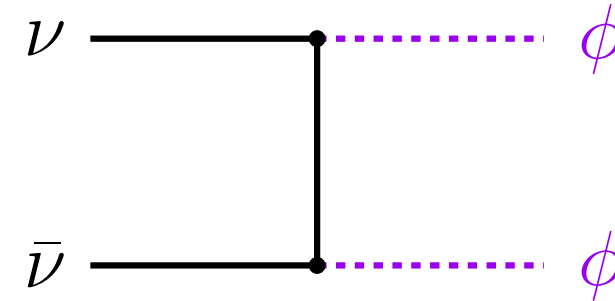
Models

Neutrino Decays



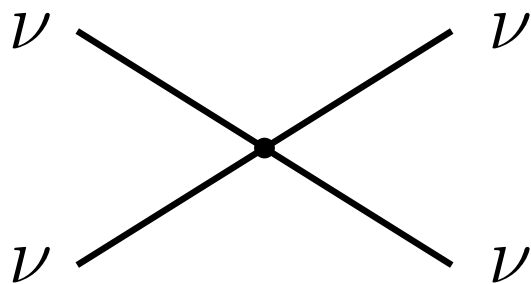
Hannestad & Raffelt [hep-ph/0509278]
Basboll, Bjaelde, Hannestad & Raffelt [0806.1735]
Escudero & Fairbairn [1907.05425]
Chacko, Dev, Du, V. Poulin and Y. Tsai [1909.05275]
Barenboim, Chen, Hannestad, Oldengott, Tram & Wong [2011.01502]
Abellán, Chacko, Dev, Du, Poulin & Tsai [2112.13862]
Chen, Oldengott, Pierobon & Wong [2203.09075]

Neutrino Annihilations



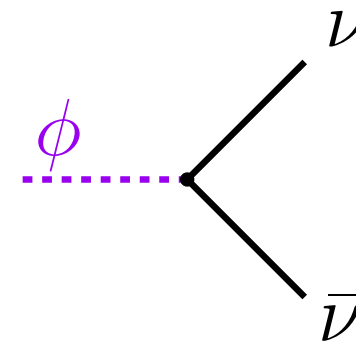
Beacom, Bell & Dodelson [astro-ph/0404585]
Hannestad [astro-ph/0411475]
Archidiacono & Hannestad [1311.3873]
Forastieri, Lattanzi & Natoli [1904.07810]

Neutrino Scatterings



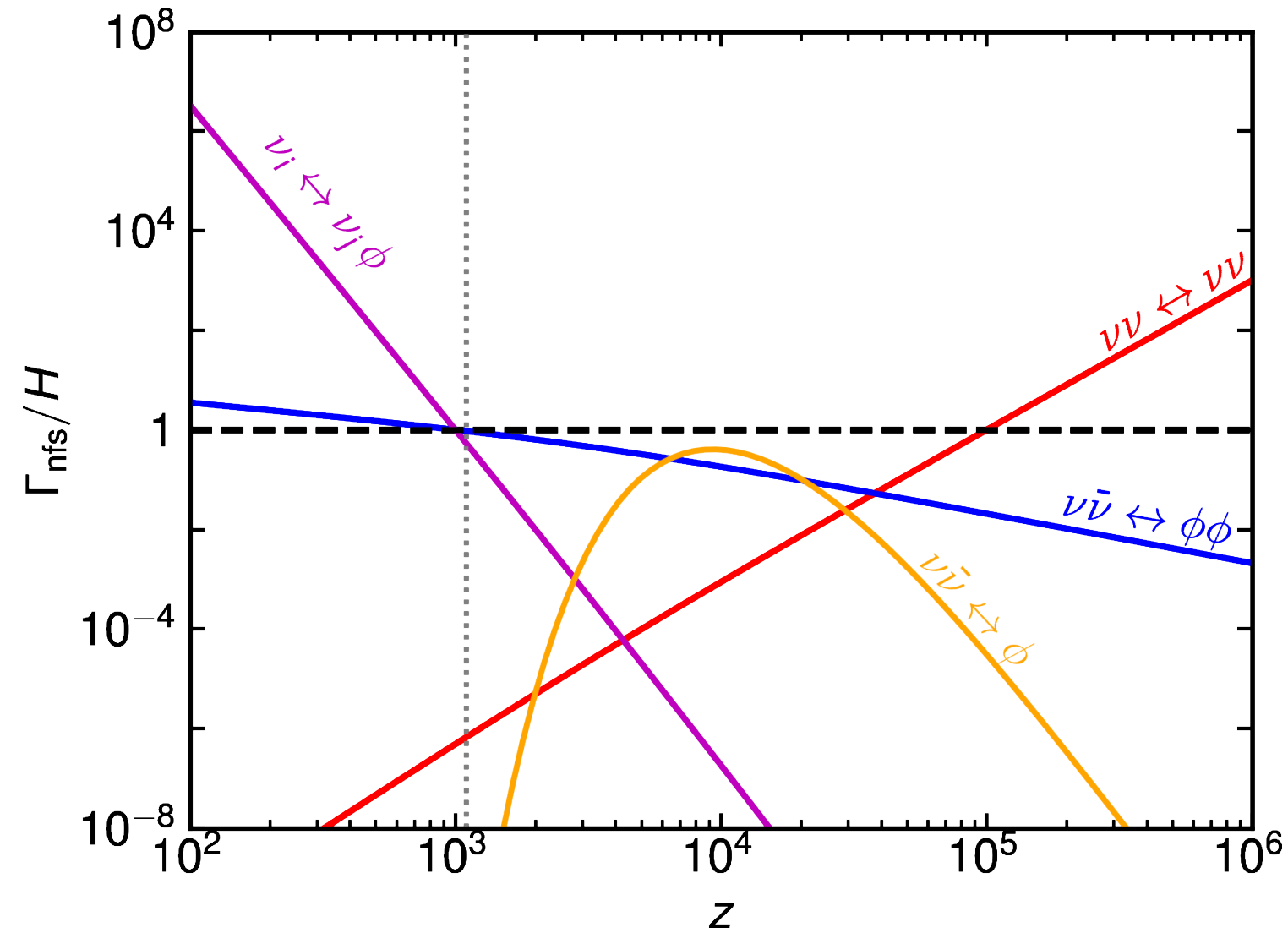
Cyr-Racine & Sigurdson [1306.1536]
Lancaster, Cyr-Racine, Knox & Pan [1704.06657]
Oldengott, Tram, Rampf & Wong [1706.02123]
Kreisch, Cyr-Racine & Doré [1902.00534]
Das & Ghosh [2011.12315]
Choudhury, Hannestad & Tram [2012.07519]
Brinckmann, Chang & LoVerde [2012.11830]

eV-scale neutrinophilic bosons

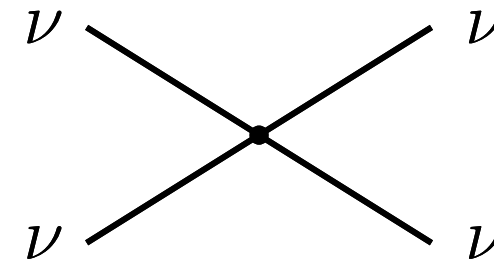


Chacko, Hall, Okui & Oliver [hep-ph/0312267]
Escudero & Witte [1909.04044]
Escudero & Witte [2103.03249]
Sandner, Escudero & Witte [2304.XXXXX]

Rates for various models

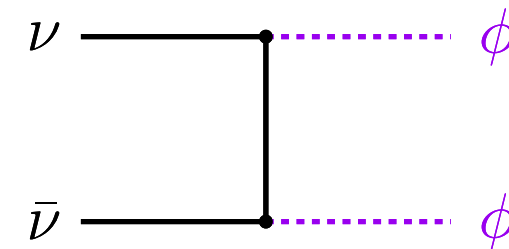


Neutrino scatterings



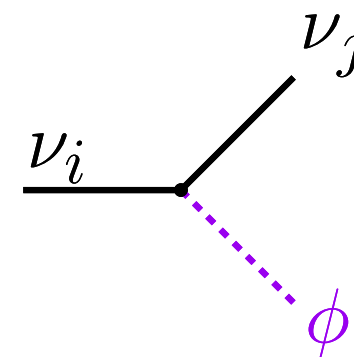
$$\Gamma_{nfs} \sim T^5$$

Neutrino annihilations



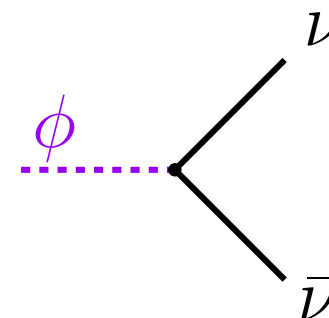
$$\Gamma_{nfs} \sim T$$

Neutrino decays



$$\Gamma_{nfs} \sim T^{-5}$$

eV-scale Neutrinophilic Bosons

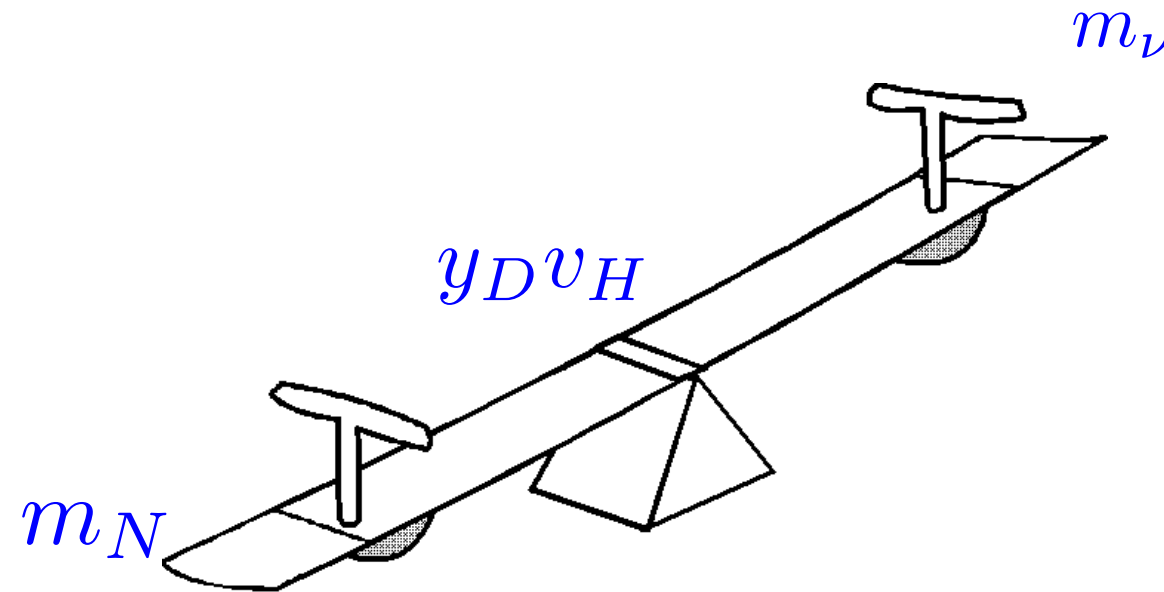


$$\Gamma_{nfs} \sim T^{-5} \quad (T > m_\phi)$$

$$\Gamma_{nfs} \sim e^{-m_\phi/T} \quad (T < m_\phi)$$

The case of the Majoron

Type-I seesaw



Neutrinos are very light Majorana particles: $m_\nu \simeq 0.03 \text{ eV} \left(\frac{y_D}{10^{-6}} \right)^2 \frac{\text{TeV}}{M_N}$

Are There Real Goldstone Bosons Associated with Broken Lepton Number? #1

Y. Chikashige (Munich, Max Planck Inst.), Rabindra N. Mohapatra (Munich, Max Planck Inst. and Munich U.), R.D. Peccei (Munich, Max Planck Inst.) (Sep, 1980)

Published in: *Phys.Lett.B* 98 (1981) 265-268

[DOI](#) [cite](#) [claim](#)

[reference search](#) [1,137 citations](#)

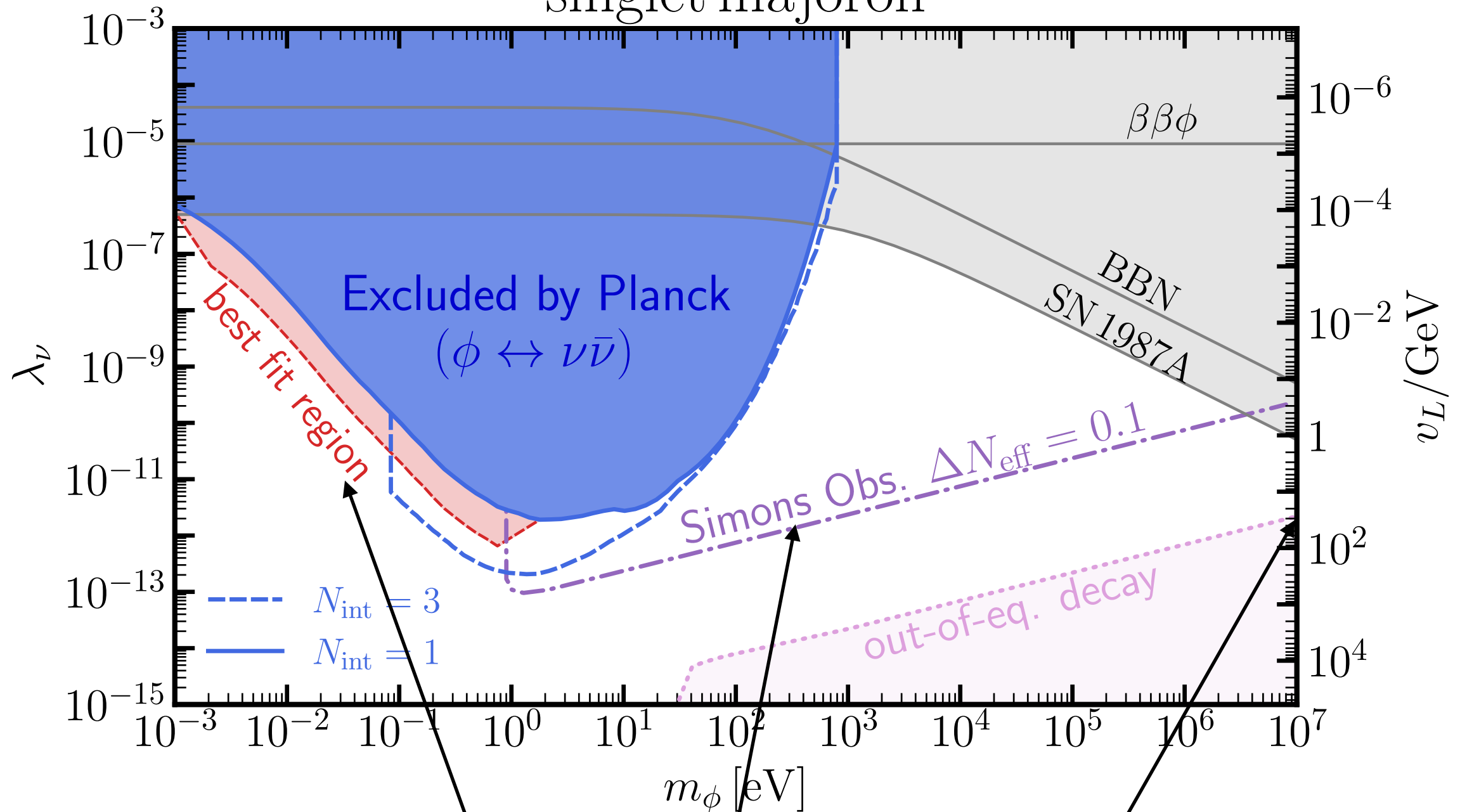
The Majoron is the pseudo-Goldstone boson associated with the spontaneous breaking of global $U(1)_L$

$$\mathcal{L} = \lambda \phi \bar{\nu} \gamma_5 \nu \quad \lambda = m_\nu / v_\phi$$

The case of the Majoron

Sandner, Escudero & Witte 2305.01692

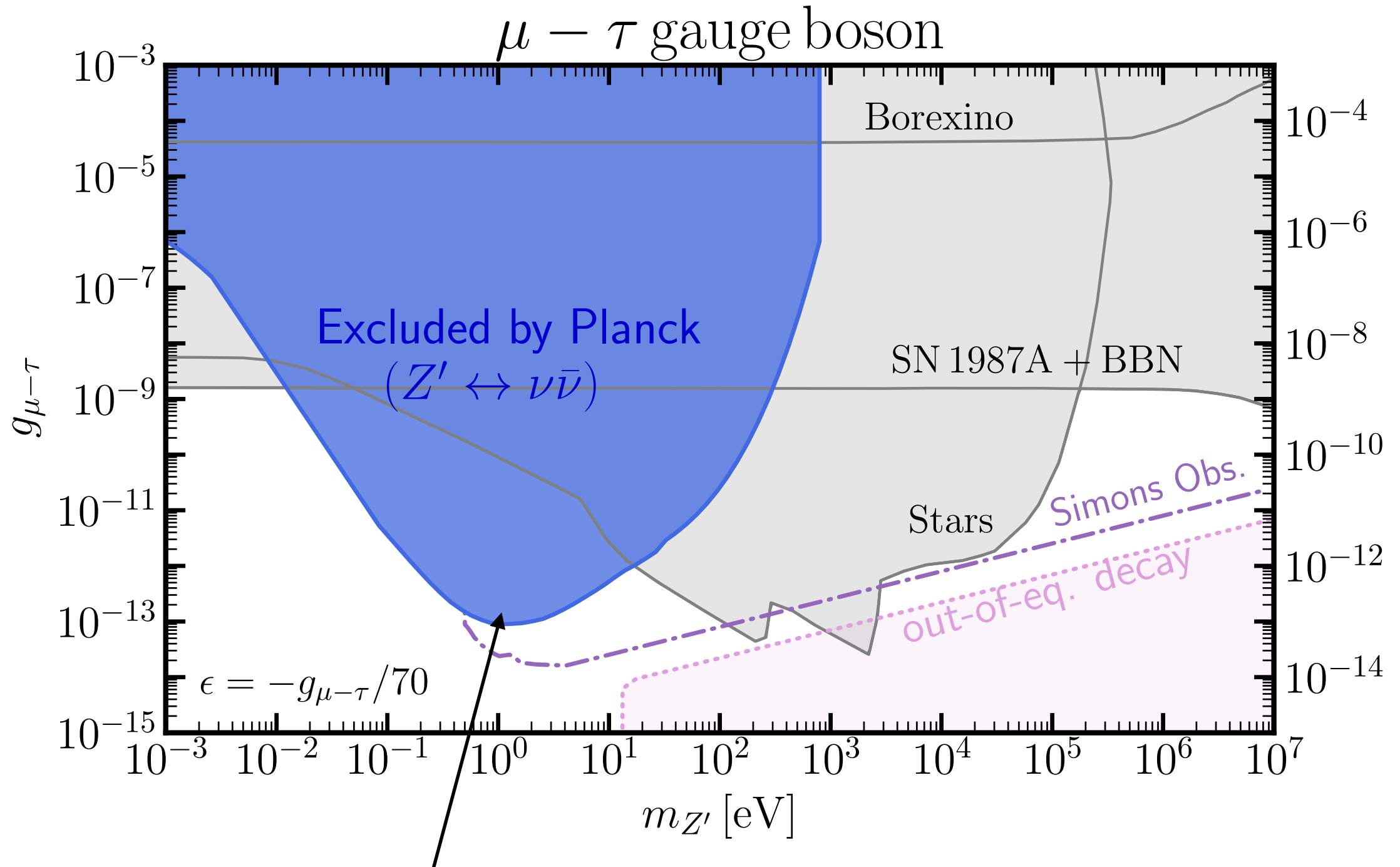
singlet majoron



- **CMB observations can test a well motivated neutrino mass model up to $\nu_L \sim \nu_H$**
- **There is a region preferred at $\sim 1\sigma$. We show that together with a primordial ΔN_{eff} the model can lower the H_0 tension to the 3σ level. This is $0.5-1\sigma$ worse than what we found in Escudero & Witte 1909.04044 and 2103.03249**
- **The Simons Observatory will test in the next ~ 6 years wide regions of parameter space!**

A light $\mu - \tau$ gauge boson

Sandner, Escudero & Witte 2305.01692



- Planck rules out couplings as small as $g_{\mu-\tau} \sim 10^{-13}$ for $m_{Z'} \sim \text{eV}$

Conclusions

Neff: Number of relativistic neutrino species

BBN $N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$

Planck+BAO $N_{\text{eff}}^{\text{CMB}} = 2.99 \pm 0.17$

Standard Model $N_{\text{eff}} = 3.043$

CMB and BBN measurements give strong evidence that the Cosmic Neutrino background should be there.

This implies:

- 1) a stringent constraint on many BSM models**
- 2) gives us confidence to test neutrino properties with cosmology**

We have performed the most accurate calculation of Neff to date and found

$$N_{\text{eff}}|_{\text{SM}} = 3.043$$

Conclusions

Neutrino Masses:

Cosmological bounds are very stringent within Λ CDM:

$$\sum m_\nu < 0.12 \text{ eV} \quad \text{at 95\% CL}$$

In addition, they are robust upon standard modifications of the model.

There are several non-standard neutrino cosmologies where this bound can be evaded

The case of a non-standard CNB to relax them

We developed a simple scenario compatible with high scale type-I seesaw

Need a large number of dark radiation species interacting with neutrinos between BBN and recombination

Parameter space of interest is $m_{Z'} \sim 10 \text{ keV}$ and $\nu_\Phi \sim 10 \text{ MeV} - 1 \text{ GeV}$

As of now a fun model building exercise but could get more relevance if we were to detect something in the lab or nothing in cosmology!

Conclusions

Neutrino Interactions:

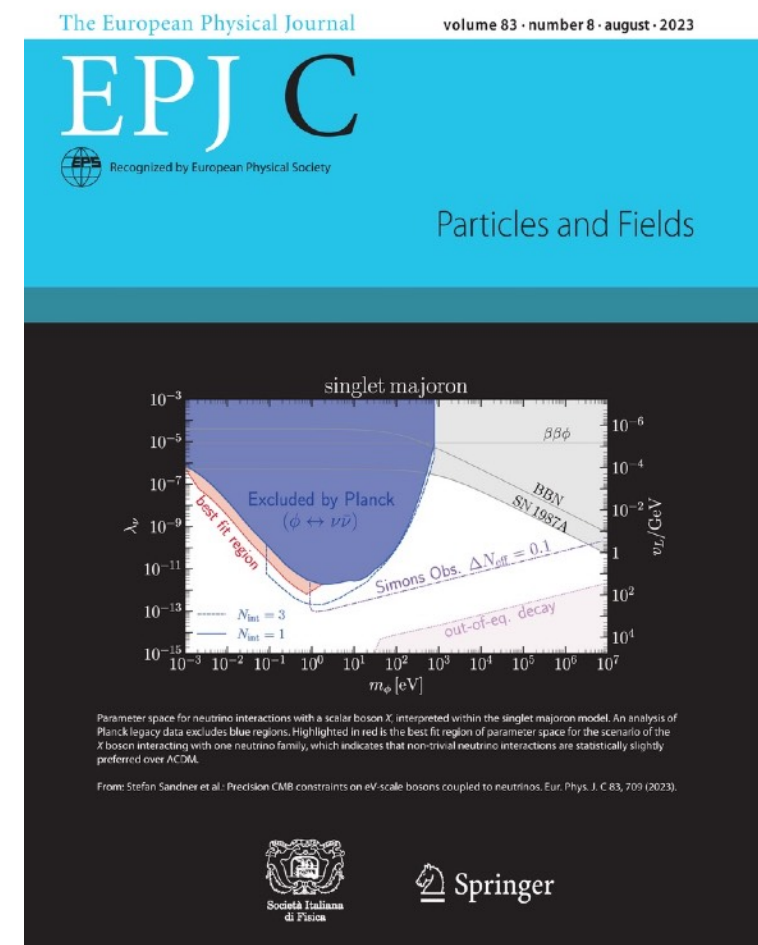
The CMB is a powerful probe of neutrino interactions

We have shown that there is a well defined redshift region where neutrinos must free stream

$$2 \times 10^3 \lesssim z \lesssim 10^5$$

These bounds are relevant for many particle physics scenarios

Including the singlet majoron model and a light mu-tau Z'



Outlook: Number of Neutrinos

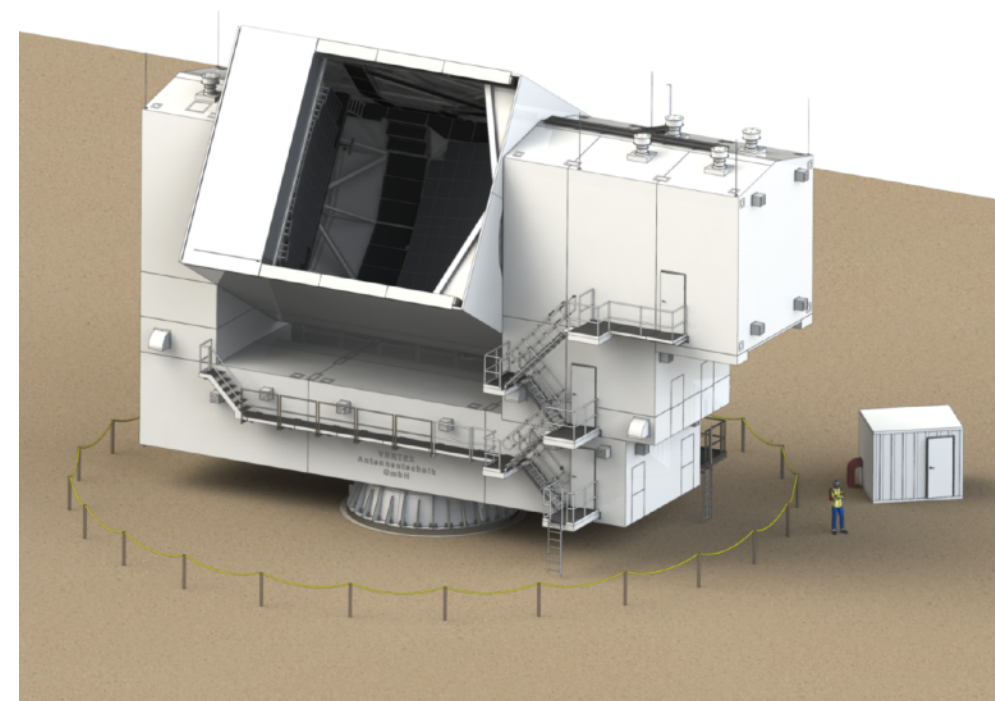
The next generation of CMB experiments are expected to significantly improve the sensitivity to N_{eff}

Simons Observatory



$$\sigma(N_{\text{eff}}) = 0.05 \sim \mathbf{2029}$$

CMB-S4



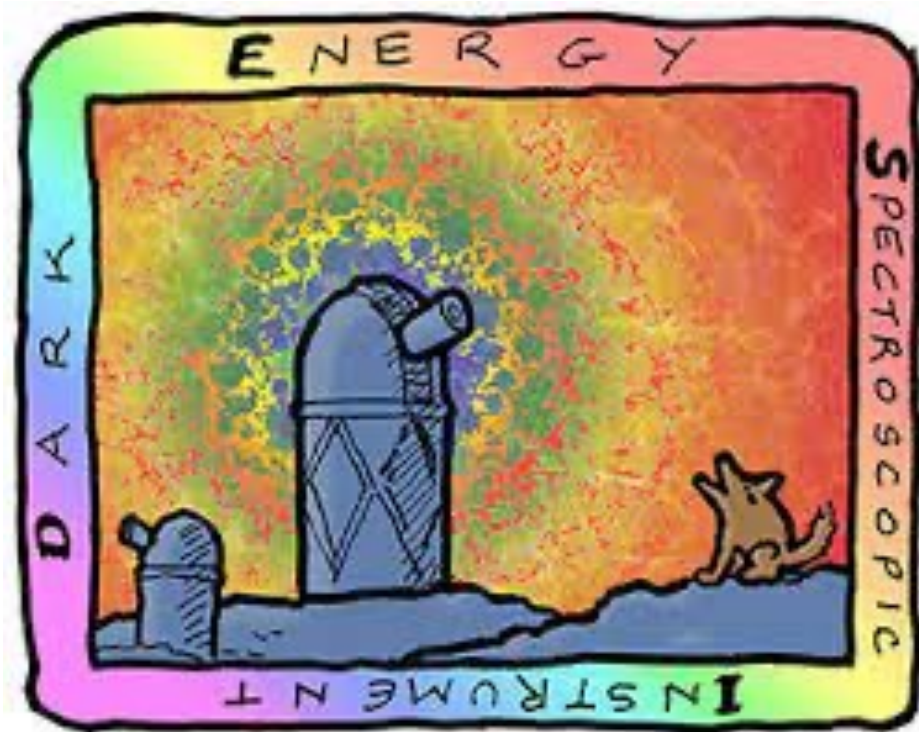
$$\sigma(N_{\text{eff}}) = 0.03 \sim \mathbf{2035?}$$

These measurements will represent an important test of the CNB and BBN in the SM and perhaps may yield a BSM signal!

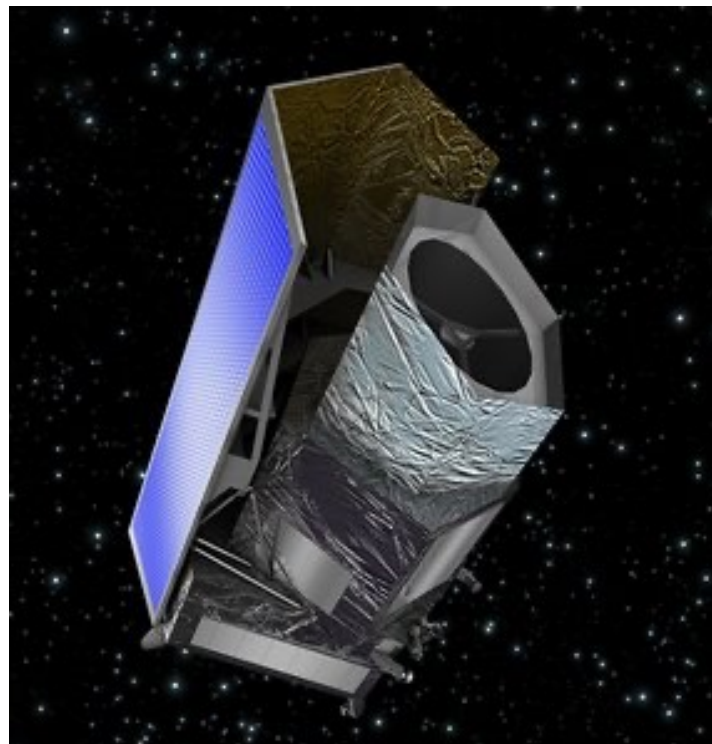
Outlook: Neutrino Masses

The next generation of galaxy surveys in combination with CMB data are expected to measure the neutrino mass if the Universe is governed by a Λ CDM cosmology

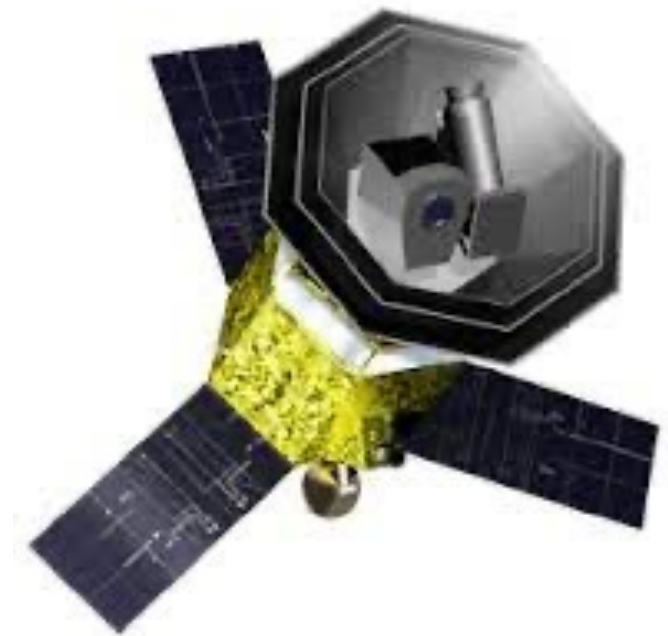
DESI



EUCLID



LiteBIRD



Why? DESI: 30M galaxies and EUCLID: 50M galaxies, but BOSS 1M galaxies

This is expected to happen in the next 3-4 years: $\sigma(\sum m_\nu) = 0.02$

In parallel, the KATRIN experiment is taking data and should reach a sensitivity of $m_{\bar{\nu}_e} \lesssim 0.2 \text{ eV}$ at 90% CL in $\sim 3\text{-}4$ years.

Hubble tension?

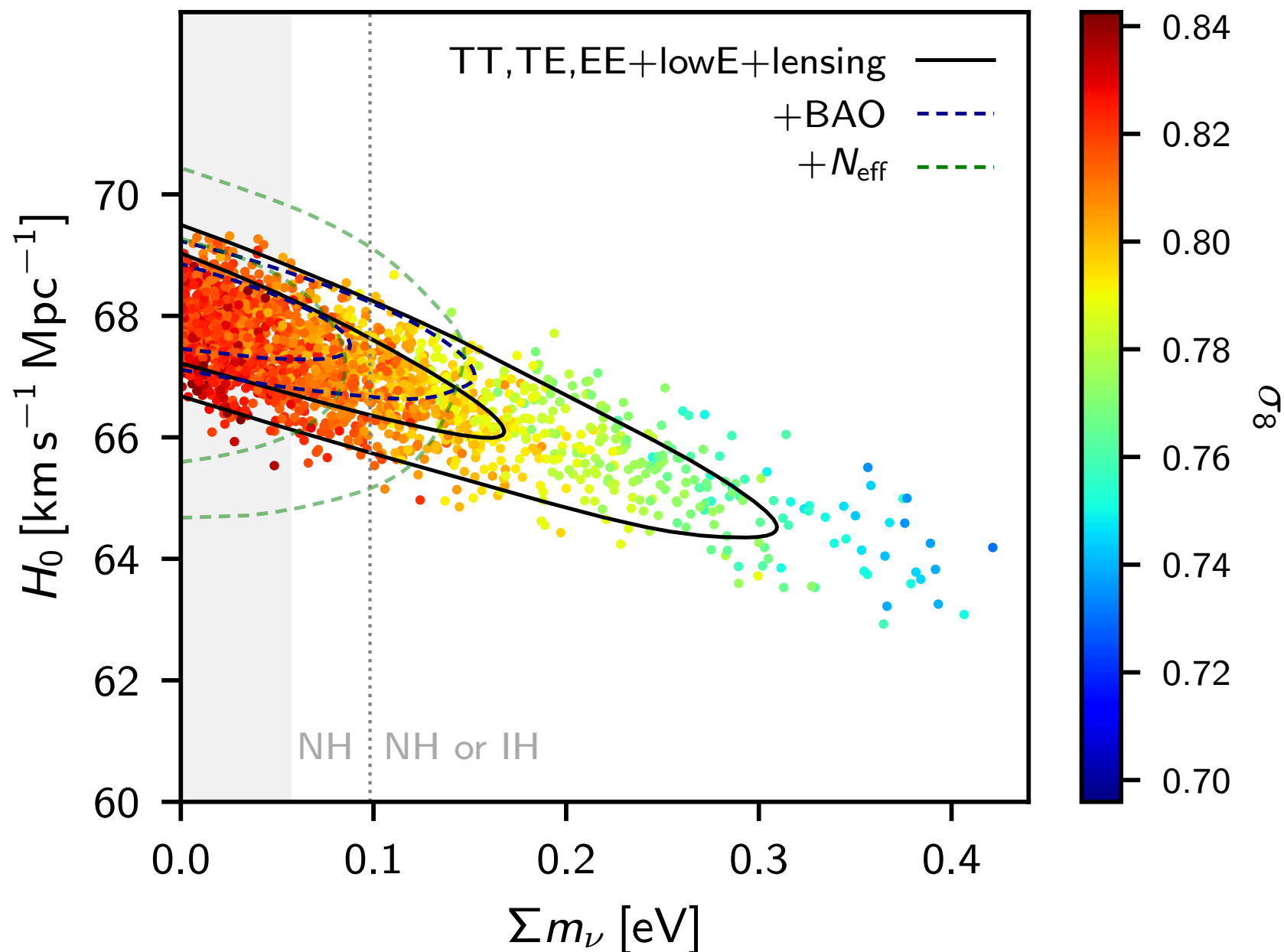
Local:

$$H_0 = 73.0 \pm 1.0 \text{ km/s/Mpc}$$

CMB+BAO:

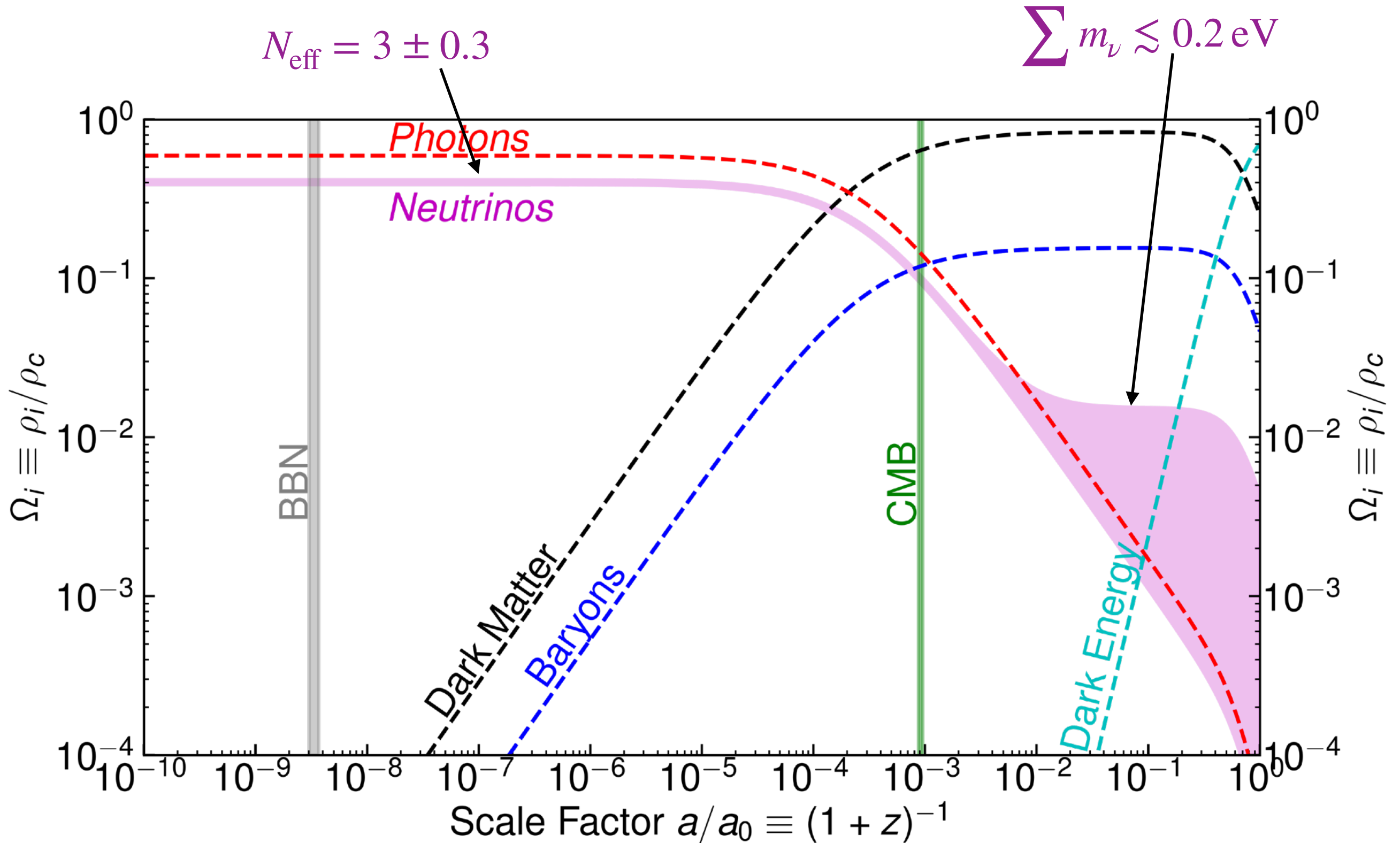
$$H_0 = 67.7 \pm 0.4 \text{ km/s/Mpc}$$

5σ discrepancy!



Global Perspective

Current knowledge:

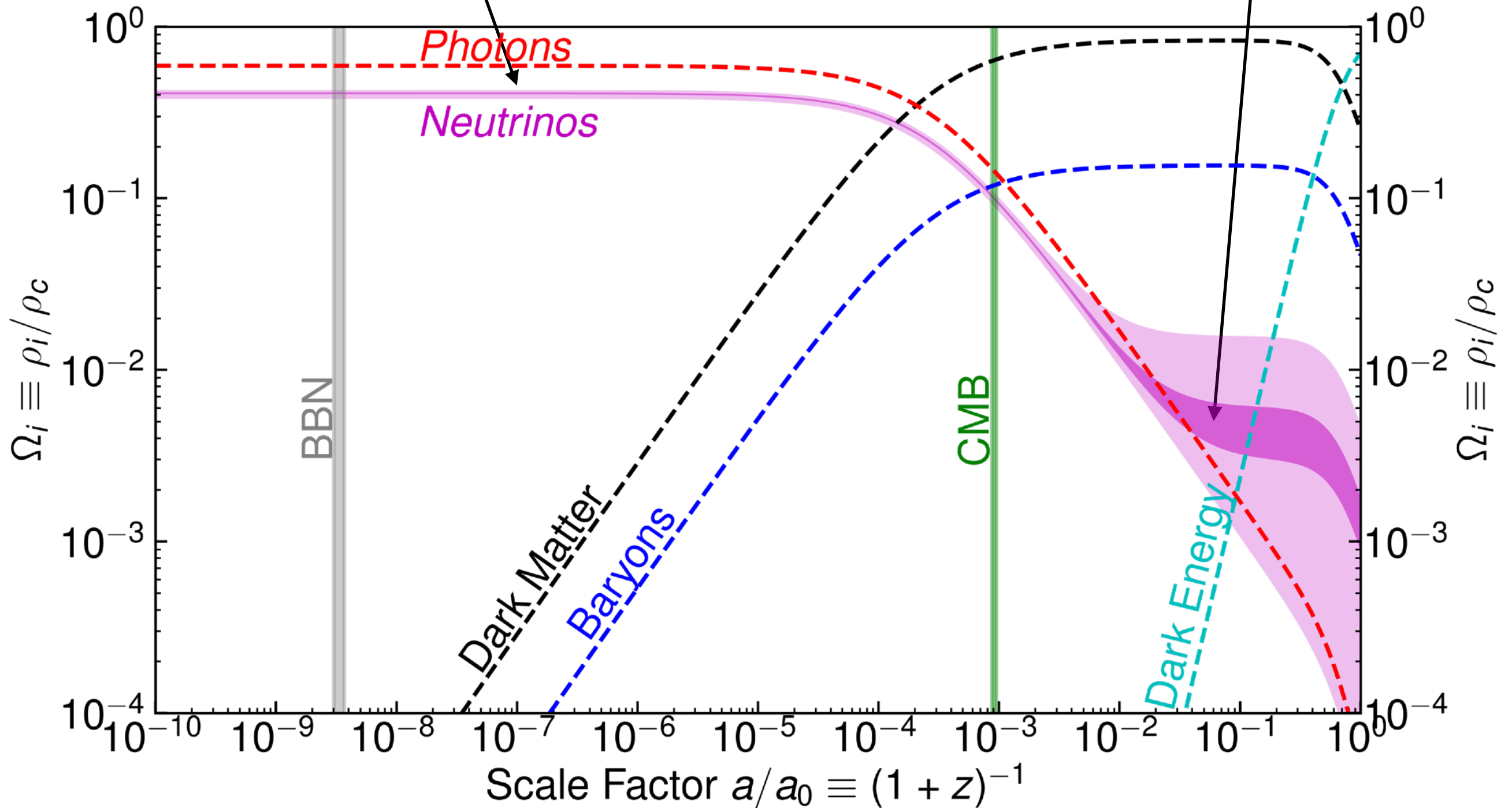


Global Perspective

In the next 5-6 years:

$$N_{\text{eff}} = 3.043 \pm 0.06$$

$$\sum m_\nu = 0.06 \pm 0.02 \text{ eV}$$



Time for Questions and Comments

Upcoming years are going to be exciting!



Thank you for your attention!

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